

Accessing the strong interaction between Λ baryons and strange hadrons with femtoscopy at LHC

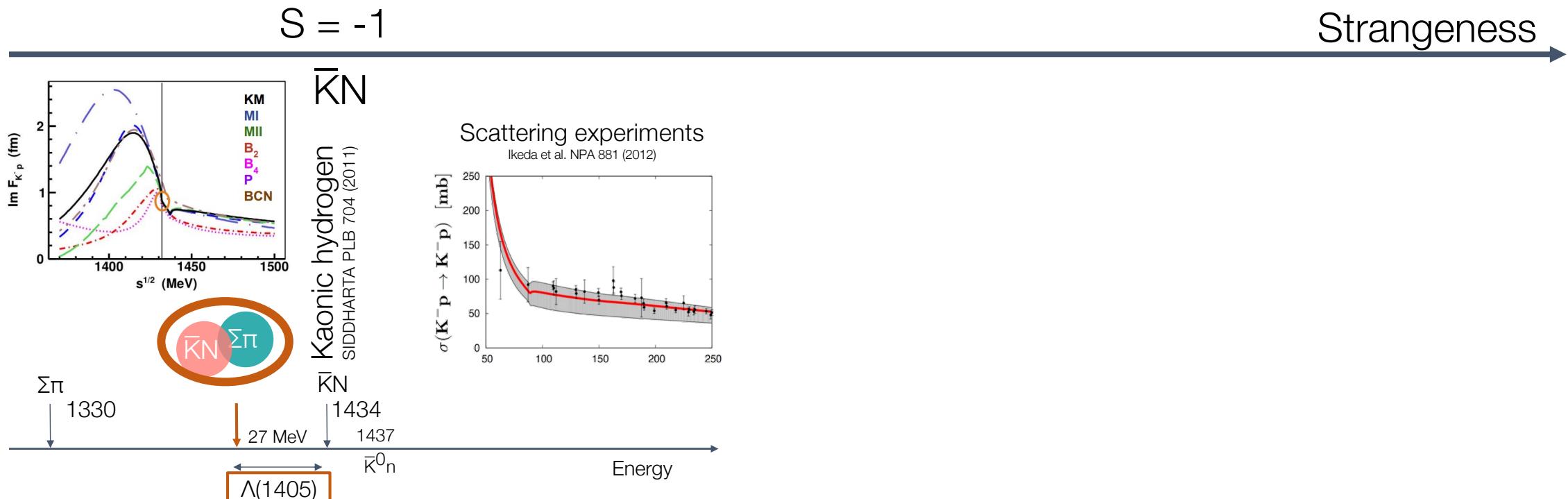
V. Mantovani Sarti (TUM) on behalf of the ALICE Collaboration

MESON, Kraków 22-27 June 2023

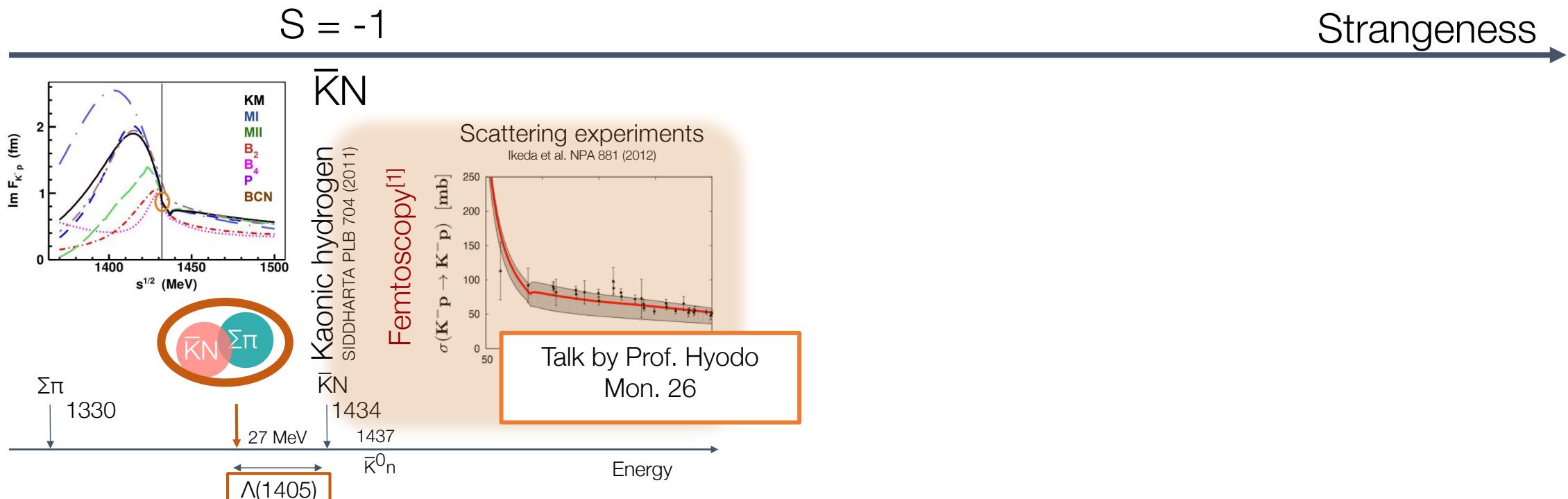


Based on:
ALICE Coll. arXiv: 2305.19093
valentina.mantovani.sarti@cern.ch

- Playground for many facets of low-energy QCD → Inelastic channels, bound and resonant states
- Still rather poorly constrained experimentally



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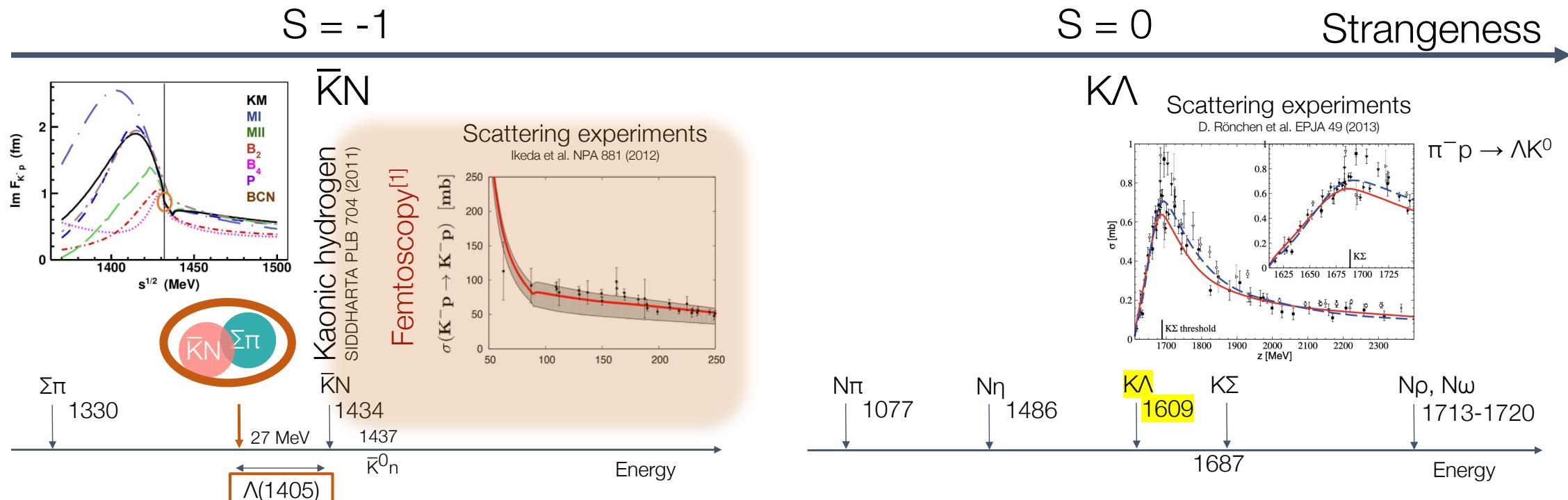


$K^- p$ correlations in different colliding systems (pp, p–Pb, Pb–Pb)
 → Improve understanding of $\Lambda(1405)$ molecular state^[2]

[1] pp: ALICE Coll. PRL 124 (2020)
 Pb–Pb: ALICE Coll. PLB 822 (2021)
 pp, p–Pb, Pb–Pb: ALICE Coll. EPJC 83 (2023)
 [2] M. Mai EPJ ST 230 (2021), 6, 1593–1607

Meson-baryon interactions with strangeness

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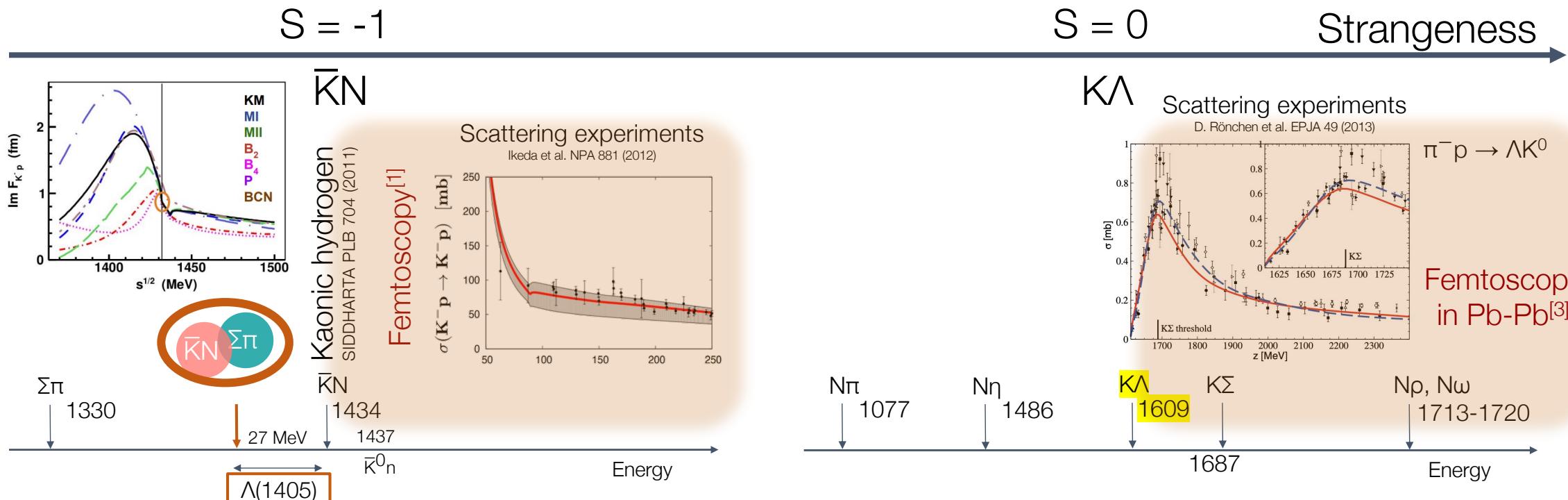
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Coupled-channel dynamics important for
the rich N^* and Δ resonance spectrum
→ Mainly driven by large amount of $N\pi$ data

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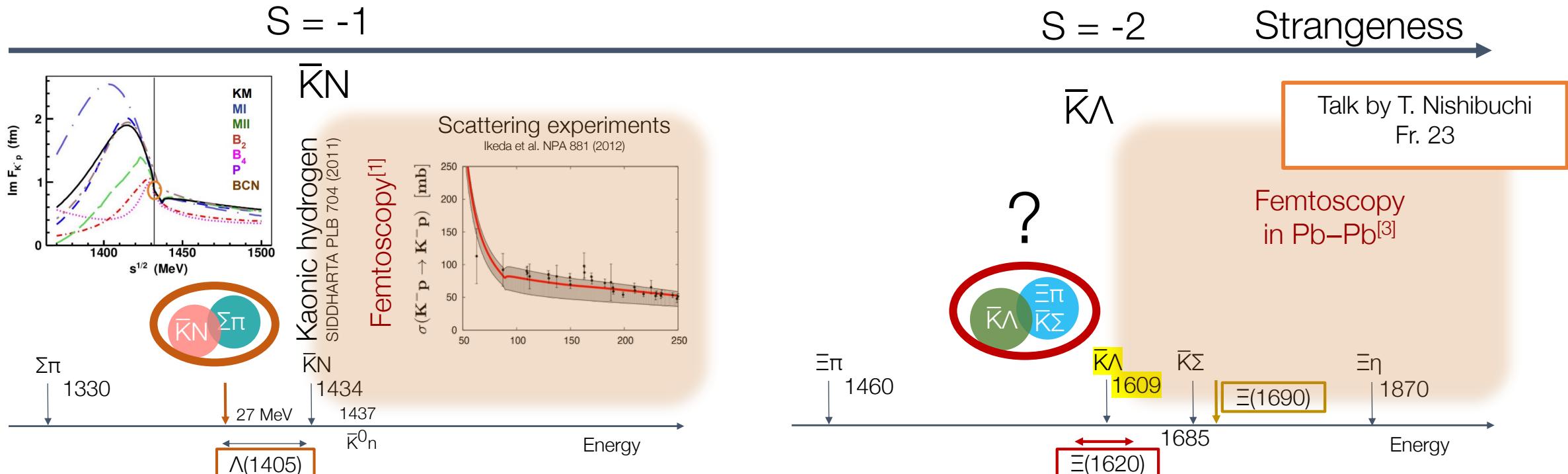
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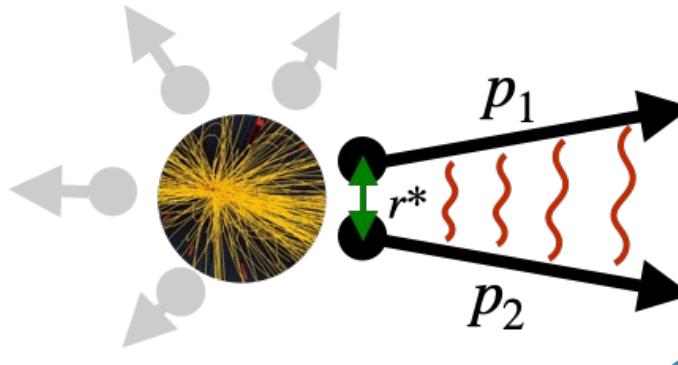
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Similar scenario for $\Xi(1620)$ state?
→ Poorly known, only observed in $\pi\Xi$ decay^[4]
→ Shed light on the nature of this state

[3] ALICE Coll. PRC 103 (2021), CMS Coll. arXiv:2301.05290
[4] Belle Coll. PRL 122 (2019)

The femtoscopy technique at ALICE

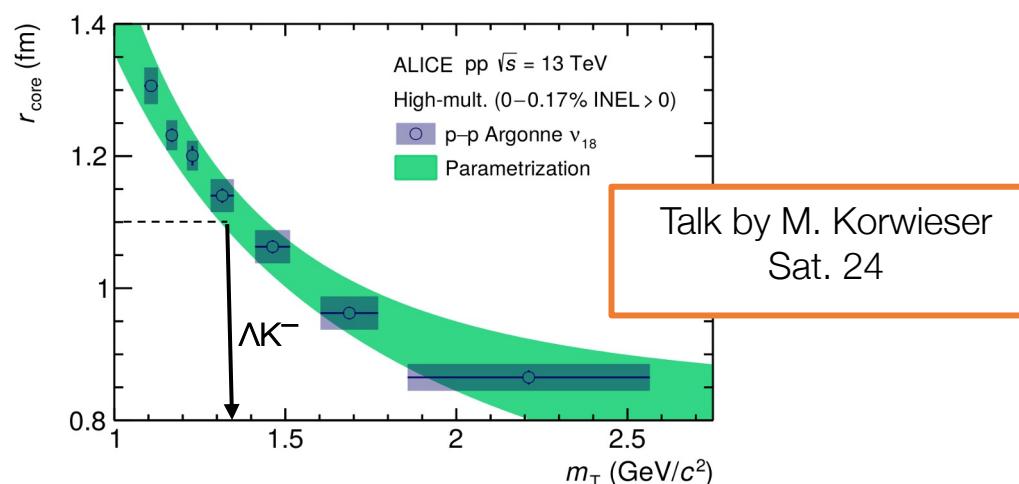


- Access to the short-range dynamics between hadrons^[1,2]:

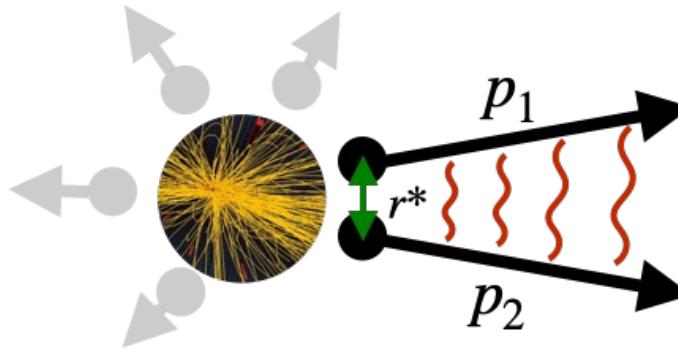
$$C(k^*) = \int S(\vec{r}^*) |\psi(\vec{k}^*, \vec{r}^*)|^2 d^3\vec{r}^* = \mathcal{N}(k^*) \frac{N_{\text{same}}(k^*)}{N_{\text{mixed}}(k^*)}$$

