

# Rapidity dependence study of charged pion production in relativistic nuclear collisions using Tsallis distribution

Oana RISTEA<sup>1</sup>, Catalin RISTEA<sup>1,2</sup>, Alexandru JIPA<sup>1</sup>,  
Marius CALIN<sup>1</sup>, Tiberiu ESANU<sup>1,3</sup>

1 University of Bucharest, Faculty of Physics, Bucharest-Magurele, Romania

2 Institute of Space Science, Bucharest-Magurele, Romania

3 National Institute of Nuclear Physics and Engineering Horia Hulubei, Bucharest-Magurele, Romania

# Outline

---

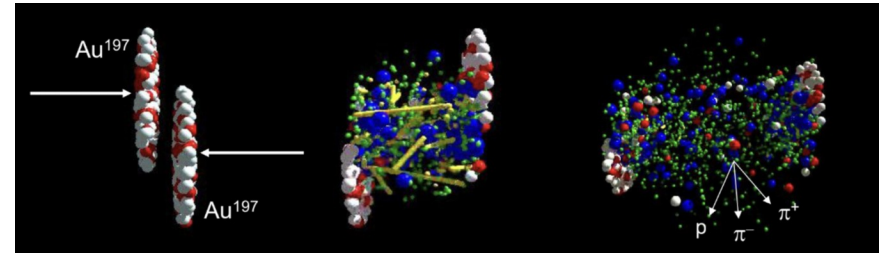
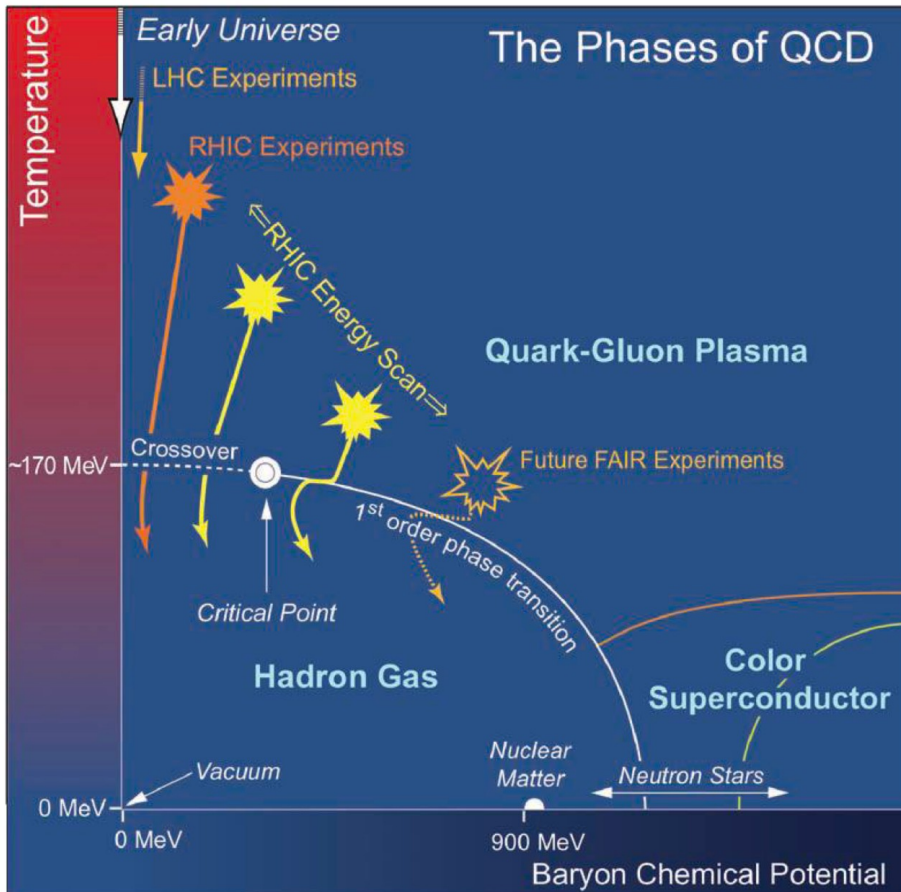
Motivation

Tsallis distribution and data used

Results

Conclusions

# Heavy-ion collision evolution



→ By changing the energy available in the collision and the projectile-target combinations, one can obtain systems characterized by various temperatures and baryon chemical potentials