ALICE measurements shown today in high-multiplicity pp collisions at 13 TeV
 → Emitting source anchored to p-p correlation data^[3]
 → Interparticle distances ~1-2 fm in pp collisions

- Long-lived strongly decaying resonances feeding to Λ and K ($\langle c\tau \rangle_{\Lambda, K} \sim 4-5$ fm)
 → Effective double Gaussian



- [1] M.Lisa, S. Pratt et al, Ann.Rev.Nucl.Part.Sci. 55 (2005), 357-402
 [2] L. Fabbietti, VMS and O. Vazquez Doce ARNPS 71 (2021), 377-402
 [3] ALICE Coll., PLB 811 (2020), 135849

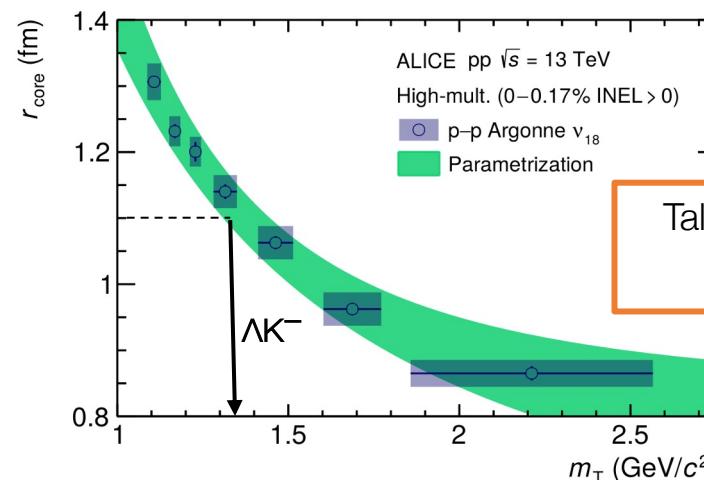


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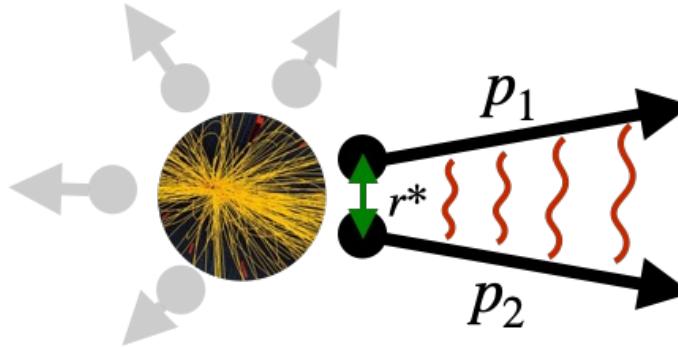
Talk by M. Korwieser
Sat. 24

Two-particle wave function^[4]
→ Profile of $C(k^*)$ vs nature of the interaction

Beyond 2-body
Talk by Prof. Fabbietti
Mon. 26

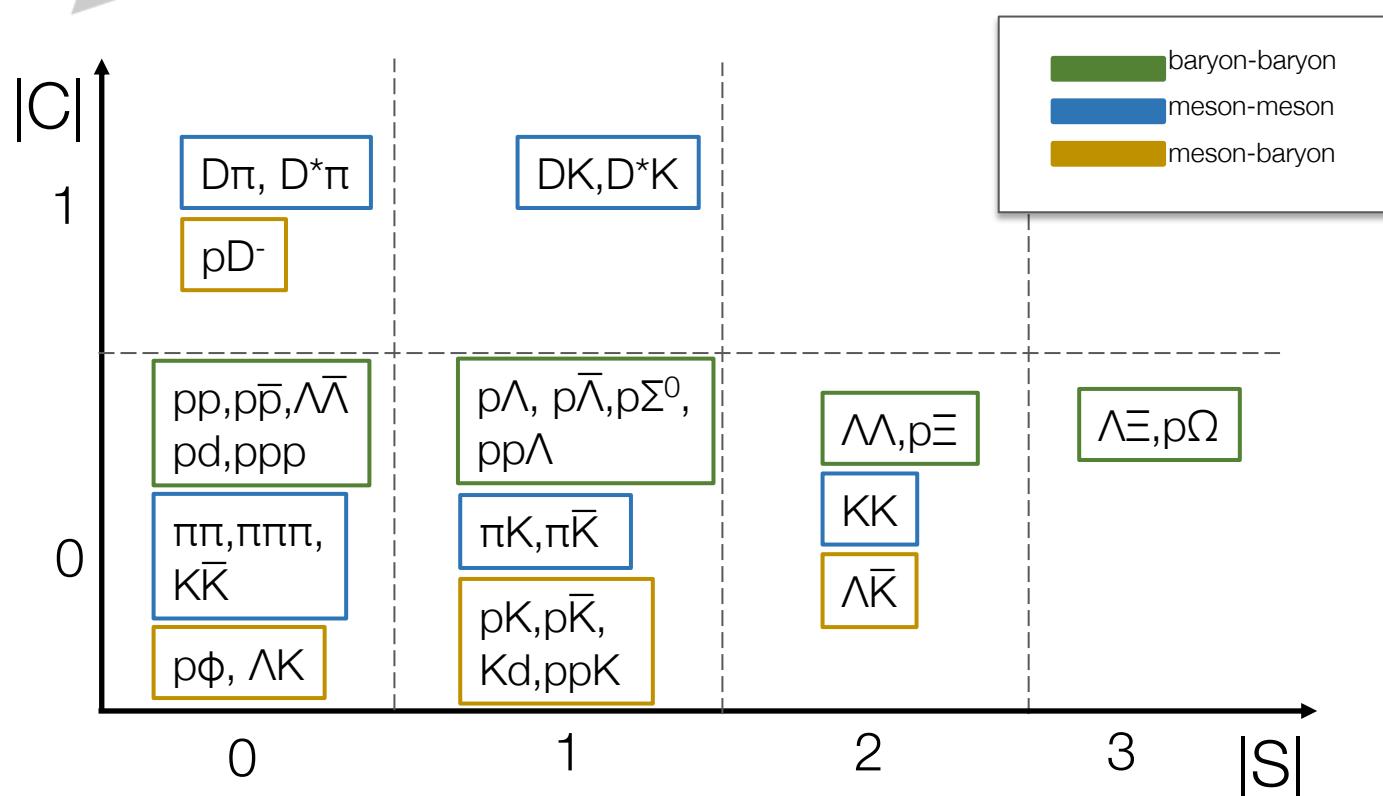
$$C(k^*) \begin{cases} > 1 & \text{Attractive} \\ < 1 & \text{Repulsive} \\ \geq 1 & \text{Bound state} \end{cases}$$

- [1] M.Lisa, S. Pratt et al, Ann.Rev.Nucl.Part.Sci. 55 (2005), 357-402
- [2] L. Fabbietti, VMS and O. Vazquez Doce ARNPS 71 (2021), 377-402
- [3] ALICE Coll., PLB 811 (2020), 135849
- [4] D. Mihaylov et al., EPJC 78 (2018), 5, 394



- Access to the **short-range dynamics** between hadrons^[1,2]:

$$C(k^*) = \int S(\vec{r}^*) |\psi(\vec{k}^*, \vec{r}^*)|^2 d^3\vec{r}^* = \mathcal{N}(k^*) \frac{N_{\text{same}}(k^*)}{N_{\text{mixed}}(k^*)}$$



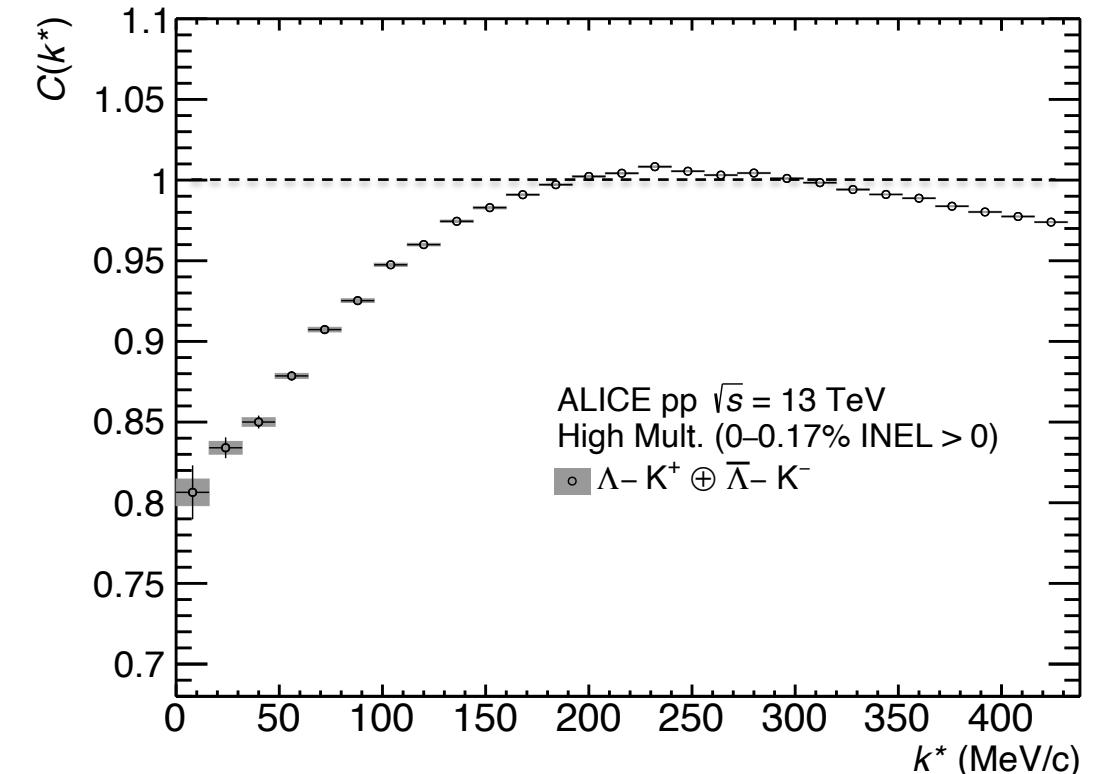
ALICE Collaboration:

- PRC 99 (2019) 2, 024001
 PLB 797 (2019) 134822
 PRL 123 (2019) 112002
 PRL 124 (2020) 09230
 PLB 805 (2020) 135419
 PLB 811 (2020) 135849
 Nature 588 (2020) 232-238
 PRL 127 (2021), 172301
 PLB 822 (2021), 136708
 PRC 103 (2021) 5, 055201
 PLB 833 (2022), 137272
 PLB 829 (2022), 137060
 PRD 106 (2022), 5, 05201
 arXiv:2204.10258 [nucl-ex], accepted in PLB
 arXiv:2206.03344 [nucl-ex], accepted in EPJA
 arXiv:2305.19093 [nucl-ex], submitted to PLB
 ... and more to come.

The ΛK^+ correlation in pp collisions

ALICE Coll. arXiv: 2305.19093

- Correlation below unity at low k^* → Repulsive interaction



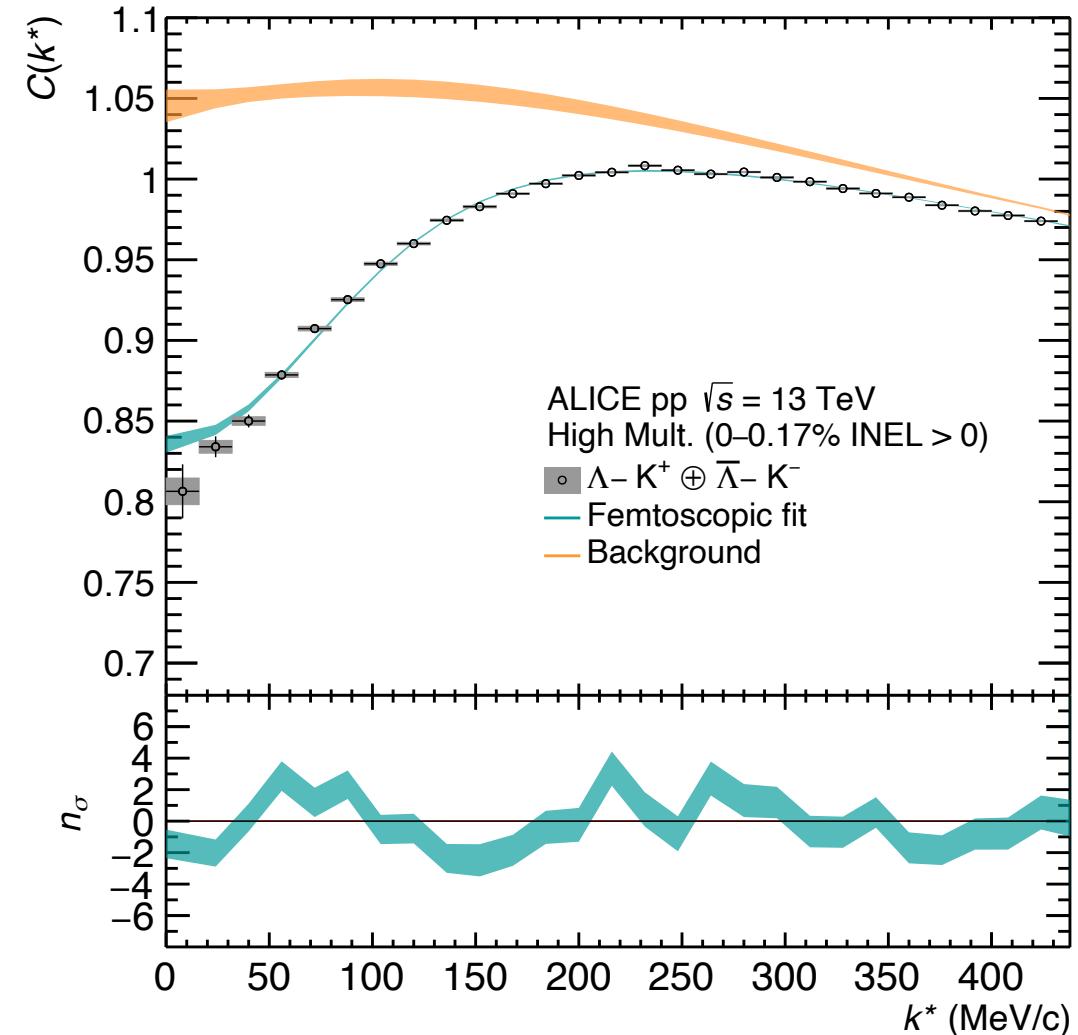
The ΛK^+ correlation in pp collisions

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- Correlation below unity at low k^* → **Repulsive interaction**

$$C(k^*) = N_D \cdot C_{model}(k^*) \cdot C_{bckg}(k^*)$$

- Residual background** due to initial parton scattering
 - Typically observed in meson-baryon and baryon-antibaryon correlations^[1,2]
 - Modeled using **Monte-carlo simulations**^[2]
- Genuine correlation** modeled with the Lednicky-Lyuboshits formula^[2,3]
 - Assuming the **scattering amplitude** within the **effective range expansion**



[1] ALICE Coll. Phys.Rev.Lett. 124 (2020)

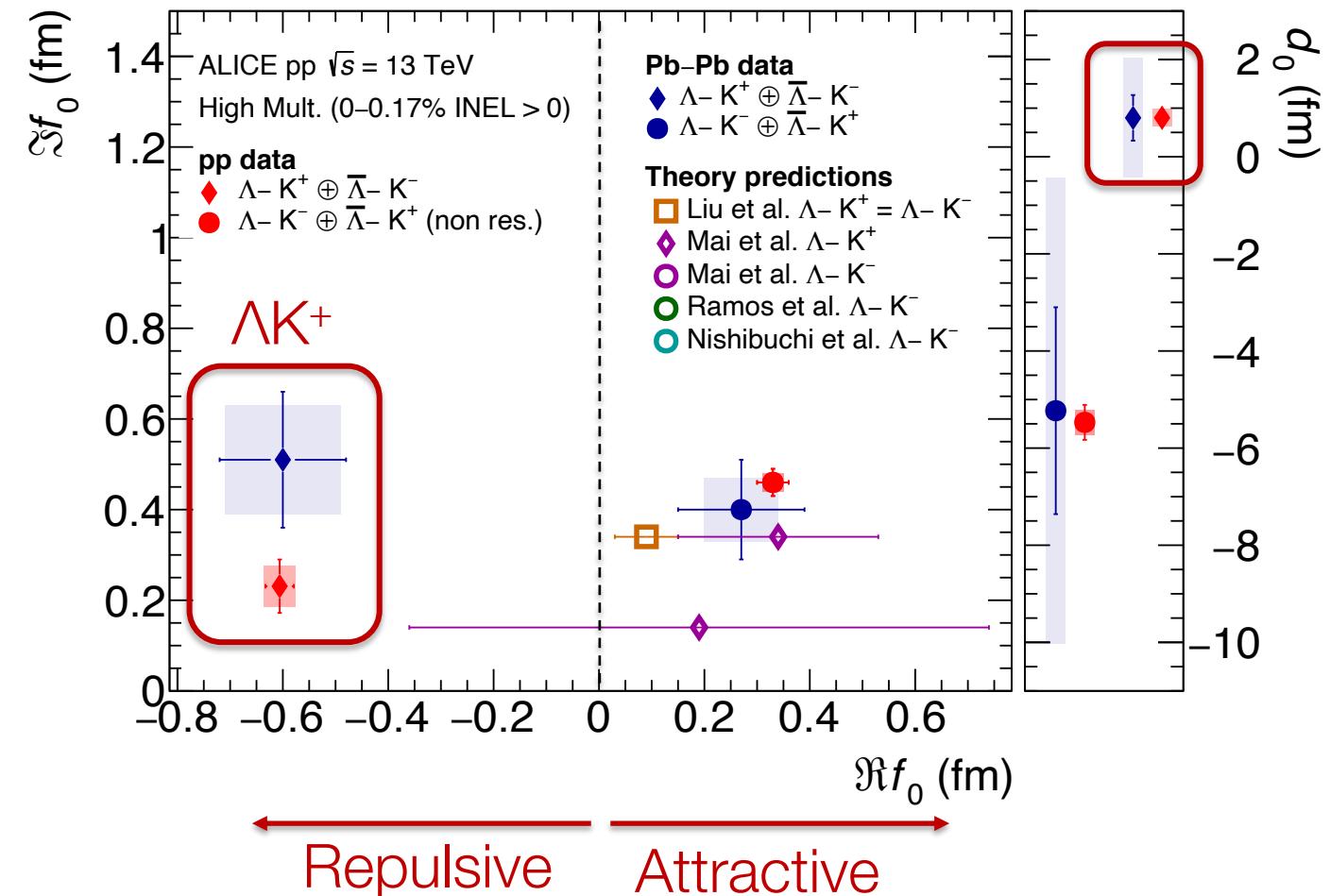
[2] ALICE Coll. PLB 829 (2022)

Scattering parameters for the ΛK^+ interaction

- Extracted scattering parameters
→ In agreement with ALICE Pb-Pb results^[1]
- Available models mainly indicating an attractive interaction
→ Parameters fixed based on SU(3) flavour symmetry, isospin symmetry
→ Mainly anchored to πN or $\bar{K}N$ data
- High-precision data to be used to constrain effective chiral theories

xPT at NLO: Liu et al. PRD 75 (2007), Mai et al. PRD 80 (2009)

ALICE Coll. arXiv: 2305.19093

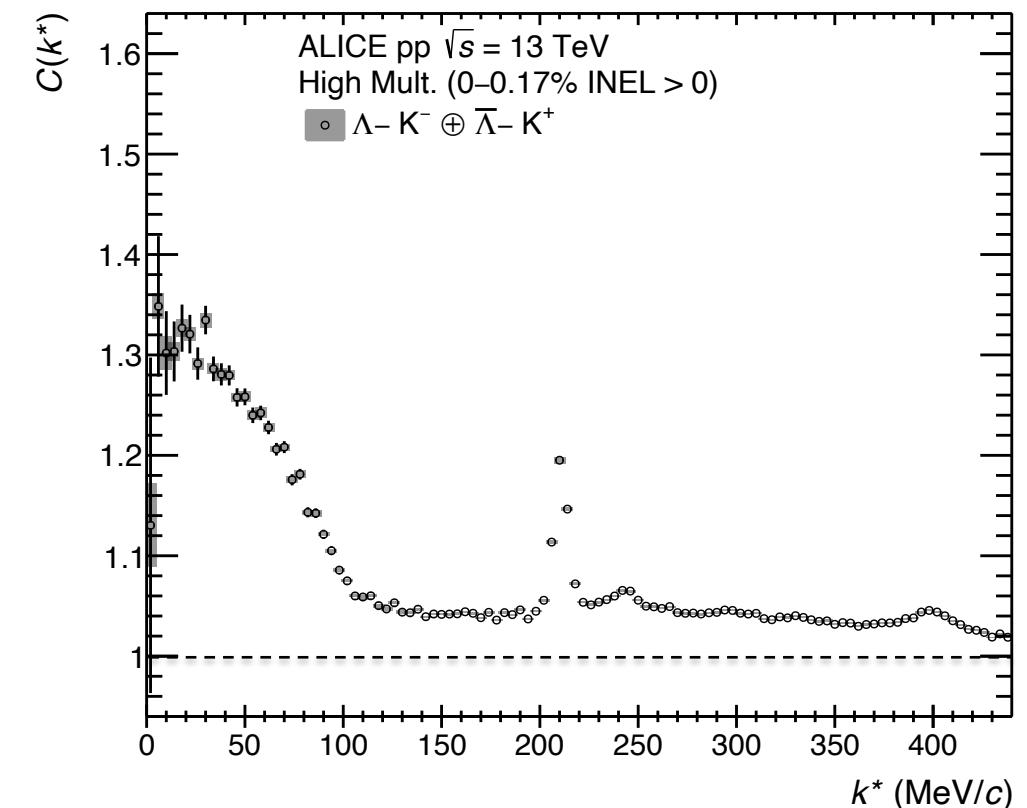


The ΛK^- correlation in pp collisions

ALICE Coll. arXiv: 2305.19093

- Correlation overall above 1 → **Attractive interaction**

$$C(k^*) = N_D \cdot C_{model}(k^*) \cdot C_{bckg}(k^*)$$

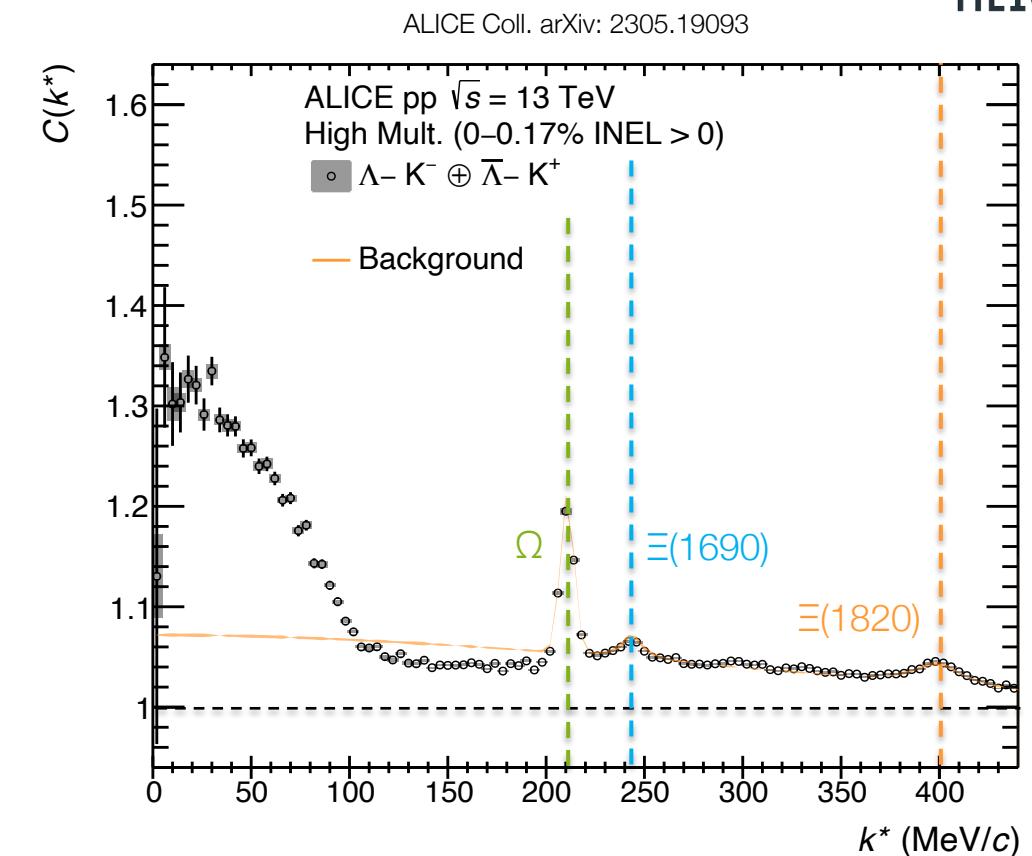


The ΛK^- correlation in pp collisions

- Correlation overall above 1 → **Attractive interaction**

$$C(k^*) = N_D \cdot C_{model}(k^*) \cdot C_{bckg}(k^*)$$

- Several peak structures in the measured correlation at $k^* > 200$ MeV/c
→ Ω , $\Xi(1690)$ and $\Xi(1820)$ included in the background
- Residual background due to initial parton scattering
→ Modeled using Monte-carlo simulations

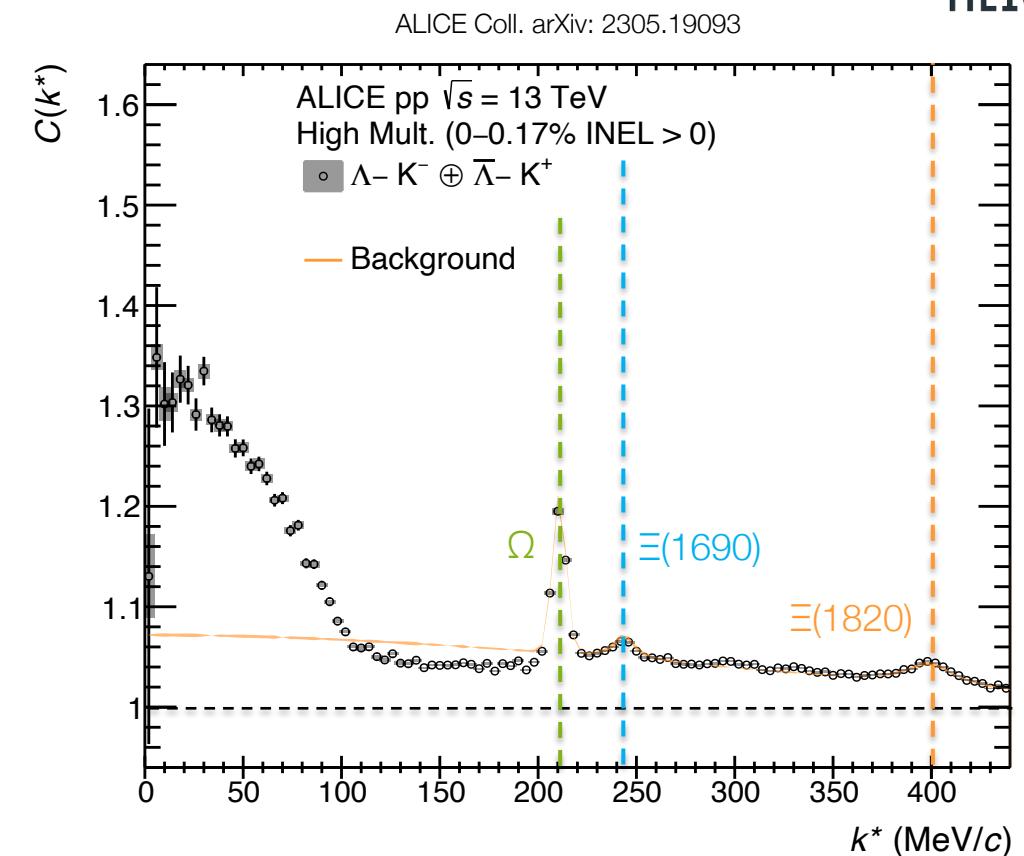


The ΛK^- correlation in pp collisions

- Correlation overall above 1 → **Attractive interaction**

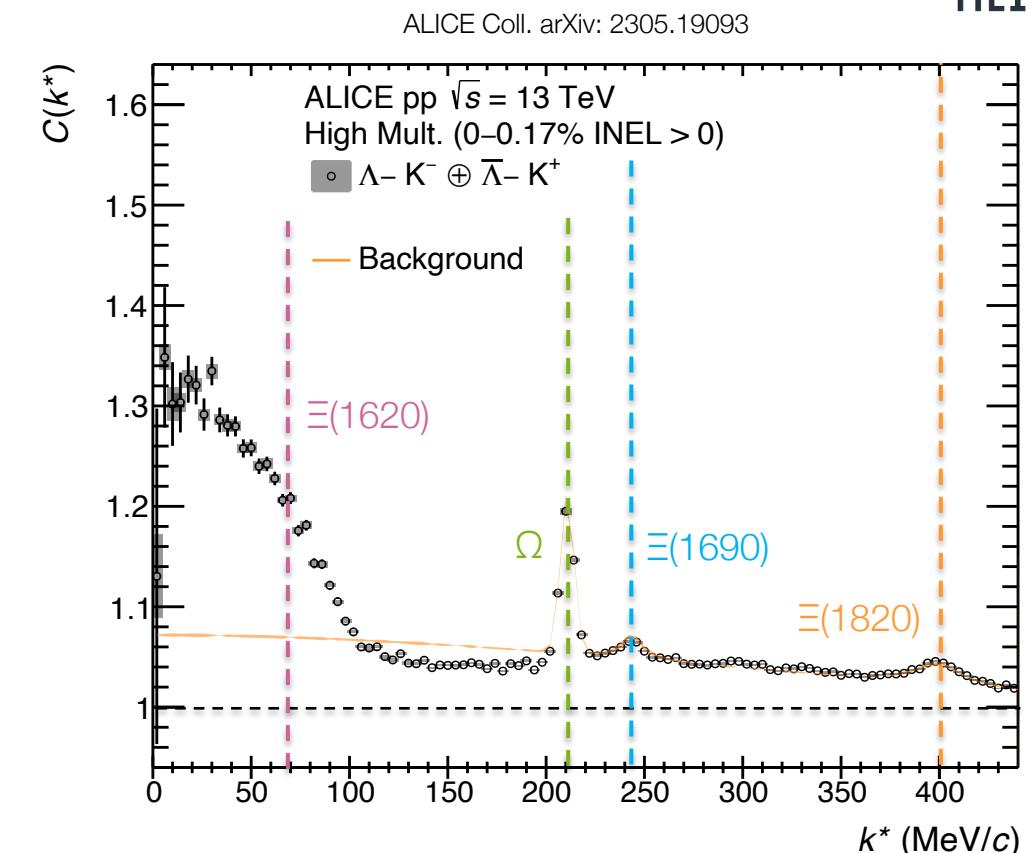
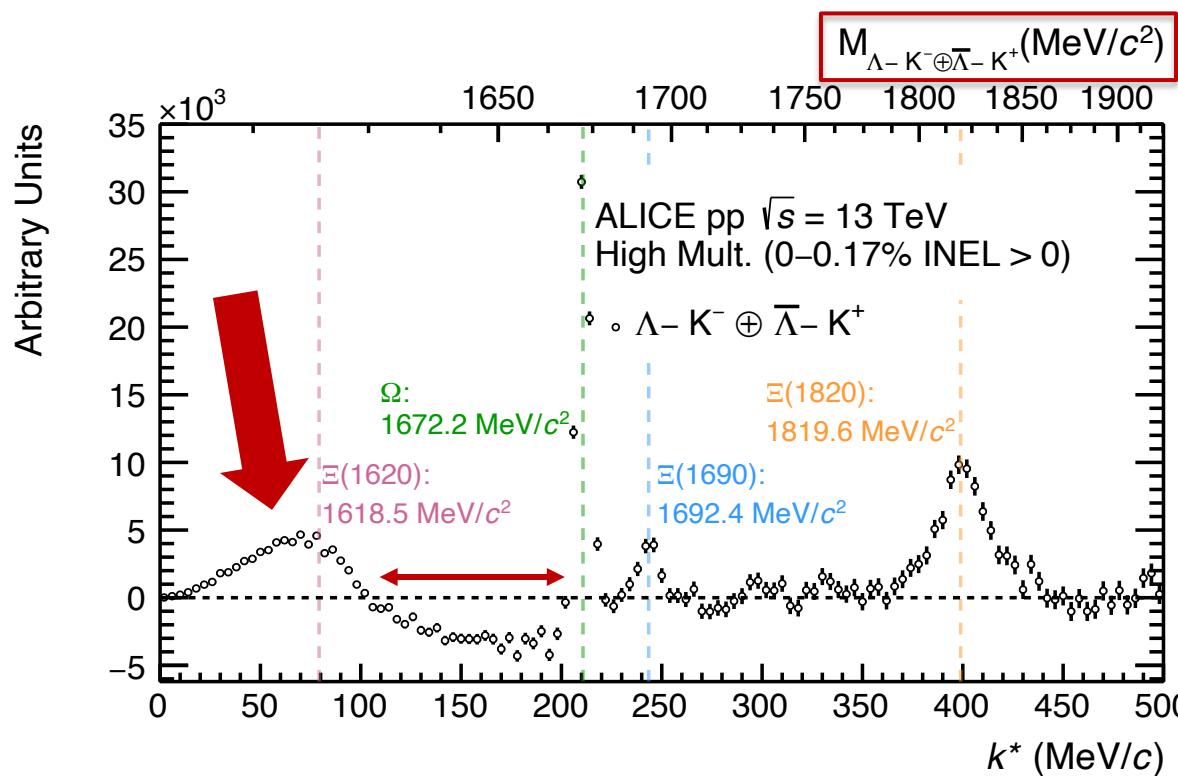
$$C(k^*) = N_D \cdot C_{model}(k^*) \cdot C_{bckg}(k^*)$$

- To understand how to model the genuine part we need to understand the signal at low k^* !