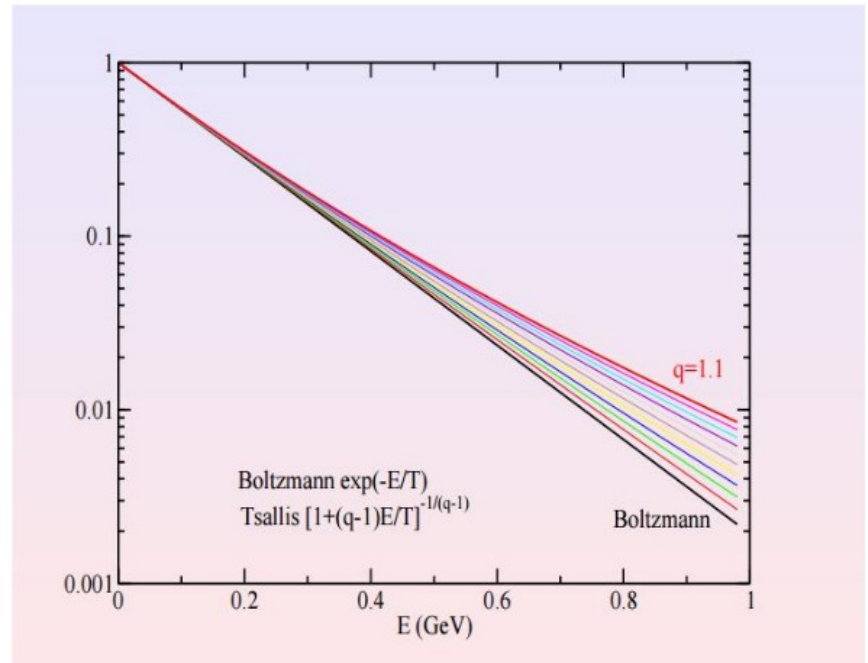
→ different regions on the phase diagram can be investigated

# Tsallis distribution

$$\frac{d^2N}{p_T dp_T dy} = gV \frac{m_T \cosh(y)}{(2\pi)^2} \times \left[ 1 + (q-1) \frac{m_T \cosh(y) - \mu}{T} \right]^{\frac{q}{1-q}}$$

where  $g$  is the degeneracy factor,  $V$  is the volume,  $y$  is the rapidity,  $m$  is the particle rest mass and  $m_T$  is the transverse mass.

→ The  $T$  parameter can be interpreted as an effective temperature at kinetic freeze-out



→ The non-extensivity parameter,  $q$  → the deviation from the standard Boltzmann distribution. This parameter can be related to the non-equilibrium degree of the system and intrinsic fluctuations in the hadronic medium

→  $q = 1$  → thermal equilibrium

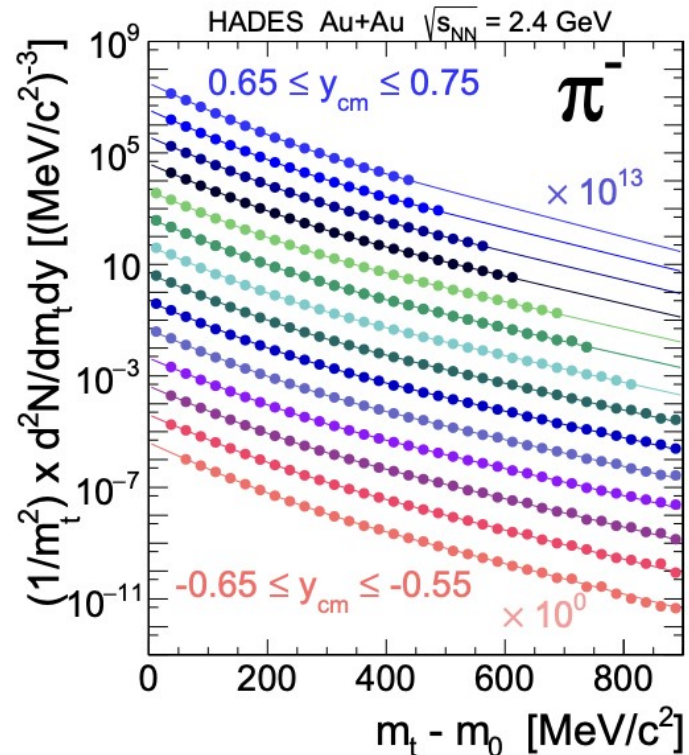
→  $q > 1$  → non-equilibrium dynamics, temperature fluctuations

# Tsallis distribution

$$\frac{d^2N}{p_T dp_T dy} = gV \frac{m_T \cosh(y)}{(2\pi)^2} \times \left[ 1 + (q-1) \frac{m_T \cosh(y) - \mu}{T} \right]^{\frac{q}{1-q}}$$

where  $g$  is the degeneracy factor,  $V$  is the volume,  $y$  is the rapidity,  $m$  is the particle rest mass and  $m_T$  is the transverse mass.