The ΛK^- correlation in pp collisions

- Correlation overall above 1 → **Attractive interaction**
- Invariant mass from same and mixed event distributions used to build the correlation
 - $\Xi(1620)$ just above the threshold
→ First experimental evidence of decay into ΛK^-

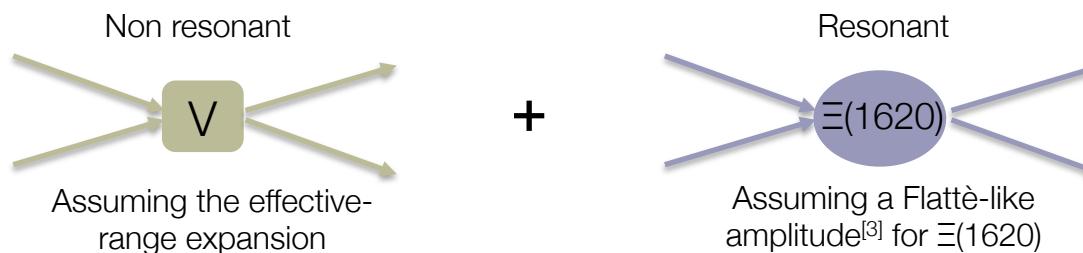


The ΛK^- correlation in pp collisions

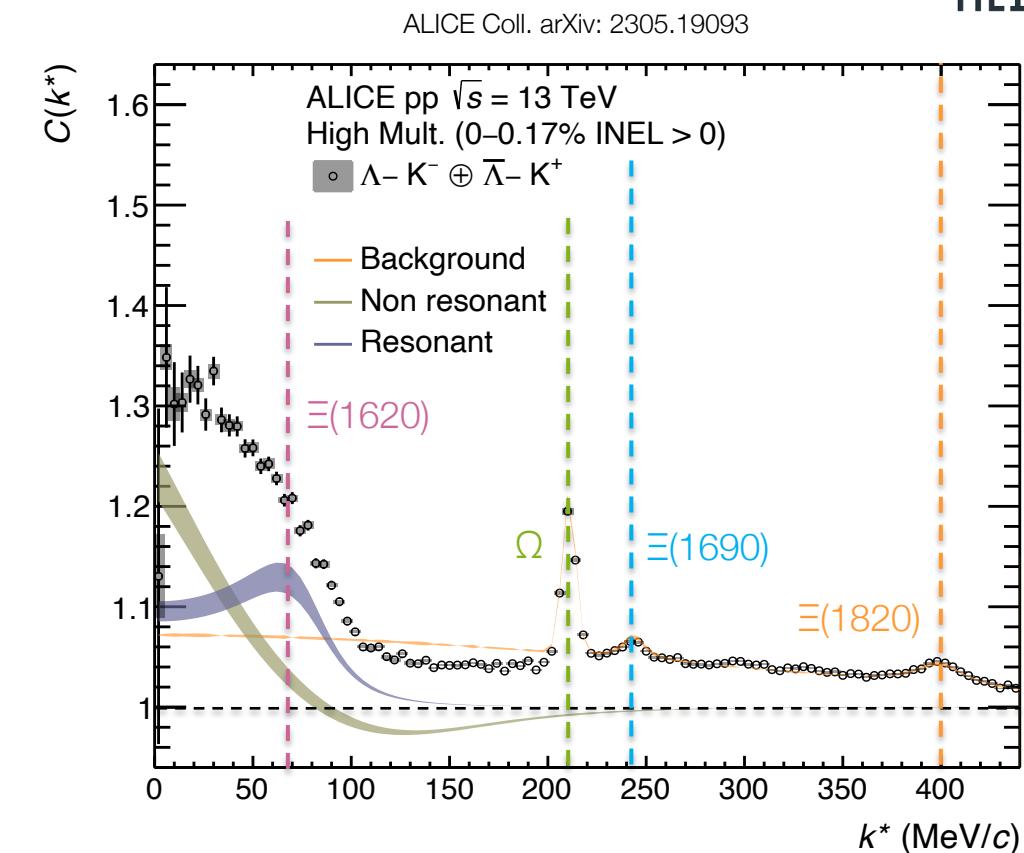
- Correlation overall above 1 → **Attractive interaction**

$$C(k^*) = N_D \cdot C_{model}(k^*) \cdot C_{bckg}(k^*)$$

- Genuine correlation of interest as weighted sum
→ Modeled with the Lednicky-Lyuboshits formula^[1,2]



Talk by Prof. Giacosa
Th. 22

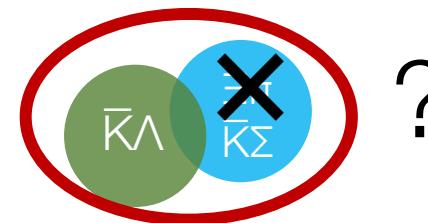


- [1] R. Lednicky, V. Lyuboshits SJNP 35 (1982)
[2] CATS: D. Mihaylov et al., EPJC 78 (2018), 5, 394
[3] F. Giacosa et al. EPJA 57 (2021), 12, 336

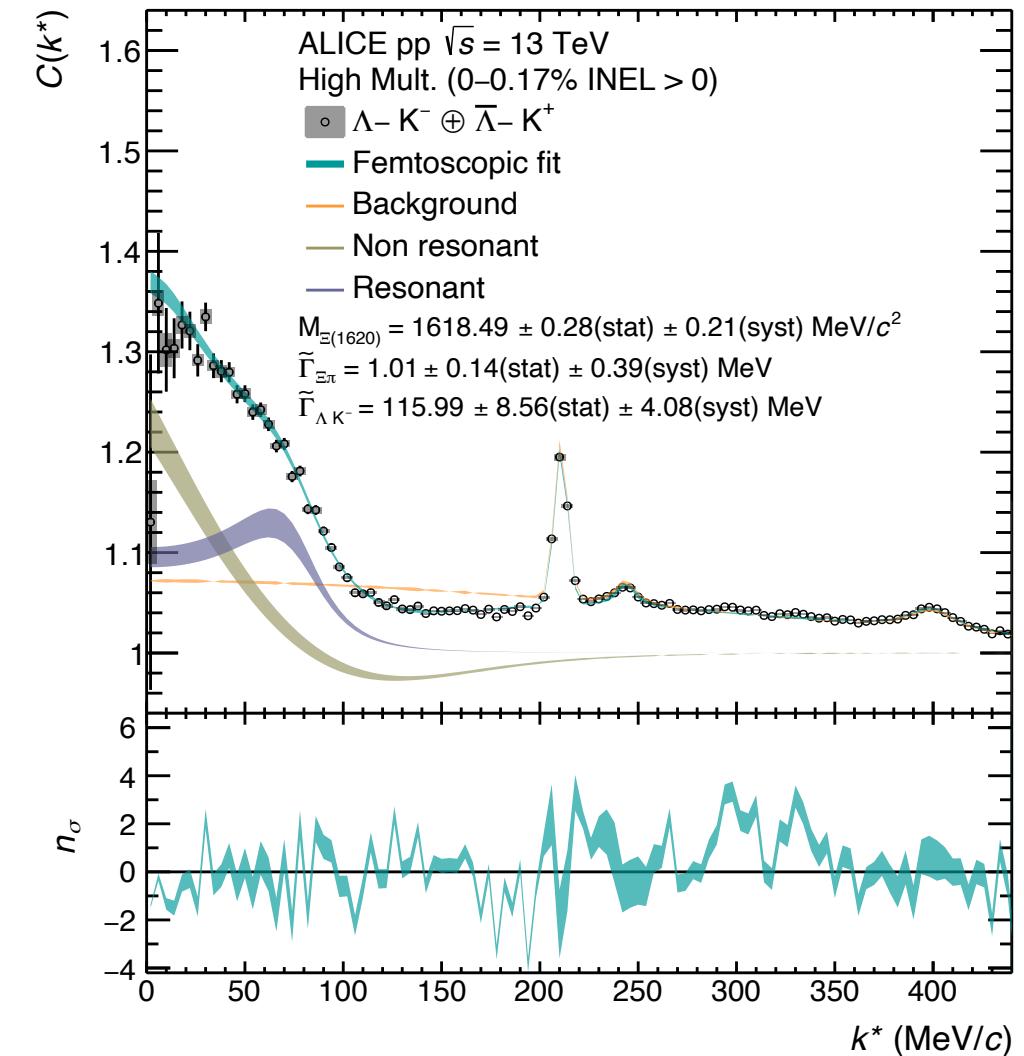
The $\Lambda\bar{K}$ - correlation in pp collisions

ALICE Coll. arXiv: 2305.19093

- Most precise data on $\Lambda\bar{K}$ interaction at low momenta
- LL model well reproduces the data in the whole k^* region
 - Interplay between resonant and non-resonant interaction
 - Non-resonant scattering parameters in agreement with Pb–Pb measurements
- $\Xi(1620)$ properties
 - Mass in agreement with Belle^[1]
 - $M_{\Xi(1620)} = 1618.49 \pm 0.28(\text{stat}) \pm 0.21(\text{syst}) \text{ MeV}/c^2$
 - Indication of a large coupling of $\Xi(1620)$ to $\Lambda\bar{K}^-$?



Can we use these femtoscopic data
to constrain low-energy QCD models?



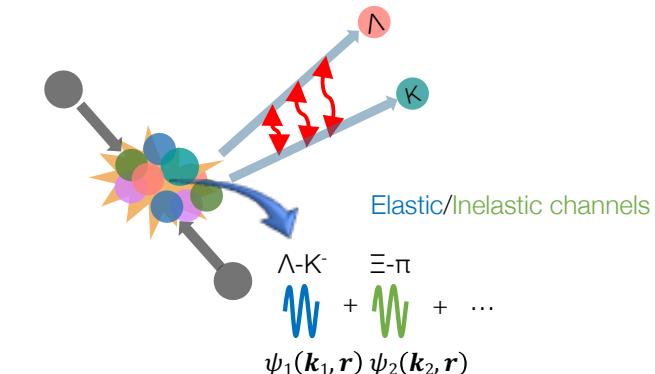
[1] Belle Coll. PRL 122 (2019), 7, 07250

- Model the genuine ΛK^- part including the coupled-channel dynamics

$$C(k^*) = \int S_1(\vec{r}^*) |\Psi_{1 \rightarrow 1}(\vec{k}^*, \vec{r}^*)|^2 d^3 r^* + \sum_{j \neq 1} \omega_j^{\text{prod}} \cdot \int S_j(\vec{r}^*) |\Psi_{j \rightarrow 1}(\vec{k}_j^*, \vec{r}^*)|^2 d^3 r^*$$

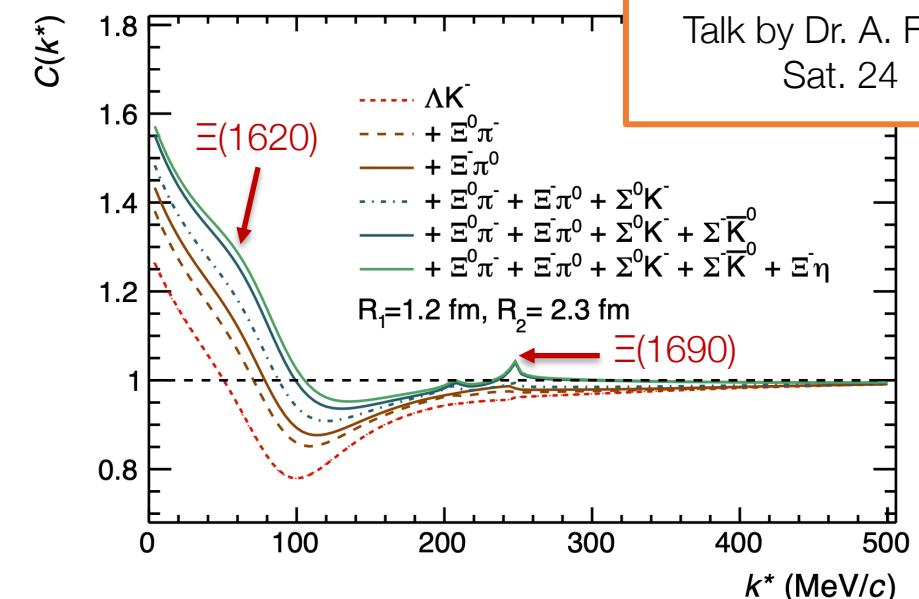
elastic $\Lambda K^- \rightarrow \Lambda K^-$

inelastic $\Xi\pi, \Sigma\bar{K}, \Xi\eta \rightarrow \Lambda K^-$



- Conversion weights ω_j^{prod} ^[1]
 - How many $\Xi\pi, \Sigma\bar{K}, \Xi\eta$ are produced as initial states?
- Wave functions $\Psi_{1,j \rightarrow 1}(\vec{k}_j^*, \vec{r}^*)$
 - State-of-the-art UxPT at NLO in CC formalism^[2]
 - Above threshold: modify the shape of CF
 - Below threshold: increase the strength of CF
 - $\Xi(1620)$ and $\Xi(1690)$ dynamically generated states

Femtoscopic data used to fix input parameters of the model!



[1] ALICE Coll. Eur.Phys.J.C 83 (2023) 4, L. Fabbietti, VMS, O. Vazquez Doce Ann.Rev.Nucl.Part.Sci. 71 (2021)

[2] A. Feijoo et al., PLB 841 (2023)

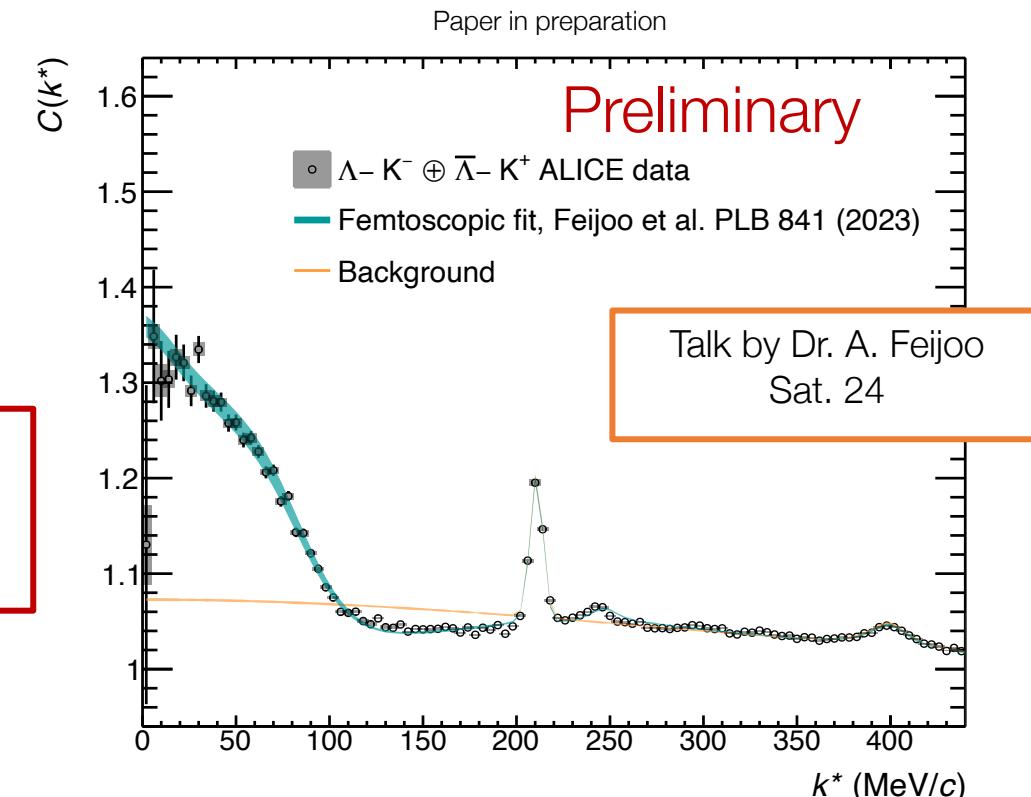
- Fit LECs and SCs to the measured ALICE $\Lambda\bar{K}$ correlation
- How does the $\Xi(1620)$ pole scenario look like?

Table 3: Poles, couplings and compositeness for model WT+NLO

$0^- \oplus \frac{1}{2}^+$ interaction in the $(I, S) = (0, -2)$ sector	
$M = 1615.46$ MeV	
$\Gamma = 20.92$ MeV	
$(-, -, +, +) \rightarrow$ pole	
$ g_i $	$ g_i^2 dG/dE $
$\pi\Xi(1456)$	0.631
$\bar{K}\Lambda(1611)$	0.919
$\bar{K}\Sigma(1689)$	2.15
$\eta\Xi(1866)$	2.75
$M = 1687.69$ MeV	
$\Gamma = 17.16$ MeV	
$(+, +, -, +) \rightarrow$ virtual	
$ g_i $	$ g_i^2 dG/dE $
$\pi\Xi(1456)$	0.581
$\bar{K}\Lambda(1611)$	0.576
$\bar{K}\Sigma(1689)$	1.54
$\eta\Xi(1866)$	0.727

$\Xi(1620)$ pole
Mainly molecular nature
composed of $\bar{K}\Sigma$
NEW PARADIGM!