→ The  $T$  parameter can be interpreted as an effective temperature at kinetic freeze-out

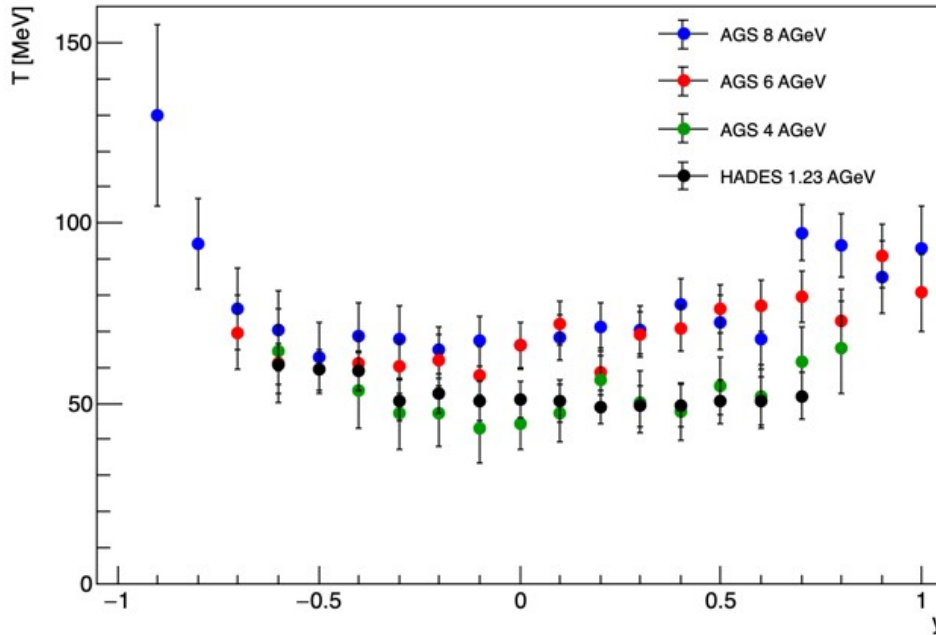
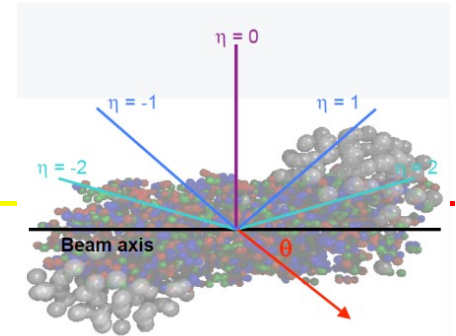


→ The non-extensivity parameter,  $q$  → the deviation from the standard Boltzmann distribution. This parameter can be related to the non-equilibrium degree of the system and intrinsic fluctuations in the hadronic medium

→  $q = 1$  → thermal equilibrium

→  $q > 1$  → non-equilibrium dynamics, temperature fluctuations

# Tsallis temperature vs rapidity



→ AGS data: Au+Au at  $E_{\text{beam}} = 4, 6$  and  $8$  AGeV ( $\sqrt{s_{\text{NN}}} = 3, 3.6$  and  $4.1$  GeV);

→ HADES data: Au+Au at  $E_{\text{beam}} = 1.23$  GeV ( $\sqrt{s_{\text{NN}}} = 2.4$  GeV);

J.L. Klay, et al., E895 Coll., Phys.Rev.C68 (2003)054905

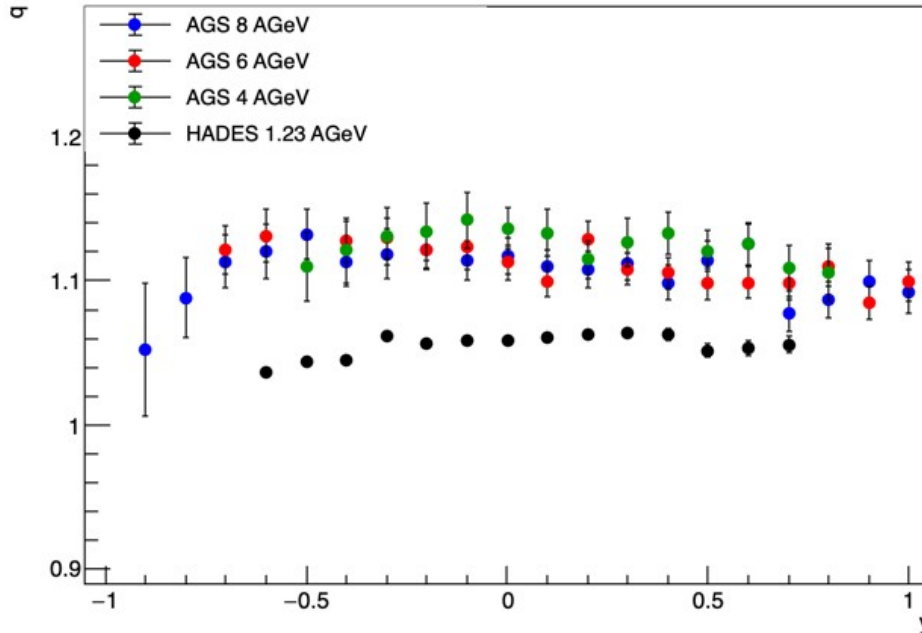
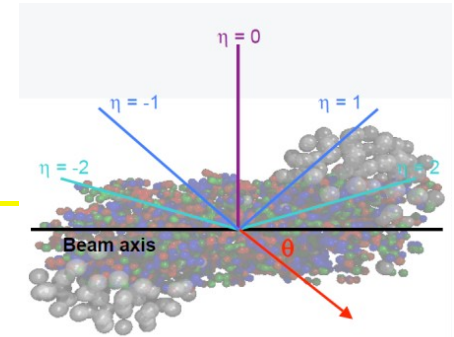
J. Adamczewski-Musch, et al., HADES Coll., Eur.Phys.J.A 56 (2020) 10, 259

→ Tsallis temperature increases from midrapidity to forward rapidities for all studied energies

→ the shape of  $m_T$  spectra is changing with rapidity → spectra are becoming harder in the forward region compared to the central region

→ at forward rapidities, particle emission is increasingly influenced by projectile fragmentation and other non-equilibrium processes

# Non-extensivity parameter vs rapidity



→  $q$  values are larger in the midrapidity region → a decreasing  $q$  with increasing  $y$  suggests that temperature fluctuations become weaker away from midrapidity

→ at AGS, the central rapidity region experiences the largest deviations from thermal equilibrium, associated with the highest particle density, intense rescattering, resonance production, and fluctuations.

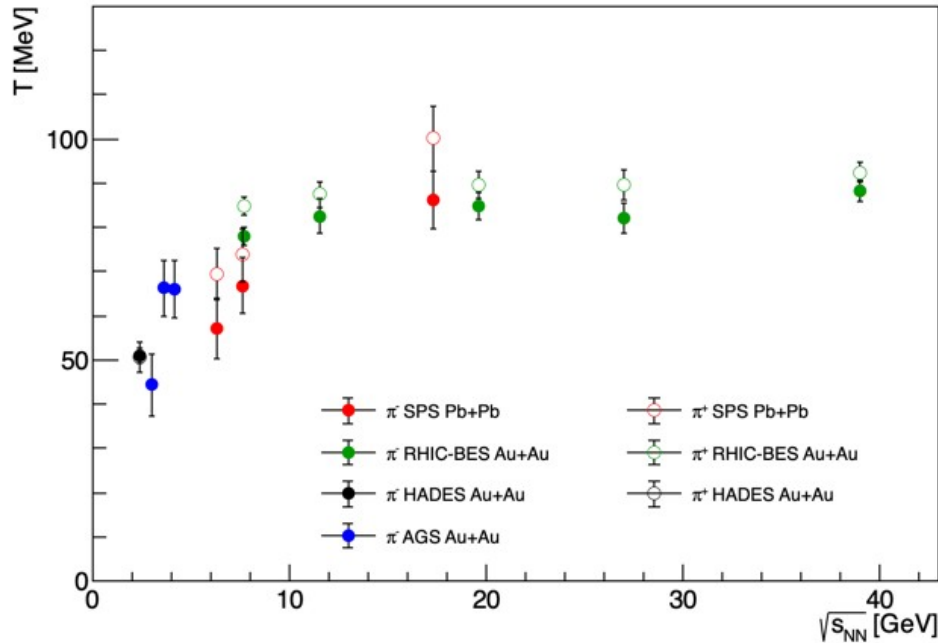
→ AGS data: Au+Au at  $E_{\text{beam}} = 4, 6$  and  $8$  AGeV ( $\sqrt{s_{\text{NN}}} = 3, 3.6$  and  $4.1$  GeV);

→ HADES data: Au+Au at  $E_{\text{beam}} = 1.23$  GeV ( $\sqrt{s_{\text{NN}}} = 2.4$  GeV);