$\Xi(1690)$ pole
Virtual state
Mainly coupled to $\bar{K}\Sigma$



Work in collaboration with:
Dr. A. Feijoo, Dr. I. Vidana, Prof. A. Ramos,
Prof. F. Giacosa,
Prof. T. Hyodo and Dr. Y. Kamiya

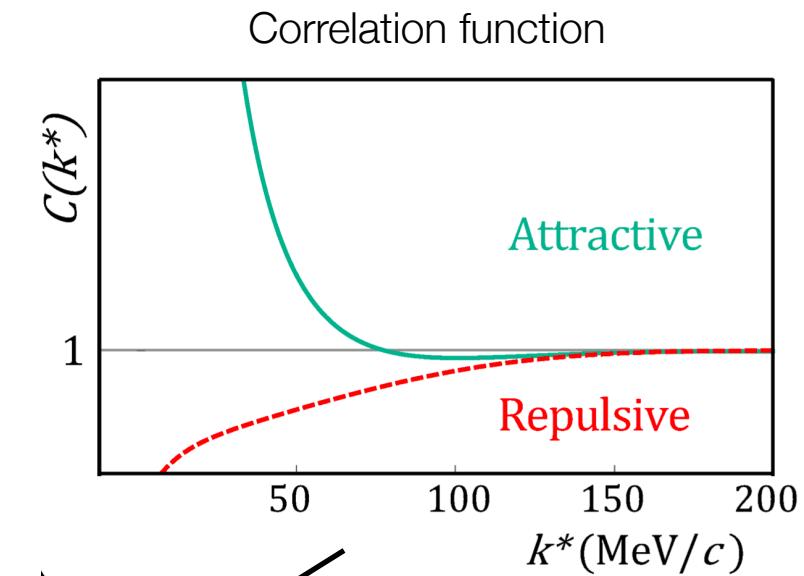
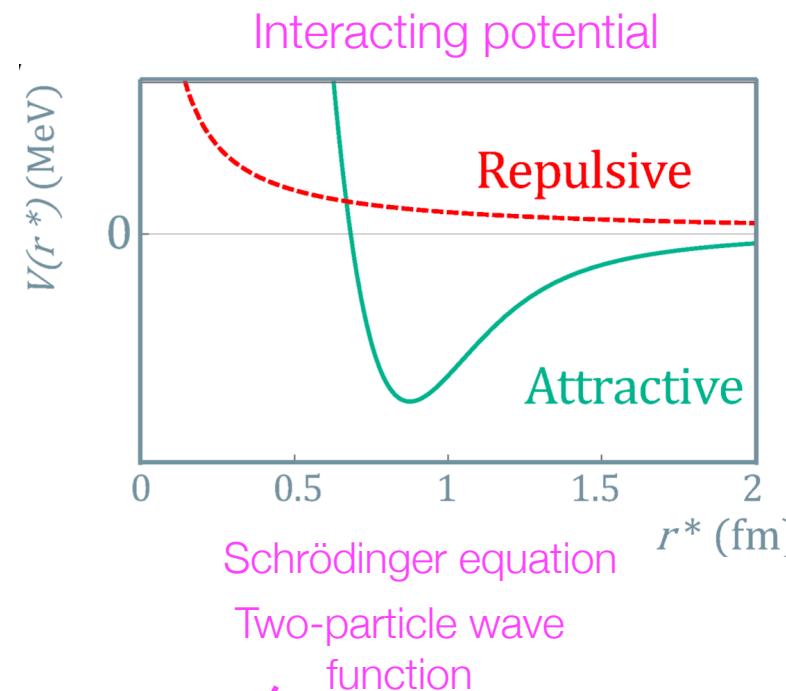
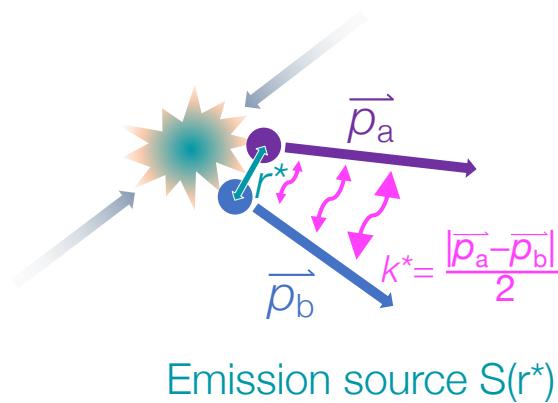
- Femtoscopy as a complementary tool to access the strange sector of meson-baryon interaction
→ **Most precise ΛK^+ and ΛK^- data at low momenta and first evidence of $\Xi(1620)$ in the ΛK^- decay channel**
- **First use of femtoscopic data to tune chiral effective potentials**
→ Possibility to constrain the QCD low energy sector for interactions not experimentally available before
- Extend the measurements in pp collisions to $\pi\Xi$
→ Provide additional constraints to models and check consistency with other available measurements

Femtoscopy@Meson 2023

- “*Towards many body nuclear interaction studies at the LHC*”, Prof. L. Fabbietti Mon. 26
- “*Towards a common particle-emitting source in small systems for mesons and baryons with ALICE*”, M. Korwieser Sat. 24
- “*Correlation Function constraints on S=-2 meson-baryon interaction from UChPT*”, Dr. A. Feijoo Sat. 24
- “*A study of K^-d and K^+d interactions via femtoscopy technique*”, W. Rzeszka Sat. 24
- “*Femtoscopy of $p - \Lambda$ system obtained in heavy-ion collision in the HADES experiment*”, Narendra Rathod Sat. 24

Additional slides

The femtoscopy technique at ALICE



$$C(k^*) = \int \underbrace{S(\vec{r}^*)}_{\text{Schrödinger equation}} \underbrace{|\psi(\vec{k}^*, \vec{r}^*)|^2}_{\text{Two-particle wave function}} d^3\vec{r}^* = \mathcal{N}(k^*) \frac{N_{\text{same}}(k^*)}{N_{\text{mixed}}(k^*)}$$

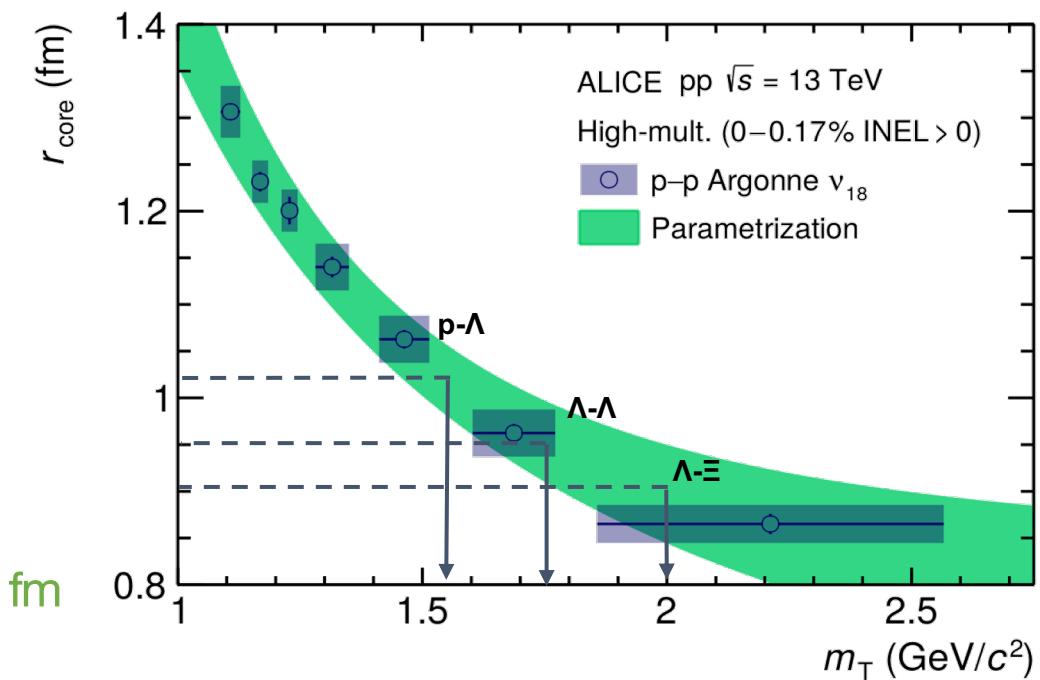
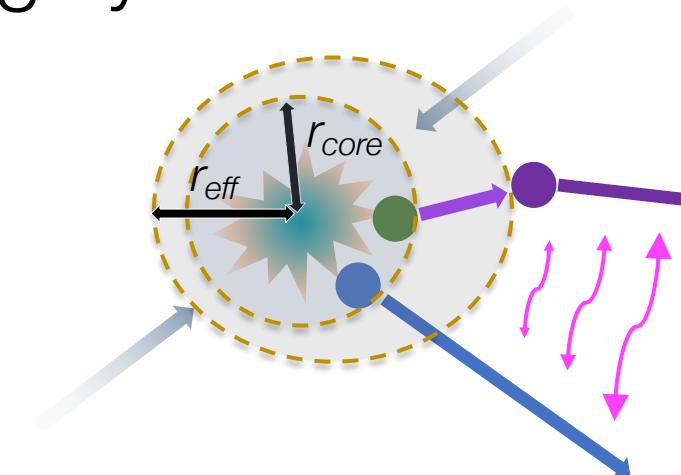
Measuring $C(k^*)$, fixing the source $S(r^*)$, study the interaction

The emitting source in small colliding systems

- Data-driven analysis on p-p and p- Λ pairs
 - Possible presence of collective effects $\rightarrow m_T$ scaling of the core radius
 - Contribution of **strongly decaying resonances with $c\tau \sim 1$ fm (*)**
- Common universal core source for baryons
- Core constrained from p-p pairs
 - Fixing of the source at corresponding $\langle m_T \rangle$
 \Rightarrow direct access to the interaction

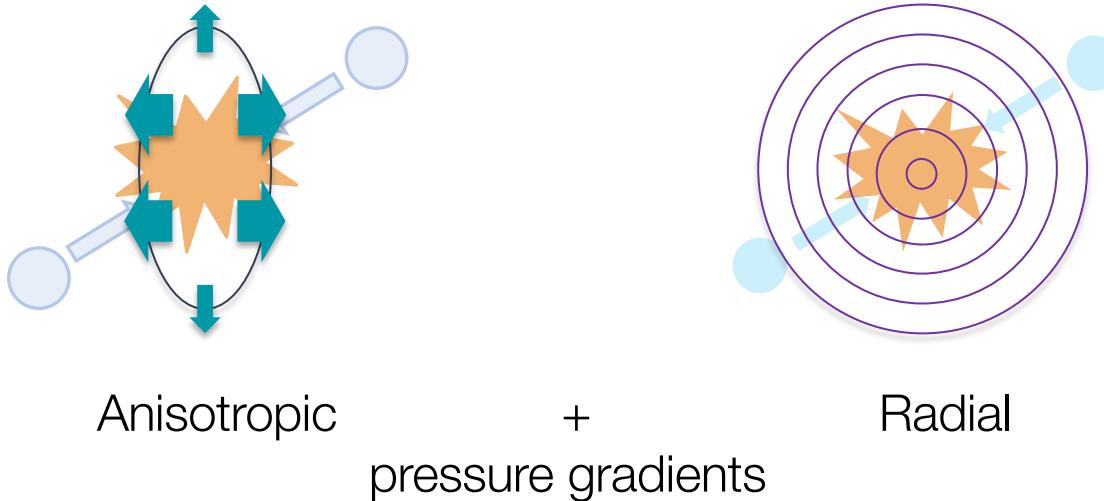
Particle	Res.	$\langle c\tau \rangle$ (fm)
p	Δ, N^*	1.6
Λ	Σ, Σ^*	4.7

$$r_{\text{eff}} = 1 - 1.25 \text{ fm}$$



Based on ALICE Coll. PLB 811 (2020) 135849

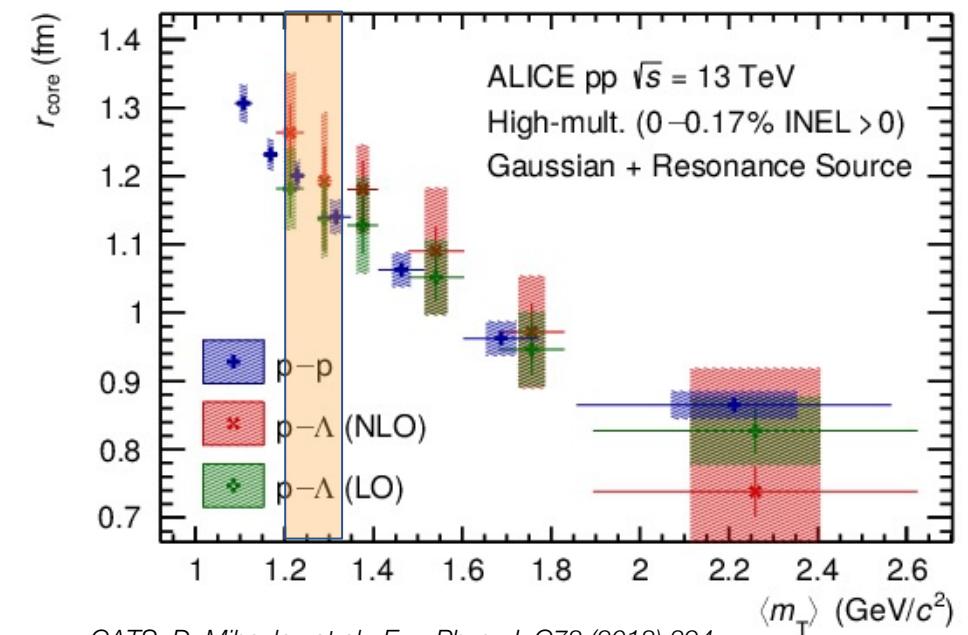
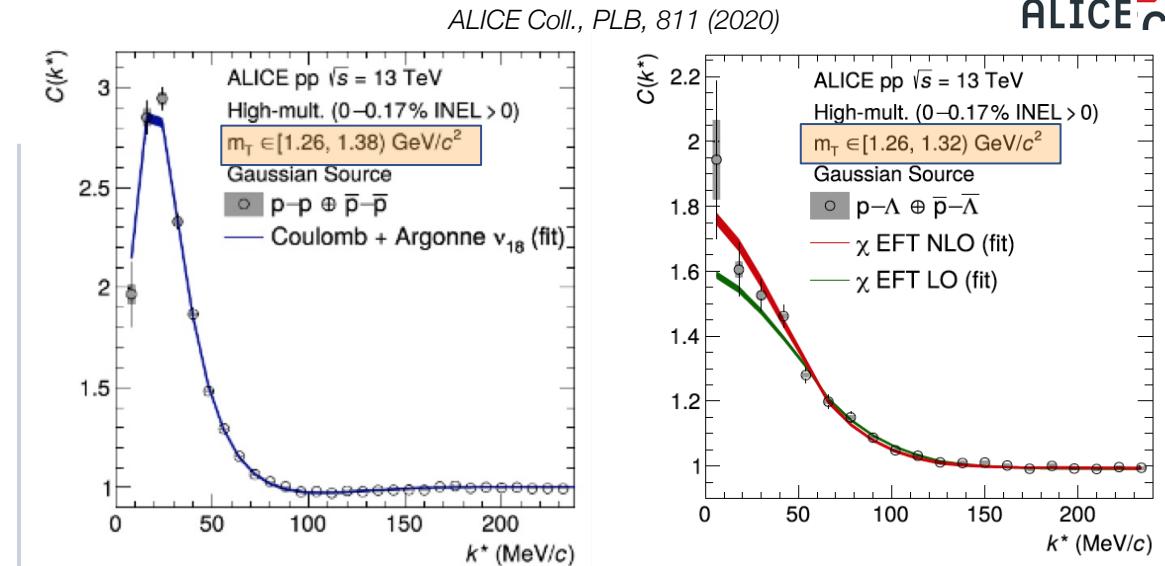
(*) U. A. Wiedemann, U. W. Heinz, Phys.Rept. 319, 145-230 (1999)



Different effect on different masses

$$C(k^*) = \int [S(r)] |\psi(\vec{k}^*, \vec{r})|^2 d^3r$$

$$S(r) = G(r, r_{core}(m_T)) = \frac{1}{(4\pi r_{core}^2)^{3/2}} \exp\left(-\frac{r^2}{4r_{core}^2}\right) \otimes \frac{1}{s} \exp\left(-\frac{r}{s}\right)$$



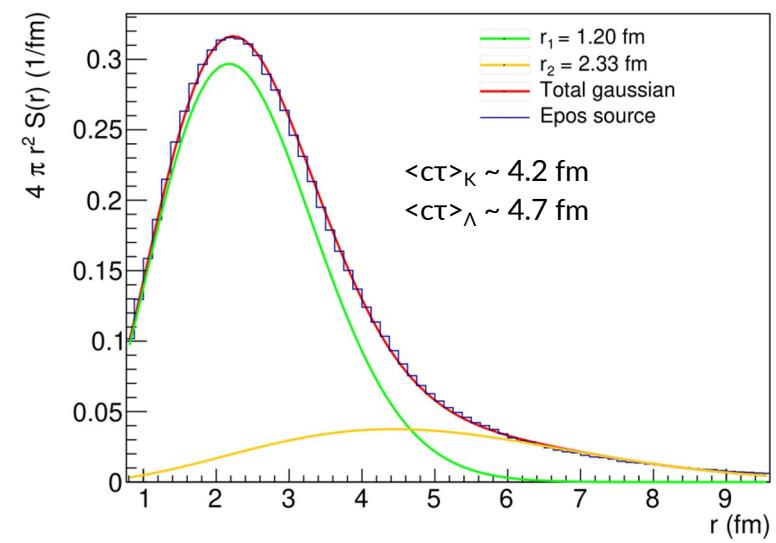
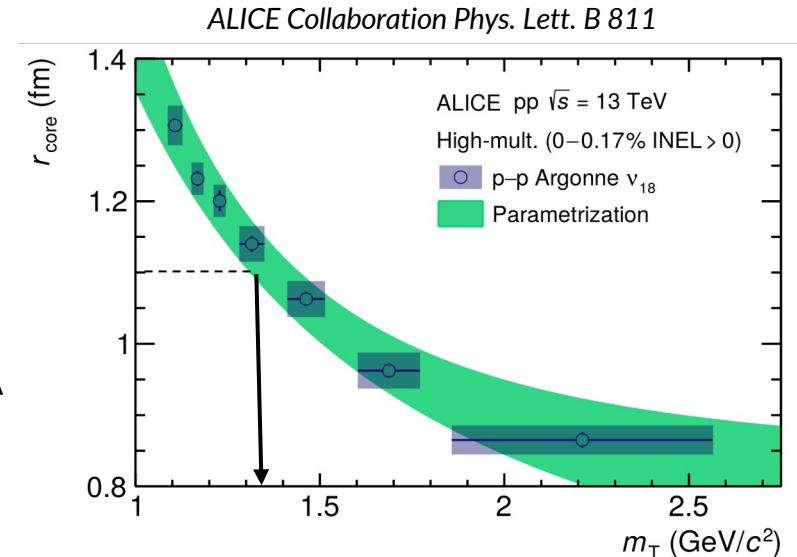
CATS: D. Mihaylov et al., Eur. Phys. J. C78 (2018) 394

Fixing the source

- Core-halo resonance model anchored to p-p CF
 - $r_{\text{core}} = 1.11 \pm 0.04$ ($\langle m_T \rangle_{\Lambda K} = 1.35 \text{ GeV}/c^2$)
- Long-lived strongly decaying resonances feeding to Λ
 - fit with effective double gaussian

$$S_{\text{tot}}(r) = \lambda_s [\omega_S \cdot S(r_1) + (1-\omega_S) \cdot S(r_2)]$$

Parameter	Value
r_{core} [fm]	$1.11^{+0.04}_{-0.04}$
$r_{1,\text{eff}}$ [fm]	$1.202^{+0.043}_{-0.042}$
$r_{2,\text{eff}}$ [fm]	$2.330^{+0.050}_{-0.045}$
ω	$0.7993^{+0.0037}_{-0.0027}$
λ	$0.9806^{+0.0006}_{-0.0008}$



Overview on S=-2 meson-baryon sector

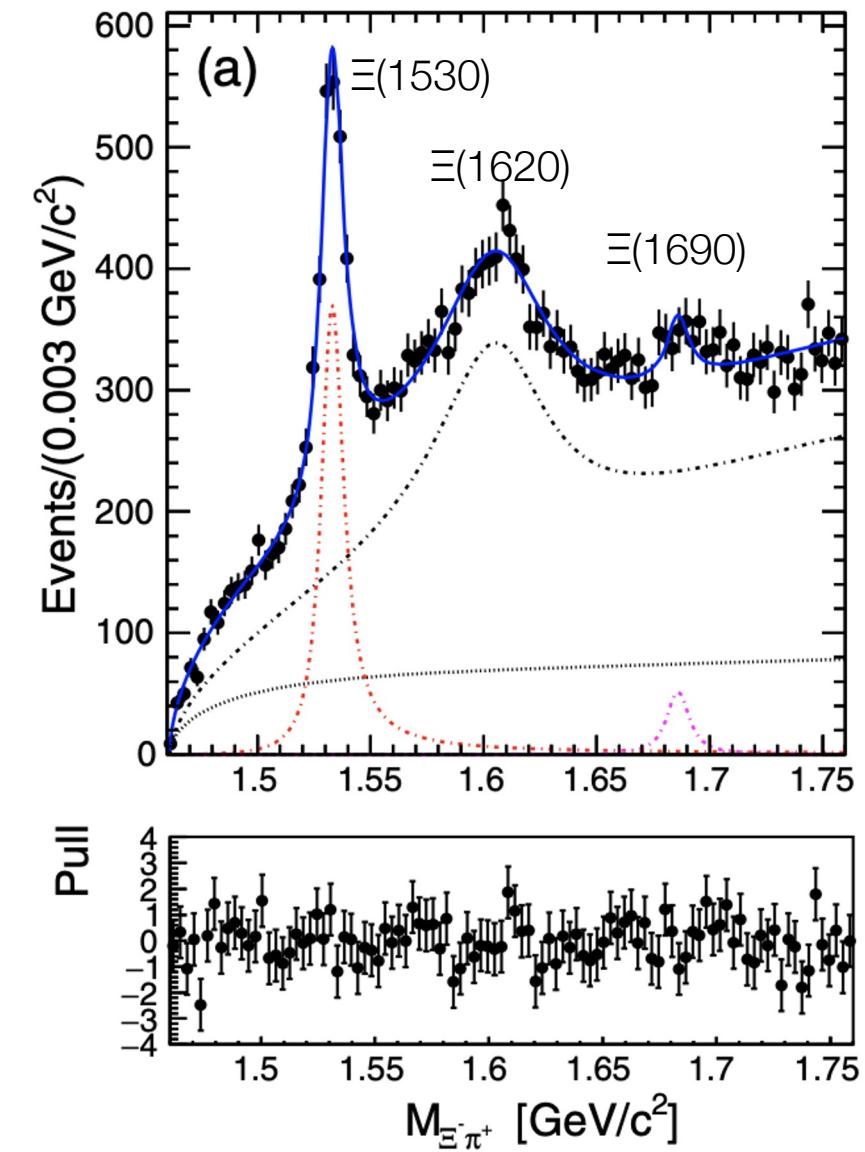
- Poorly constrained experimentally
→ Effective lagrangians anchored to S=-1 sector^[1]
- $\Xi(1620)$ and $\Xi(1690)$ ^[2] dynamically generated states within coupled-channel approaches
→ $\Xi(1620)$ observed by Belle in $\pi\bar{\Xi}$ decay but currently only 1 star in PDG
- Recent development of chiral calculations at NLO^[3]

Need for high-precision data to constrain models' parameters

- [1] A. Ramos et al. PRL 89 (2002)
[2] LHCb Coll. Sci.Bull. 66 (2021)
[3] A. Feijoo et al. PLB 841 (2023)

- LO xPT calculations:
C. Garcia-Recio et al. PLB 582 (2004)
D. Gamermann et al. PRD 84 (2011)
T. Sekihara PTEP 2015 (9) (2015)
T. Nishibuchi and T. Hyodo, EPJ Web Conf 271 (2022)

Belle Coll. PRL 122 (2019)

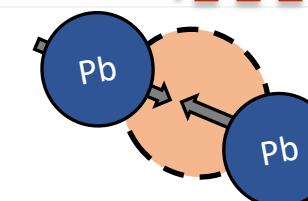
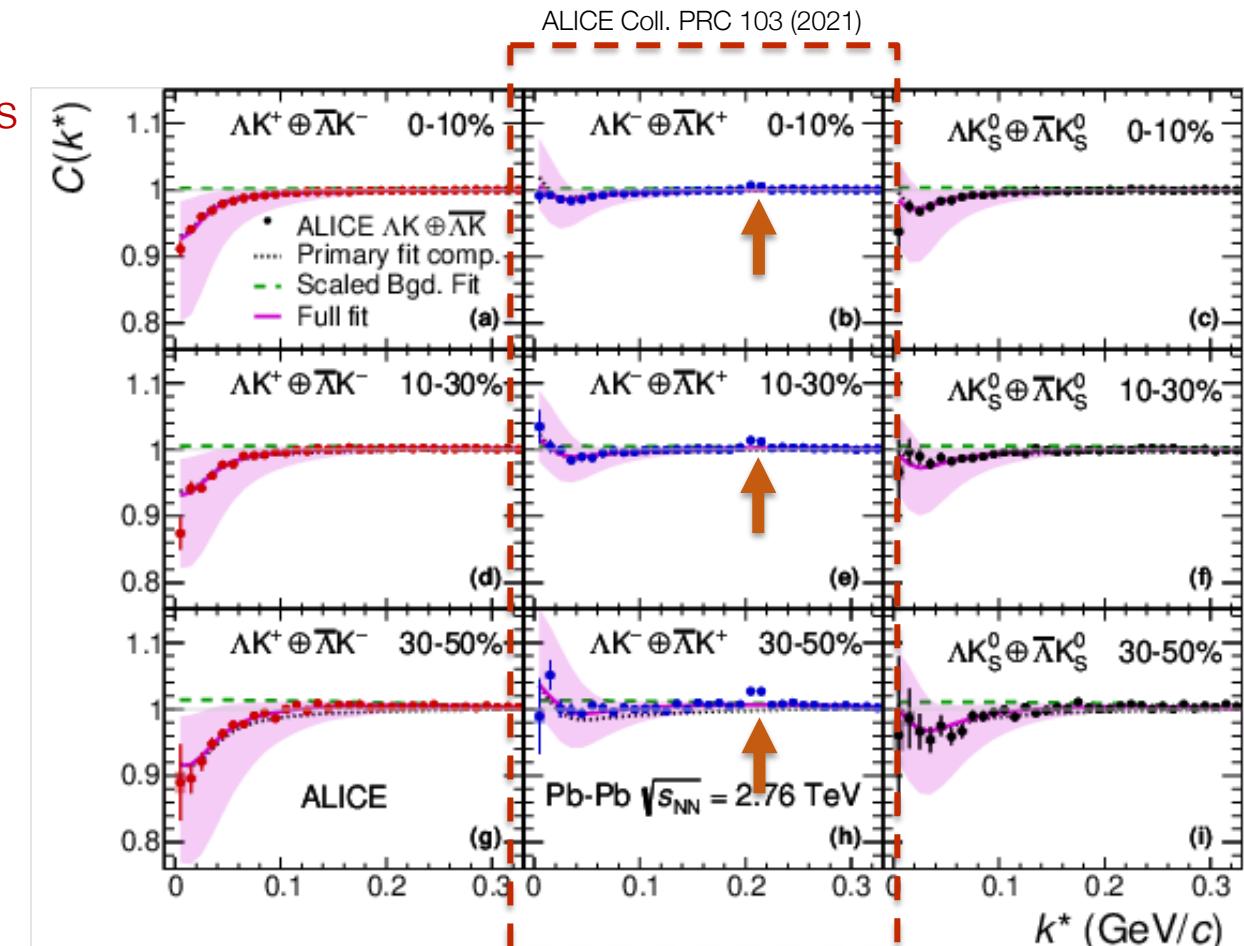


The ΛK^- correlation in large colliding systems

- Delivered the **first measured scattering parameters**
- Femtoscopy in large colliding systems as Pb–Pb
 → **Negligible effects of inelastic channels**
 → Presence of **resonances suppressed**
 (evidence of Ω only in most periph. events)

How does the correlation look like in pp collisions?

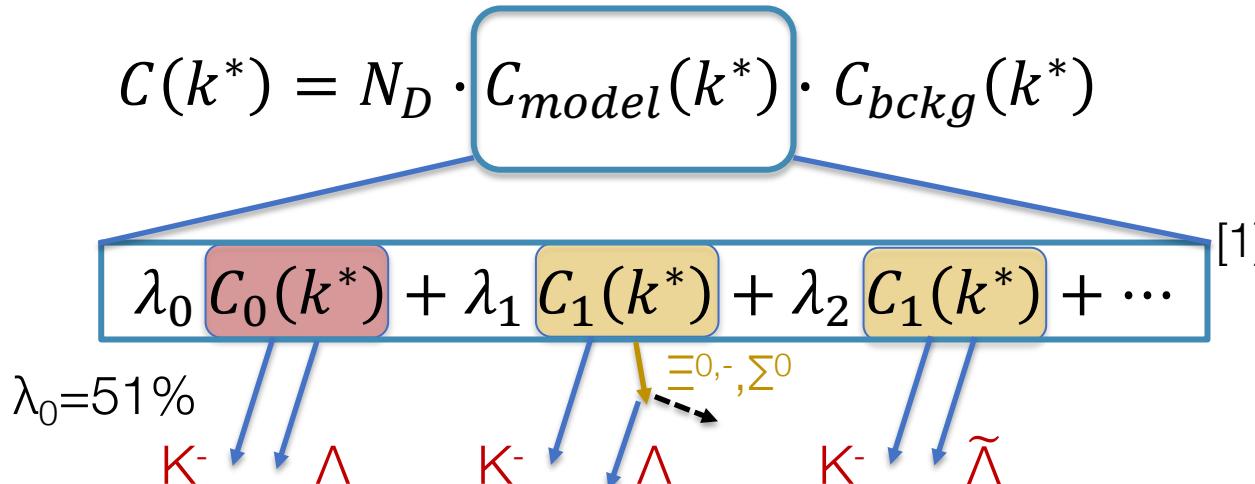
Presence of $\Xi(1620)$ and $\Xi(1690)$?



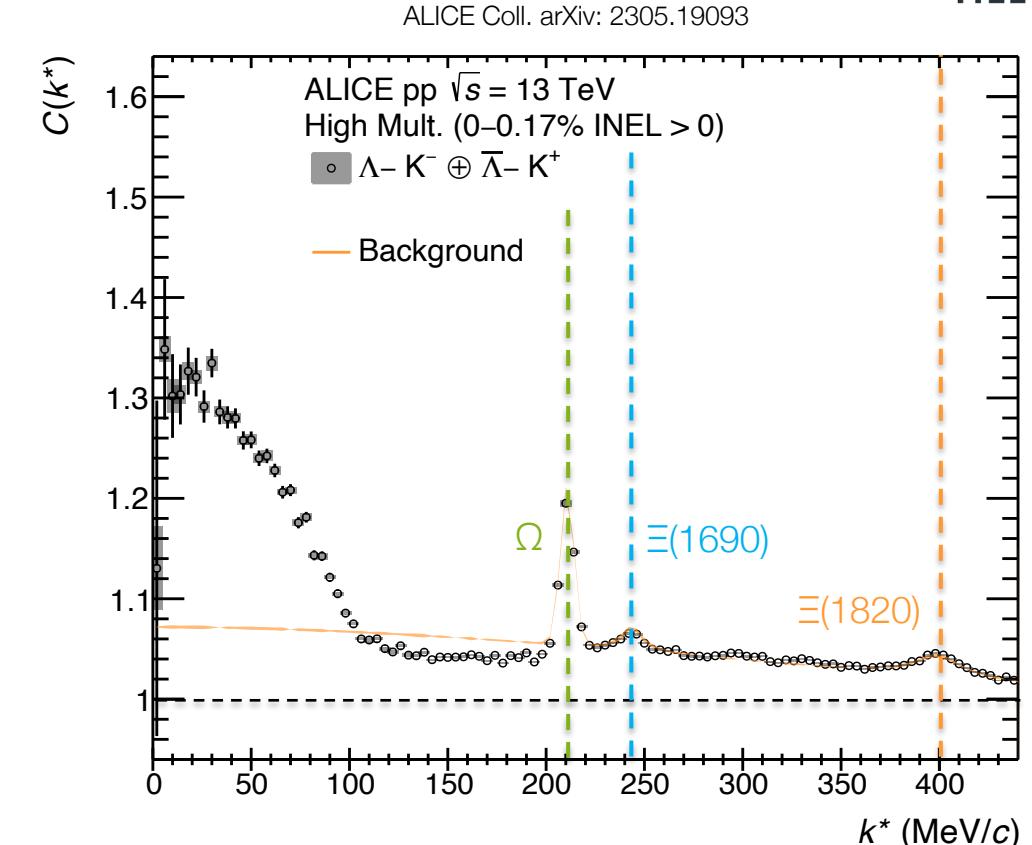
$r \sim 4-10$ fm

The ΛK^- correlation in pp collisions

- Correlation overall above 1 → Attractive interaction



- Contributions from secondaries, impurities, etc..
 → Modeled when possible^[2] or assumed flat
- To understand how to model the genuine part we need to understand the signal at low k^*