J.L. Klay, et al., E895 Coll., Phys.Rev.C68 (2003)054905

J. Adamczewski-Musch, et al., HADES Coll., Eur.Phys.J.A 56 (2020) 10, 259

# Tsallis temperature vs energy



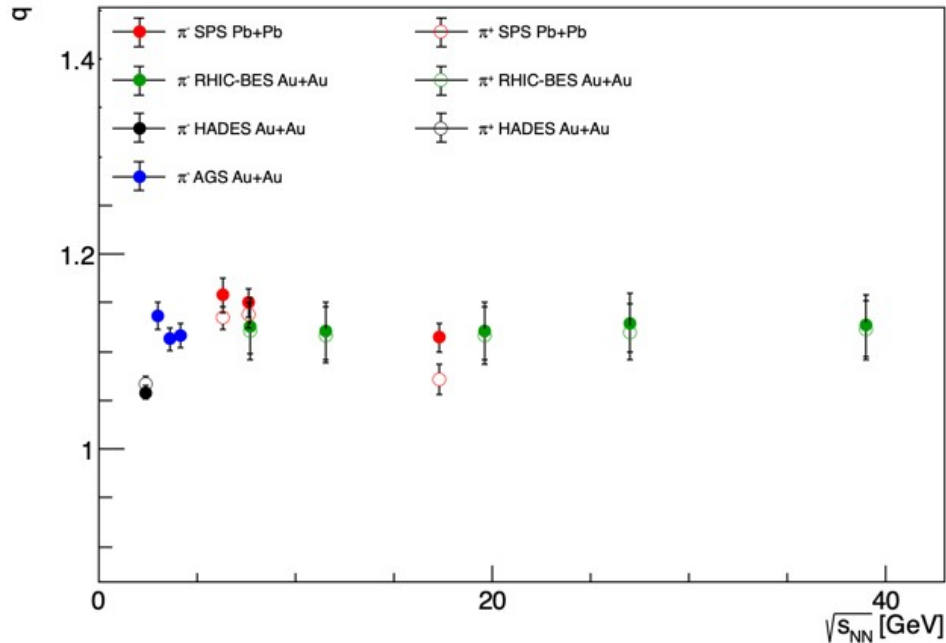
→  $T$  increases from HADES and AGS energies and saturates above  $\sqrt{s_{NN}} \sim 11$  GeV

→ this behaviour could be related to the transition between the baryon-rich regime explored at AGS and HADES energies and the higher-energy regime explored at the RHIC Beam Energy Scan energies

AGS data: Au+Au at  $E_{\text{beam}} = 4, 6$  and  $8$  AGeV ( $\sqrt{s_{NN}} = 3, 3.6$  and  $4.1$  GeV); HADES data: Au+Au at  $E_{\text{beam}} = 1.23$  GeV ( $\sqrt{s_{NN}} = 2.4$  GeV); SPS data: Pb+Pb data at  $E_{\text{beam}} = 20, 30$  and  $158$  AGeV ( $\sqrt{s_{NN}} = 6.3, 7.6$  and  $17$  GeV) and STAR-BES data: Au+Au at  $\sqrt{s_{NN}} = 7.7, 11.5, 19.6, 27$  and  $39$  GeV

J.L. Klay, et al., E895 Coll., Phys.Rev.C68 (2003)054905; J. Adamczewski-Musch, et al., HADES Coll., Eur.Phys.J.A 56 (2020) 10, 259; L. Adamczyk et al., Phys. Rev. C 96, 044904 (2017); S. Afanasiev et al. (NA49 Collaboration), Phys. Rev. C 66, 054902 (2002); I. Bearden et al. (NA44 Collaboration), Phys. Rev. C 66, 044907 (2002)

# q parameter vs energy



→ q quantifies deviations from thermal equilibrium, reflecting fluctuations, correlations, and non-equilibrium dynamics.

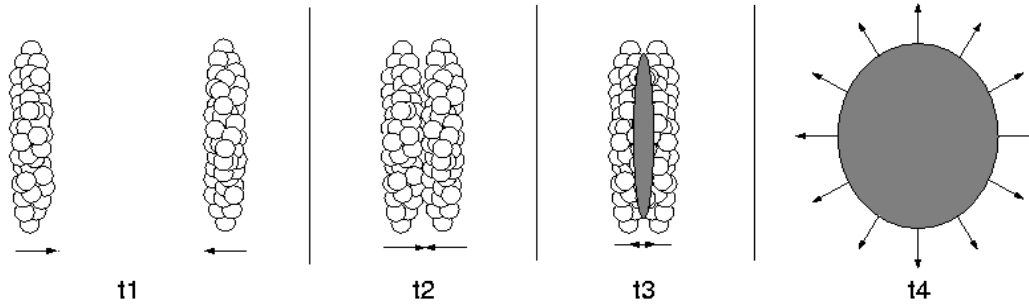
→ at RHIC-BES energies the degree of non-equilibrium no longer increases despite the higher collision energy.

AGS data: Au+Au at  $E_{\text{beam}} = 4, 6$  and  $8$  AGeV ( $\sqrt{s_{NN}} = 3, 3.6$  and  $4.1$  GeV); HADES data: Au+Au at  $E_{\text{beam}} = 1.23$  GeV ( $\sqrt{s_{NN}} = 2.4$  GeV); SPS data: Pb+Pb data at  $E_{\text{beam}} = 20, 30$  and  $158$  AGeV ( $\sqrt{s_{NN}} = 6.3, 7.6$  and  $17$  GeV) and STAR-BES data: Au+Au at  $\sqrt{s_{NN}} = 7.7, 11.5, 19.6, 27$  and  $39$  GeV