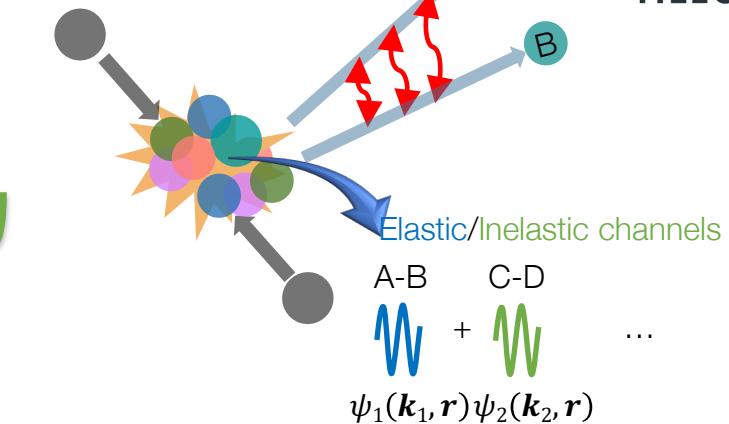
[1] ALICE Coll. Phys.Rev. C99 (2019)

[2] CATS: D. Mihaylov et al., EPJC 78 (2018), 5, 394

Coupled-channels in femtoscopy

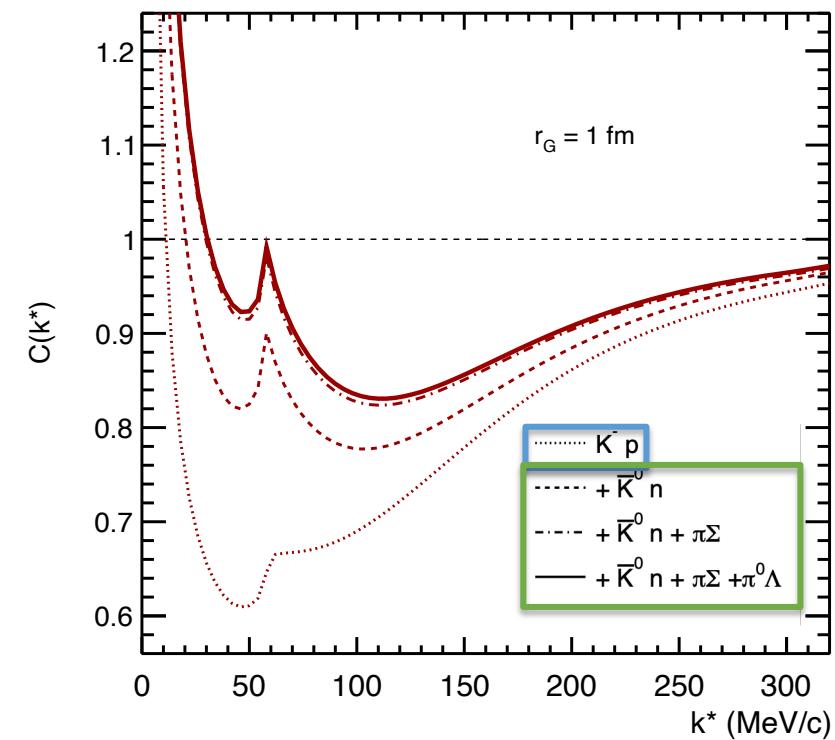
$$C(k^*) = \int S(\vec{r}^*) |\psi_{1 \rightarrow 1}(\vec{k}^*, \vec{r}^*)|^2 d^3 r^* + \sum_{j \neq 1} w_j \int S(\vec{r}^*) |\psi_{j \rightarrow 1}(\vec{k}_j^*, \vec{r}^*)|^2 d^3 r^*$$

 elastic inelastic
 A-B → A-B C-D → A-B, ...



- Wave functions $\Psi_{j \rightarrow 1}(k_j^*, r^*)$
 - Above threshold: modify the shape of CF
→ cusp structure e.g. $\bar{K}^0 n$
 - Below threshold: increase the strength of CF
→ shift upward of CF e.g. $\Sigma \pi$

How does this scenario changes with the source size?



- Lednicky-Lyuboshits analytical formula
 - assuming a gaussian source
 - relies on the asymptotic behaviour of wf
→ **scattering parameters** as inputs, eff.
- CATS framework
 - local potentials, wavefunctions, gaussian and beyond sources
 - relies on the exact wavefunction
→ behaviour at short-distances

$$\Psi_k(\vec{r}) = R_k(r)Y(\theta),$$

$$C_{\text{LL}}(k^*) = 1 + \frac{1}{2} \left| \frac{f}{r_0} \right|^2 \left[1 - \frac{d_0}{2\sqrt{\pi}r_0} \right] + \frac{2\mathcal{R}[f]F_1(2k^*r_0)}{\sqrt{\pi}r_0} - \frac{\mathcal{I}[f]F_2(2k^*r_0)}{r_0}$$

$$f(k^*) \approx \left(f_0^{-1} + \frac{1}{2} d_0 k^{*2} - ik^* \right)^{-1}$$

- might break down for small sources
widely used in large colliding systems

$$C_{\text{th}}(k^*) = \int S(\vec{k}^*, \vec{r}^*) \left| \Psi(\vec{k}^*, \vec{r}^*) \right|^2 d^3 r^*.$$

- works for small and large sources

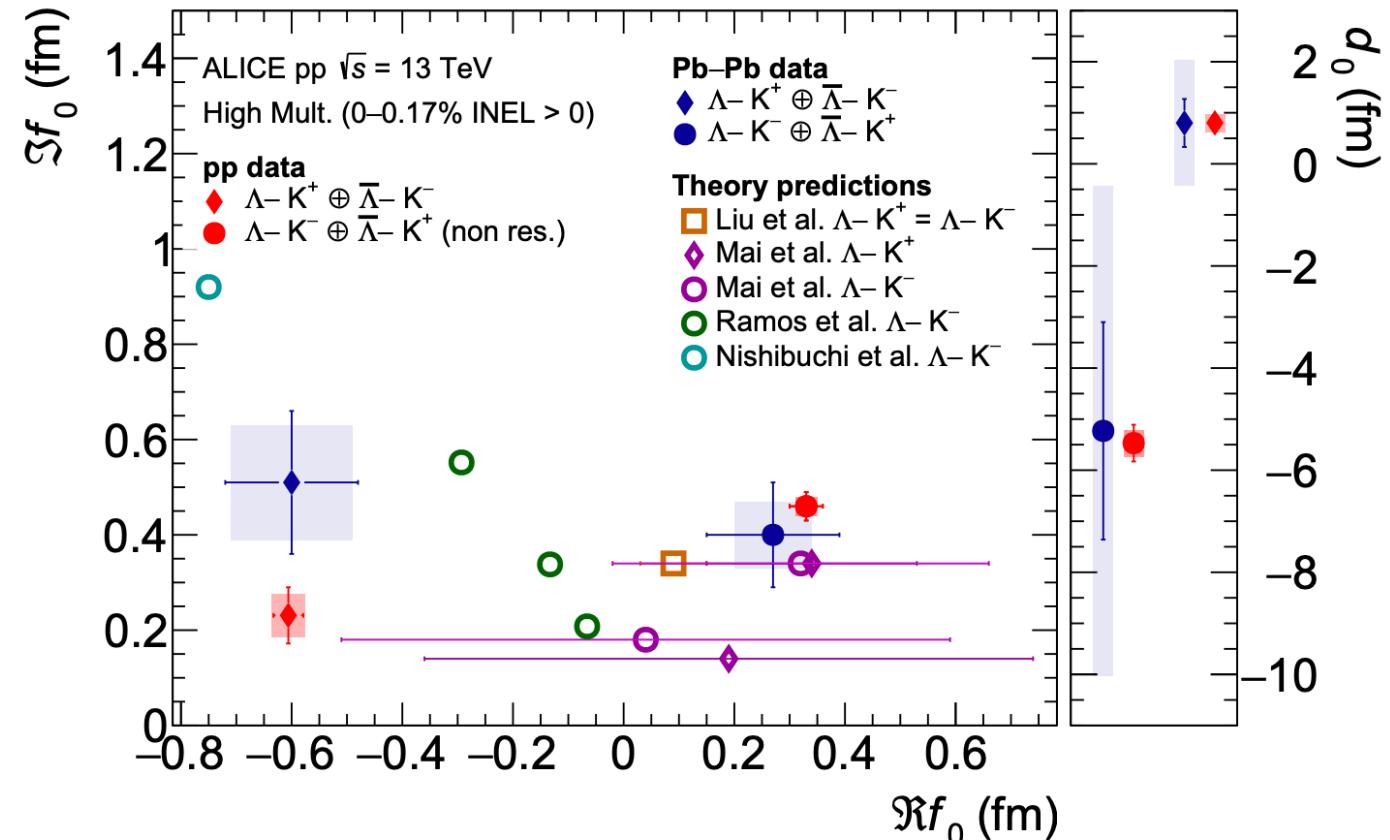
R. Lednicky, V.L. Lyuboshits Sov. J. Nucl. Phys. 35 (1982)

CATS Framework: D. Mihaylov et al., Eur. Phys. J. C78 (2018) 394

Scattering parameters for ΛK^-

ALICE Coll. arXiv: 2305.19093

- Indication of an attractive non-resonant interaction
→ In agreement with ALICE Pb–Pb results^[1]
- Available models far from converging on similar results
 - Parameters fixed based on SU(3) flavour symmetry, isospin symmetry
 - Mainly anchored to πN or $\bar{K}N$ data
 - $\Xi(1620)$ typically lying below threshold
- High-precision data to constrain effective chiral theories and to understand the $\Xi(1620)$ nature



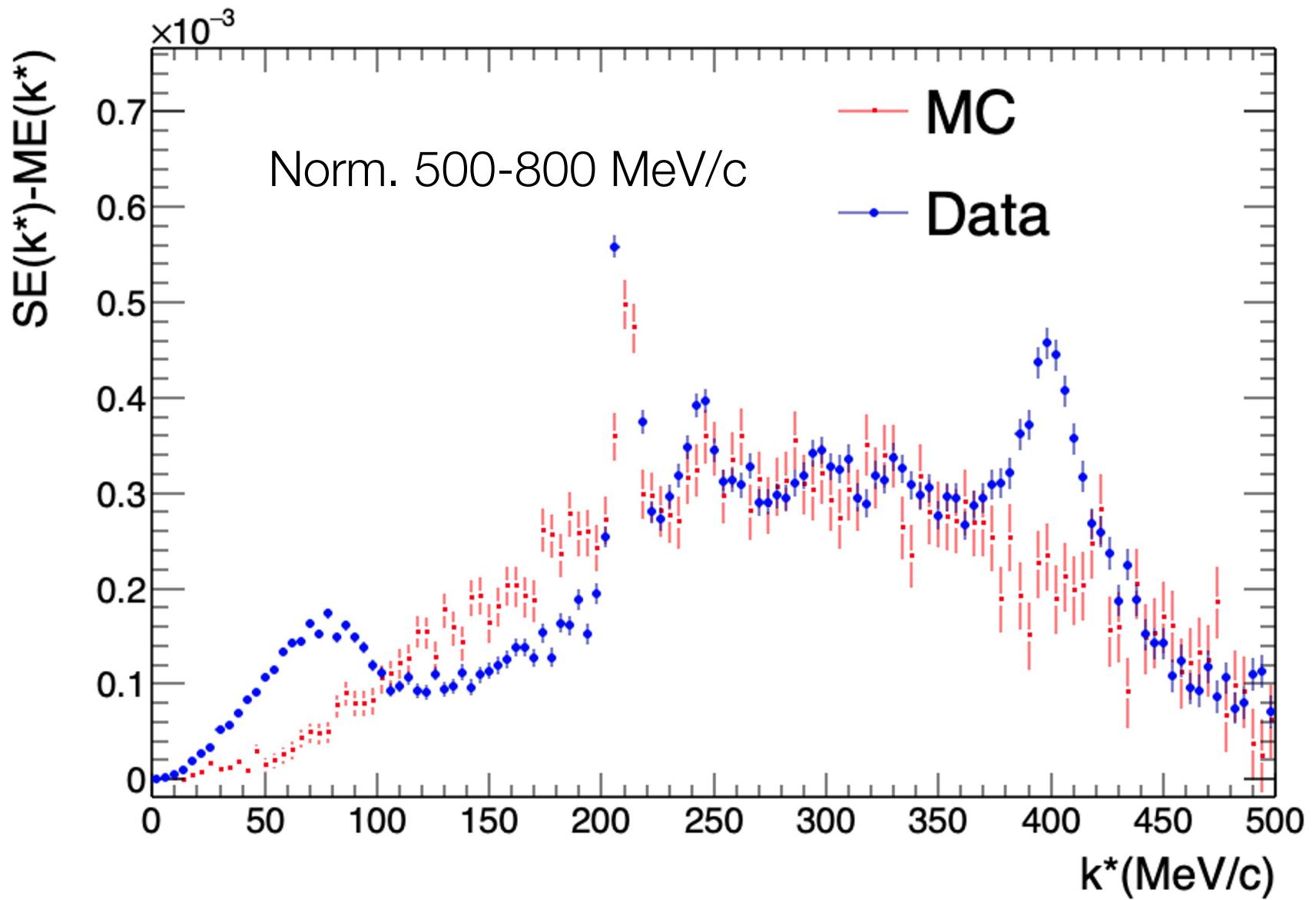
U_XPT at LO: Ramos et al. PRL 89 (2002), Nishibuchi et al. EPJ Web Conf 271 (2022)
 xPT at NLO: Liu et al. PRD 75 (2007), Mai et al. PRD 80 (2009)

The approach adopted is the same as in a typical resonance analysis:

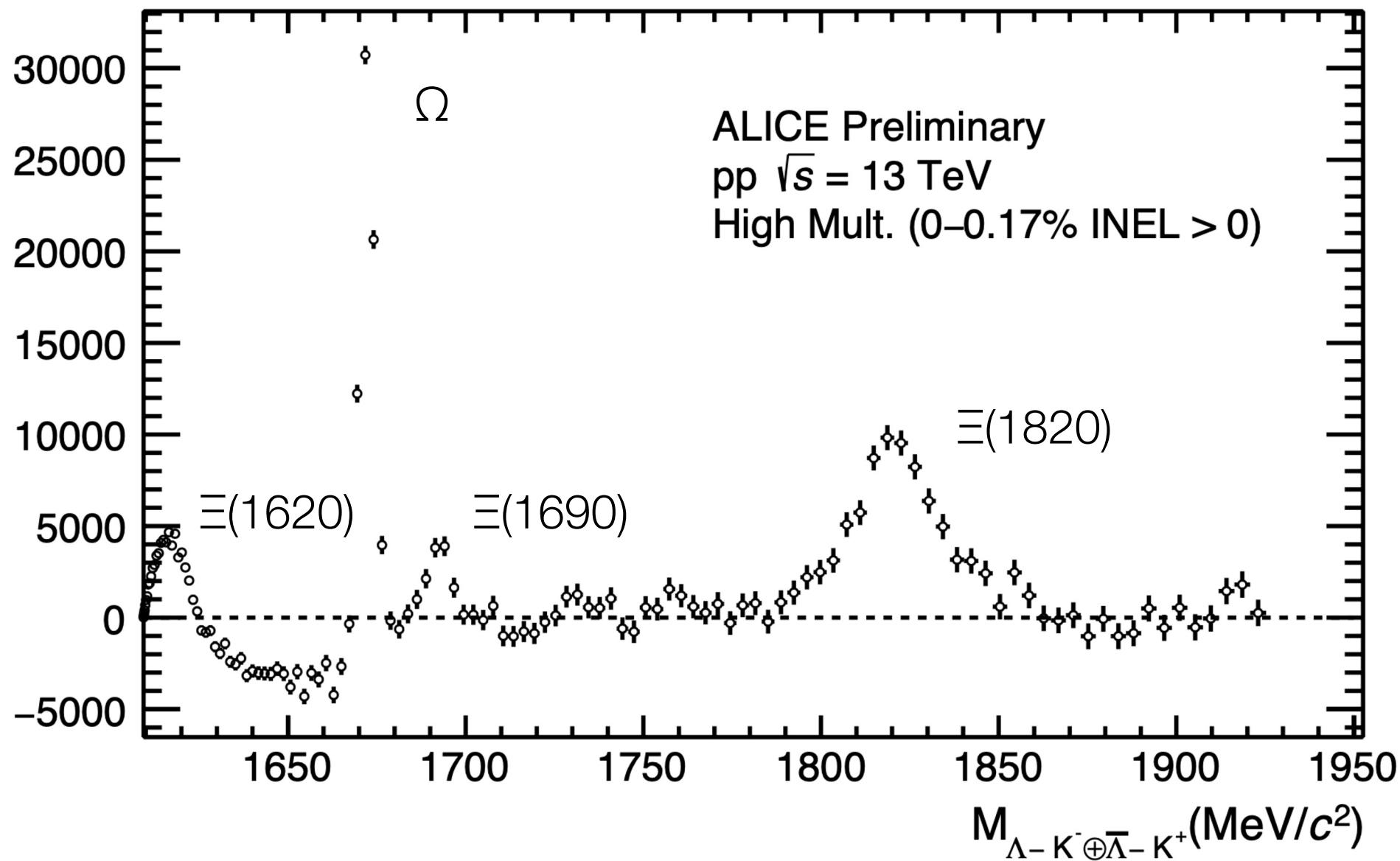
1. Raw signal: Pair (ΛK^-) and AntiPair ($\bar{\Lambda} \bar{K}^+$) for SE data
2. Subtract the uncorrelated background: as in resonance analysis we use the ME data, normalized to SE in a region outside the signal (k^* in 500-800 MeV/c)
3. Subtract the uncorrelated background
 - a. $(SE-ME)_{\text{data}}$
4. The residual background from mini-jet is left, and we can use MC to subtract it
 - a. as done for data we obtain $(SE-ME)_{\text{MC}}$
 - b. fitting $(SE-ME)_{\text{MC}}$ with Pol4 and subtracting it from $(SE-ME)_{\text{data}}$
5. Obtaining the final IM spectrum as a function of the energy