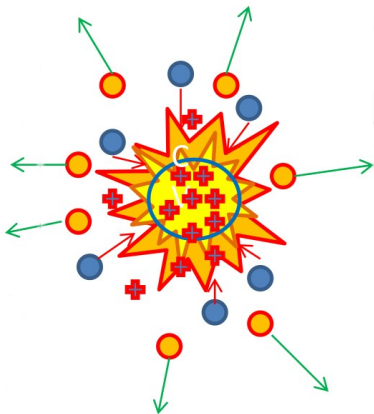
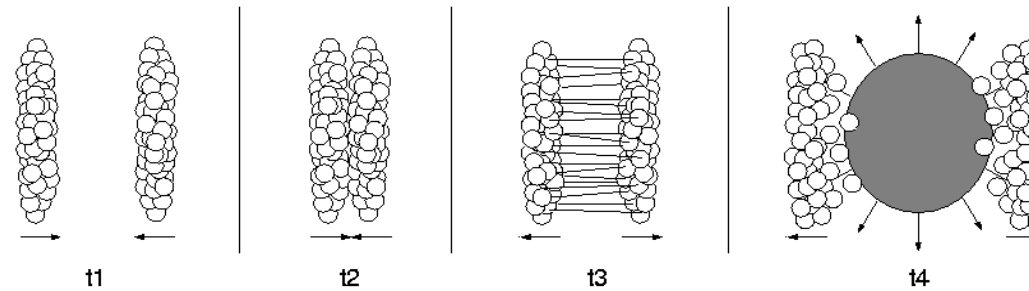
J.L. Klay, et al., E895 Coll., Phys.Rev.C68 (2003)054905; J. Adamczewski-Musch, et al., HADES Coll., Eur.Phys.J.A 56 (2020) 10, 259; L. Adamczyk et al., Phys. Rev. C 96, 044904 (2017); S. Afanasiev et al. (NA49 Collaboration), Phys. Rev. C 66, 054902 (2002); I. Bearden et al. (NA44 Collaboration), Phys. Rev. C 66, 044907 (2002)

# From full stopping to transparency...

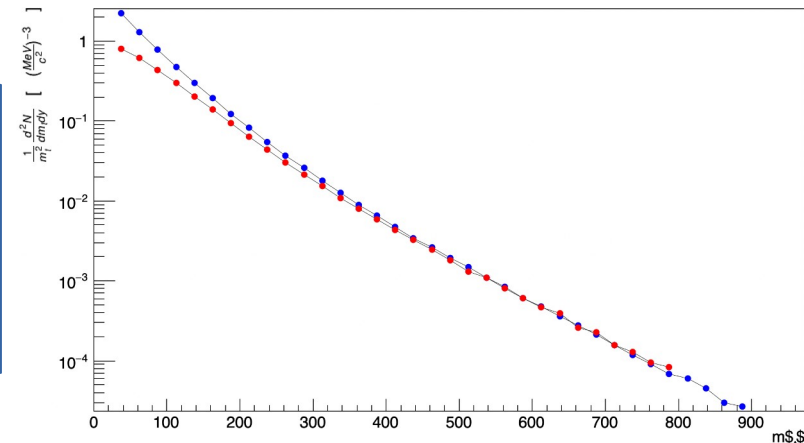
Full stopping:



Transparency



Coulomb interaction between the pions and positive net-charge →  $\pi^+$  ( $\pi^-$ ) are accelerated (decelerated) by the Coulomb field generated by the positive charge



## Model\*:

- considers that during the fireball expansion, a charged pion will receive a momentum change due to the Coulomb interaction or “Coulomb kick”,  $p_c$

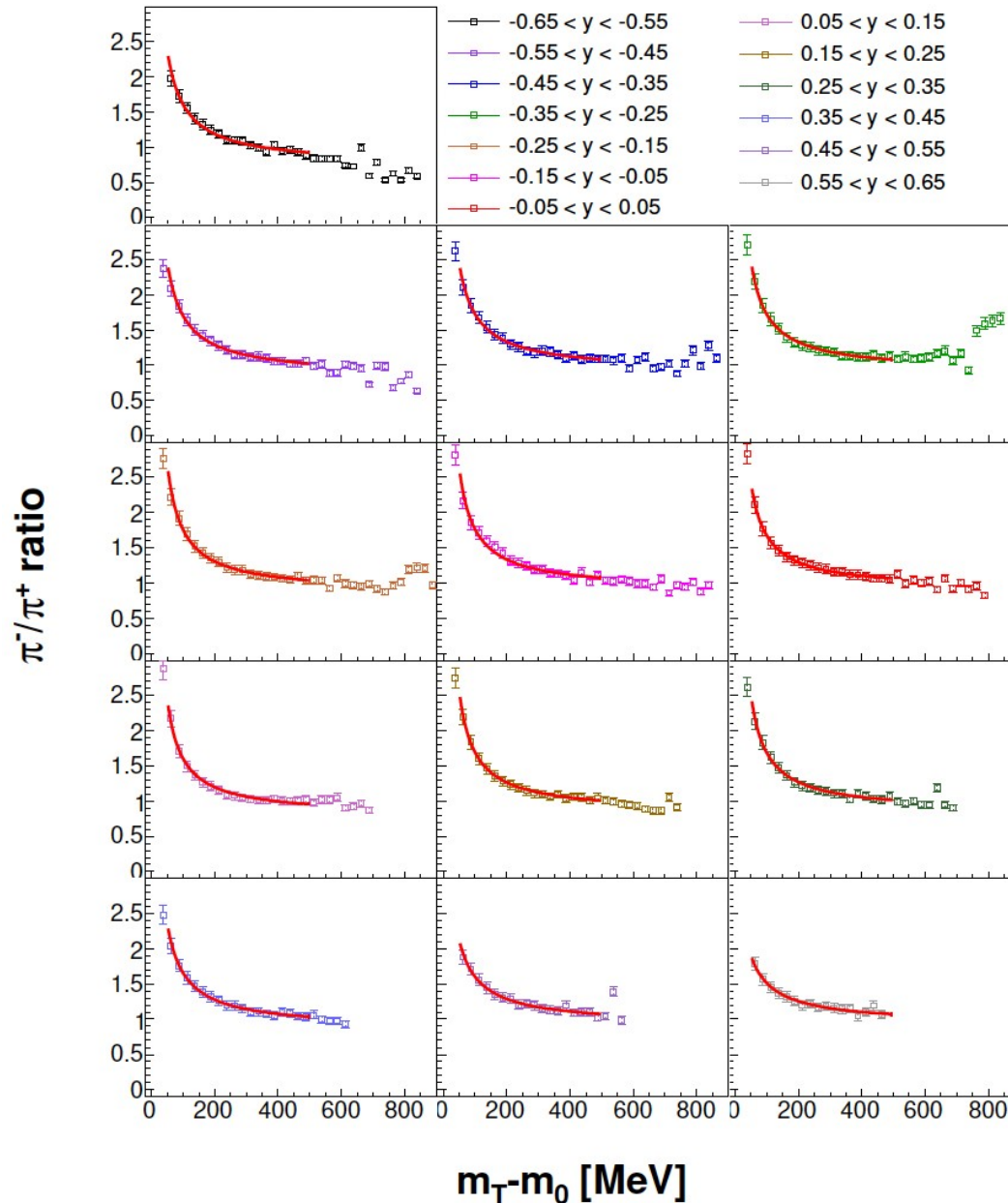
$$p_c \equiv |p_T - p_{T,0}| \cong 2e^2 \frac{dN^{\text{ch}}}{dy} \frac{1}{R_f}$$

where  $p_{T,0}$  is the transverse momentum at freeze-out,  $p_T$  is the final transverse momentum,  $dN_{\text{ch}}/dy$  is the net-charge distribution and  $R_f$  is the kinetic freeze-out radius. The charged pion ratio is:

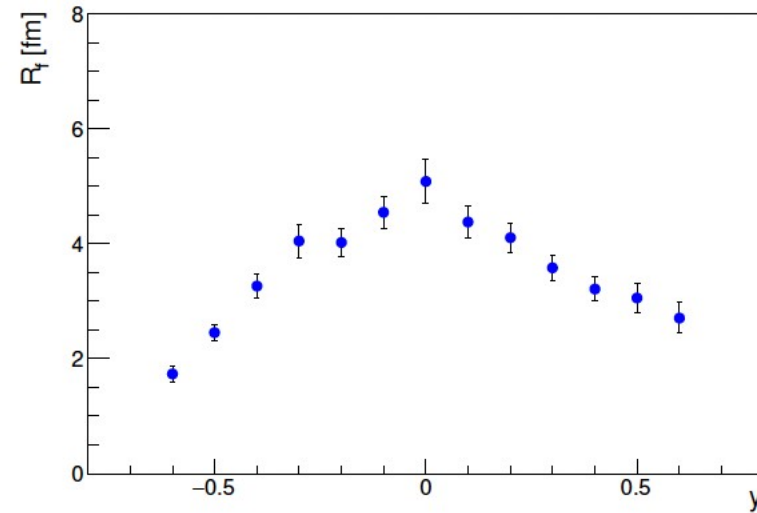
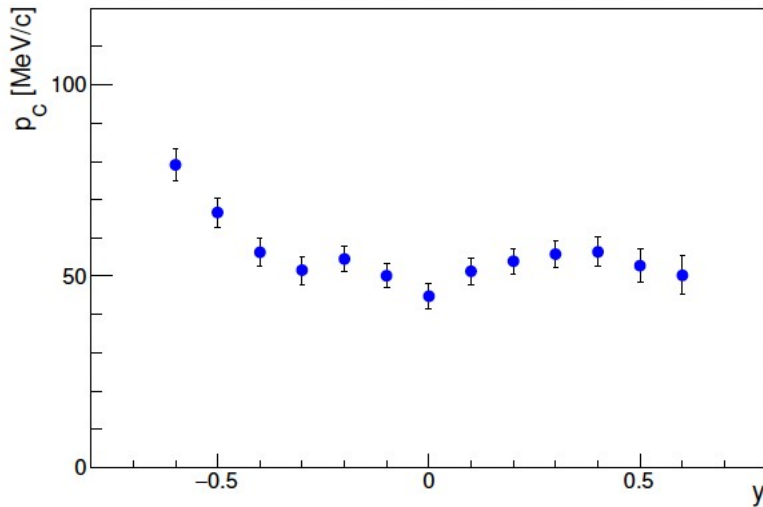
$$\frac{\pi^-}{\pi^+} = \left\langle \frac{\pi^-}{\pi^+} \right\rangle \frac{p_T + p_c}{p_T - p_c} \exp\left(\frac{m_T^- - m_T^+}{T}\right)$$