- Values in agreement with PDG and recent measurements
 - $\Xi(1690)/\Xi(1820)$ in $\Xi_b \rightarrow J/\Psi \Lambda K^-$ by *LHCb Coll. Sci.Bull.* 66 (2021) 1278-1287
 - $\Xi(1820)$ by *BESIII Phys. Rev. Lett.* 124 (2020)
 - $\Xi(1820)$ preliminary ALICE Coll. in pp HM

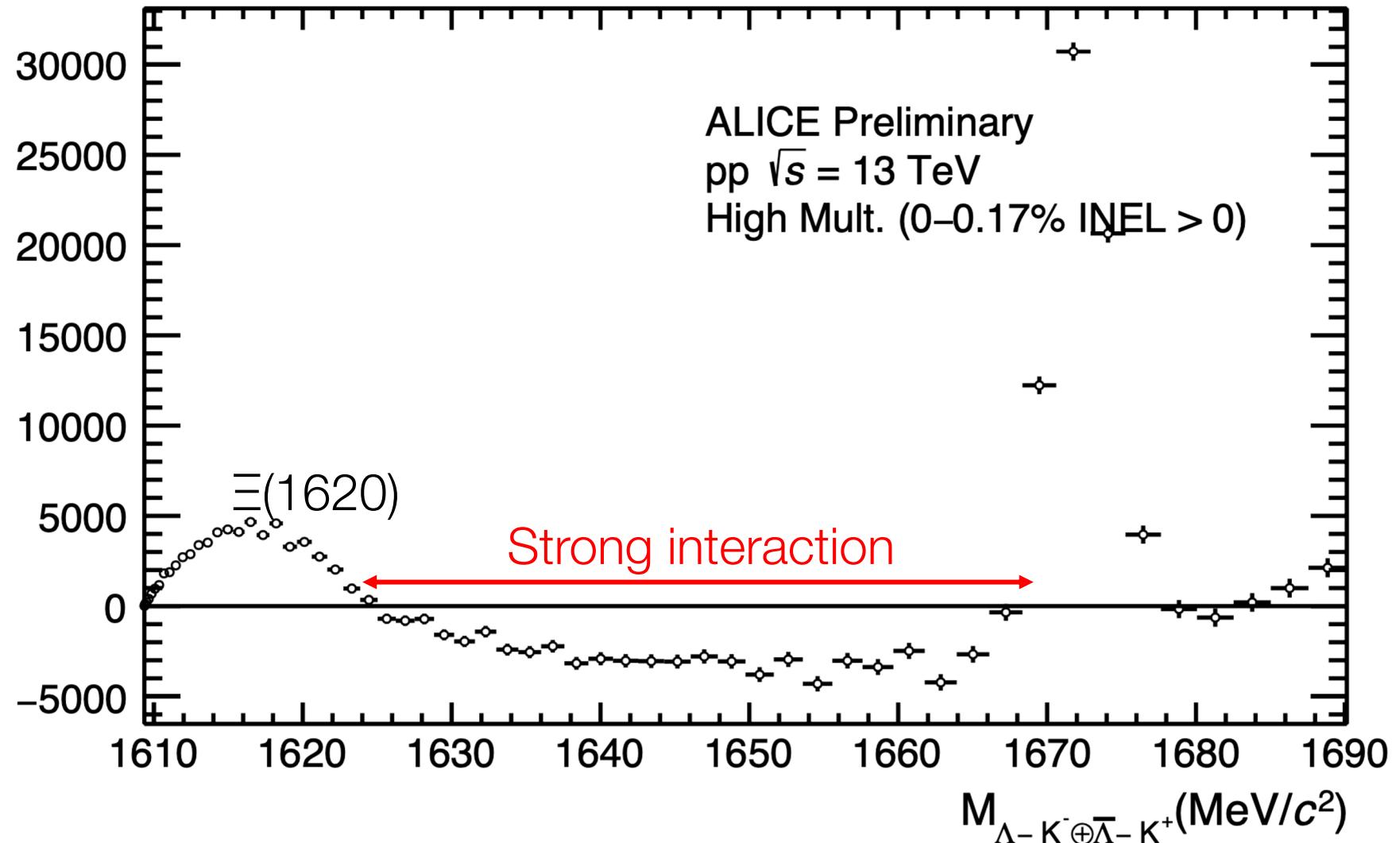
Res.	Mass (MeV)	Width	PDG
Ω	$1672.19 \pm 0.06(\text{stat}) \pm 0.02(\text{syst})$	$3.5 \pm 0.06(\text{stat}) \pm 0.01(\text{syst})$	$M=1672.45 \pm 0.29$
$\Xi(1690)$	$1692.39 \pm 0.6(\text{stat}) \pm 0.18(\text{syst})$	$14.17 \pm 2.8(\text{stat}) \pm 1.1(\text{syst})$	$M=1690 \pm 10$ $\Gamma=20 \pm 15$
$\Xi(1820)$	$1819.57 \pm 0.6(\text{stat}) \pm 0.2(\text{syst})$	$31.73 \pm 2.8(\text{stat}) \pm 1.9(\text{syst})$	$M=1823 \pm 5$ $\Gamma=24^{+15}_{-10}$

Invariant mass in k^* with residual background

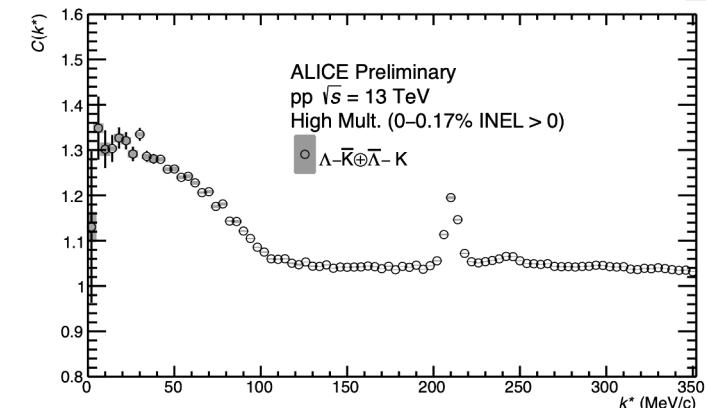
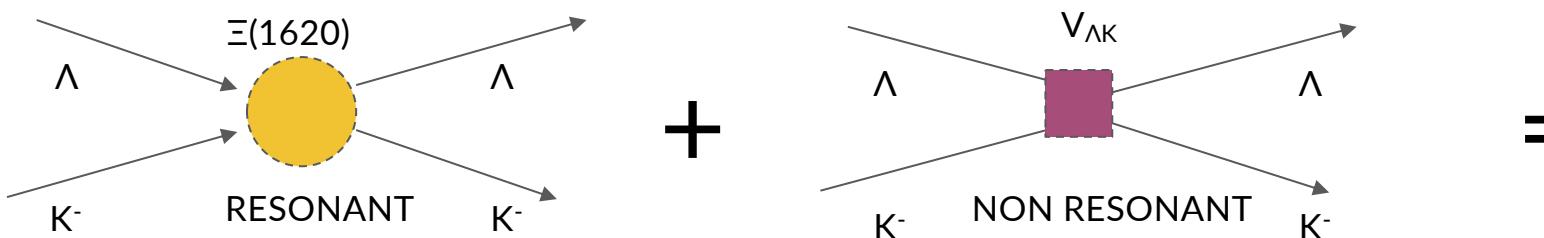
Invariant mass spectrum



Invariant mass spectrum zoom



ΛK^- correlation: including the $\Xi(1620)$ resonance



$$C_{model}(k^*) = \lambda_{gen} C_{gen}(k^*) + \lambda_{\Xi K} C_{\Xi K}(k^*) + \lambda_{flat}$$

- Modeled with Lednicky-Lyuboshits analytical formula

$$C_{gen}(k^*) = w C_{non-res}(k^*) + (1 - w) C_{res}(k^*)$$

- $C_{non-res} \rightarrow$ LL with ERE scatt. amplitude
- $C_{res} \rightarrow$ LL with Flattè-like scatt. amplitude ([F. Giacosa et al. Eur.Phys.J.A 57 \(2021\) 12, 336](#))

$$f(k^*) = \frac{-2\tilde{\Gamma}_2}{E^2 - M^2 + i\tilde{\Gamma}_1 \sqrt{E^2 - E_{thr.1}^2} + i\tilde{\Gamma}_2 \sqrt{E^2 - E_{thr.2}^2}}$$

ch. 1 = π^-
ch. 2 = ΛK^-

Constraining the S=-2 meson-baryon sector

- State-of-the-art U_XPT at NLO available^[1]
 - Low energy constants (LECs) fixed to S=-1 sector^[2]
 - Two sets of subtraction constants (SCs) values
 - Widths of $\Xi(1620)$ too large wrt to Belle's results^[3]

**Use the high-precision femtoscopic data
to fix LECs and SCs!**

Work in collaboration with:

Dr. A. Feijoo, Dr. I. Vidana, Prof. A. Ramos,
Prof. F. Giacosa,
Prof. T. Hyodo and Dr. Y. Kamiya

[1] A. Feijoo et al., PLB 841 (2023)

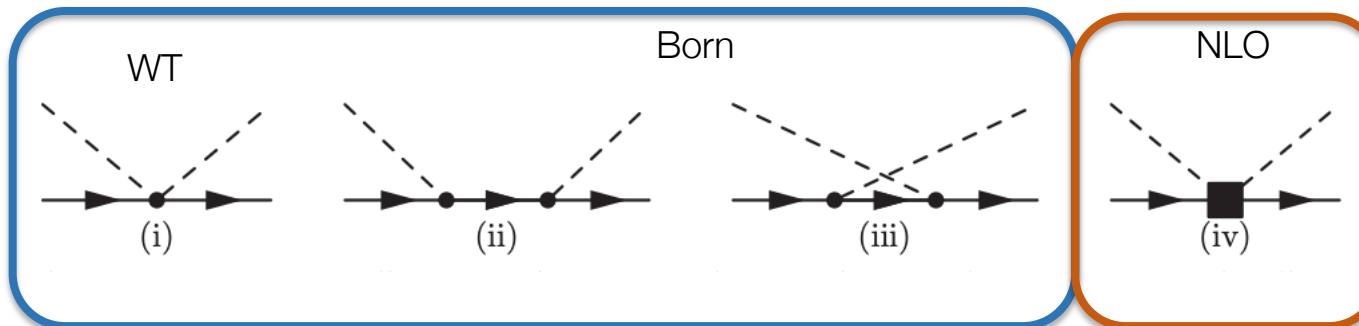
Table 3: Comparison of the pole positions between the models: Model I and Model II (in MeV) with their couplings g_i and the corresponding modulus found in $J^P = \frac{1}{2}^-$, $(I, S) = (\frac{1}{2}, -2)$.

Model I	$\Xi(1620)$		$\Xi(1690)$	
M [MeV]	1599.95		1683.04	
Γ [MeV]	158.88		11.51	
	g_i	$ g_i $	g_i	$ g_i $
$\pi\Xi$	$2.09 + i1.00$	2.32	$-0.30 - i0.12$	0.33
$\bar{K}\Lambda$	$-2.11 - i0.09$	2.11	$-0.49 + i0.05$	0.50
$\bar{K}\Sigma$	$-0.90 + i0.34$	0.97	$1.57 - i0.24$	1.59
$\eta\Xi$	$-0.23 + i0.13$	0.26	$0.74 - i0.11$	0.74

Model II	$\Xi(1620)$		$\Xi(1690)$	
M [MeV]	1608.51		1686.17	
Γ [MeV]	170.00		29.72	
	g_i	$ g_i $	g_i	$ g_i $
$\pi\Xi$	$2.11 + i1.07$	2.37	$-0.36 - i0.24$	0.43
$\bar{K}\Lambda$	$-2.10 - i0.09$	2.10	$-0.81 + i0.02$	0.81
$\bar{K}\Sigma$	$-0.86 + i0.38$	0.94	$2.26 + i0.03$	2.26
$\eta\Xi$	$-0.19 + i0.12$	0.23	$1.04 - i0.07$	1.04

[2] A. Feijoo et al PRC 99 (2019)

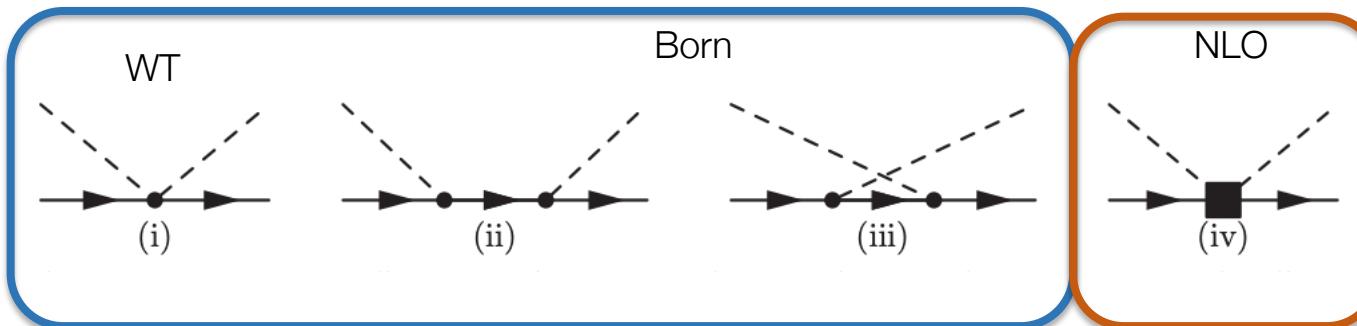
[3] Belle Coll. PRL 122 (2019)



$$\begin{aligned} \mathcal{L}_{\phi B}^{(1)} = & i\langle \bar{B}\gamma_\mu [D^\mu, B] \rangle - M_0\langle \bar{B}B \rangle - \frac{1}{2}D\langle \bar{B}\gamma_\mu\gamma_5\{u^\mu, B\} \rangle \\ & - \frac{1}{2}F\langle \bar{B}\gamma_\mu\gamma_5[u^\mu, B] \rangle, \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{\phi B}^{(2)} = & b_D\langle \bar{B}\{\chi_+, B\} \rangle + b_F\langle \bar{B}[\chi_+, B] \rangle + b_0\langle \bar{B}B \rangle\langle \chi_+ \rangle \\ & + d_1\langle \bar{B}\{u_\mu, [u^\mu, B]\} \rangle + d_2\langle \bar{B}[u_\mu, [u^\mu, B]] \rangle \\ & + d_3\langle \bar{B}u_\mu \rangle\langle u^\mu B \rangle + d_4\langle \bar{B}B \rangle\langle u^\mu u_\mu \rangle. \end{aligned}$$

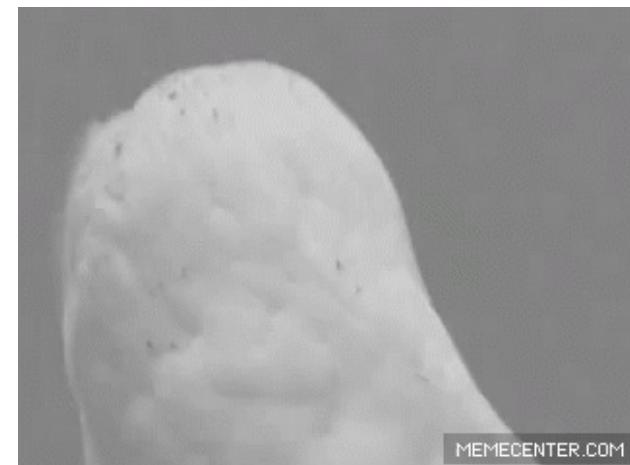
Talk by Dr. A. Feijoo
Sat. 24



$$\begin{aligned} \mathcal{L}_{\phi B}^{(1)} = & i\langle \bar{B}\gamma_\mu[D^\mu, B] \rangle - M_0\langle \bar{B}B \rangle - \frac{1}{2}D\langle \bar{B}\gamma_\mu\gamma_5\{u^\mu, B\} \rangle \\ & - \frac{1}{2}F\langle \bar{B}\gamma_\mu\gamma_5[u^\mu, B] \rangle, \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{\phi B}^{(2)} = & b_D\langle \bar{B}\{\chi_+, B\} \rangle + b_F\langle \bar{B}[\chi_+, B] \rangle + b_0\langle \bar{B}B \rangle\langle \chi_+ \rangle \\ & + d_1\langle \bar{B}\{u_\mu, [u^\mu, B]\} \rangle + d_2\langle \bar{B}[u_\mu, [u^\mu, B]] \rangle \\ & + d_3\langle \bar{B}u_\mu \rangle\langle u^\mu B \rangle + d_4\langle \bar{B}B \rangle\langle u^\mu u_\mu \rangle. \end{aligned}$$

Overall 12 input parameters to be fixed!



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