where  $T$  thermal freeze-out temperature,  $\langle \pi^-/\pi^+ \rangle$  initial pion ratio and  $m_T^\pm = \sqrt{m^2 + (p_T \pm p_c)^2}$

\*H. W. Barz, J. P. Bondorf, J. J. Gaardhøje, and H. Heiselberg, Phys. Rev. C 57 (1998)2536–2546; H. Heiselberg, Nuclear Physics A 638 (1998) 479C



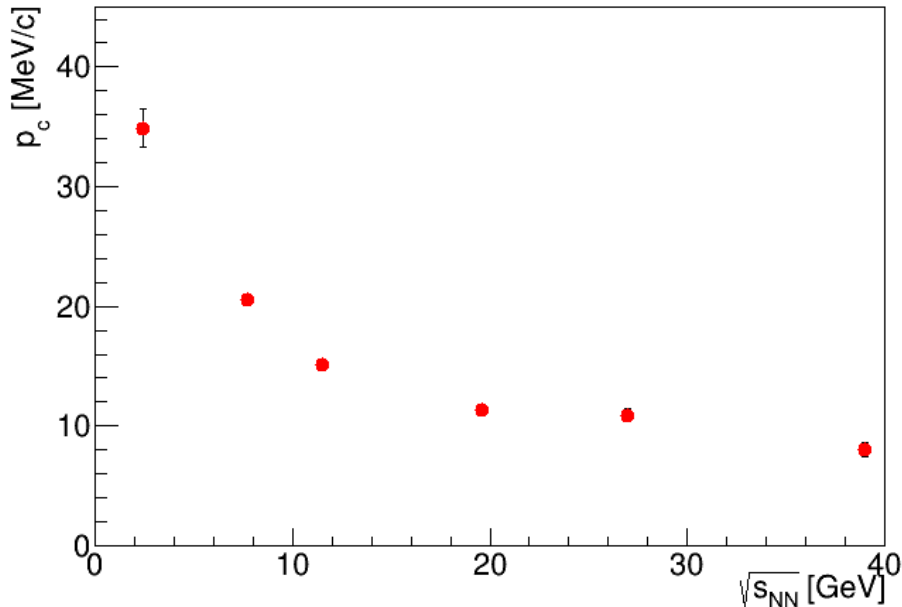
# Rapidity dependence



→ HADES  $m_T$  spectra of charged pions produced in Au+Au collisions at 2.4 GeV → Coulomb momentum has the smallest values near the mid-rapidity and increases towards large rapidities → the Coulomb effects on the charged pion spectra increase at larger rapidities, probably due to the possible influence of the target and beam spectator charges

→ The kinetic freeze-out radius decreases from the mid-rapidity region towards high rapidities, indicating that at forward and backward rapidities the system volume is smaller

# Energy dependence



→ data from HADES and RHIC-BES Au+Au collisions

→ Coulomb momentum decreases with the increase of beam energy, showing that the Coulomb interaction is stronger at lower energies.

If the colliding nuclei are fully stopped, due to a larger stopped charge in the overlap volume a stronger Coulomb field is generated, and this Coulomb field has a stronger influence on the charged pions.

# Conclusions

---

- Tsallis fits suggest that the dense baryonic matter created at HADES and AGS energies is well described by a nearly thermal source with only moderate non-equilibrium effects → the weak rapidity dependence indicates that the degree of non-equilibrium is relatively uniform throughout the system at freeze-out
- Tsallis temperature increases from HADES to AGS energies due to strong baryon stopping → at SPS and RHIC-BES energies,  $T$  saturates, reflecting the transition to a transparency-dominated regime where additional beam energy mainly drives expansion and particle production rather than further heating
- The parameter  $q$  is close to unity across all energies, indicating a system close to thermal equilibrium. Its weak energy dependence and saturation at higher energies suggest that the degree of non-equilibrium changes only slightly despite the transition from stopping to transparency
- The Coulomb effect decreases with the energy → consistent with the transition from strong baryon stopping to increasing nuclear transparency, resulting in a progressively weaker Coulomb influence on the charged pion spectra