



M E S O N 2 0 2 1

16th International Workshop on Meson Physics

17th - 20th May 2021, online via ZOOM

Organized by Jagiellonian University Kraków, Forschungszentrum Jülich,
INFN-LNF Frascati, Institute of Nuclear Physics PAS

Organizers:

Catalina Curceanu	INFN-LNF Frascati	(Chairman)
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In-medium spectral functions of vector and axial-vector mesons from aFRG flow equations

Lorenz von Smekal

Meson 2021 – 16th International Workshop on Meson Physics

18 May 2021



Bundesministerium
für Bildung
und Forschung

HFHF Helmholtz
Forschungsakademie
Hessen für FAIR

CRC-TR 211
Strong-interaction matter
under extreme conditions

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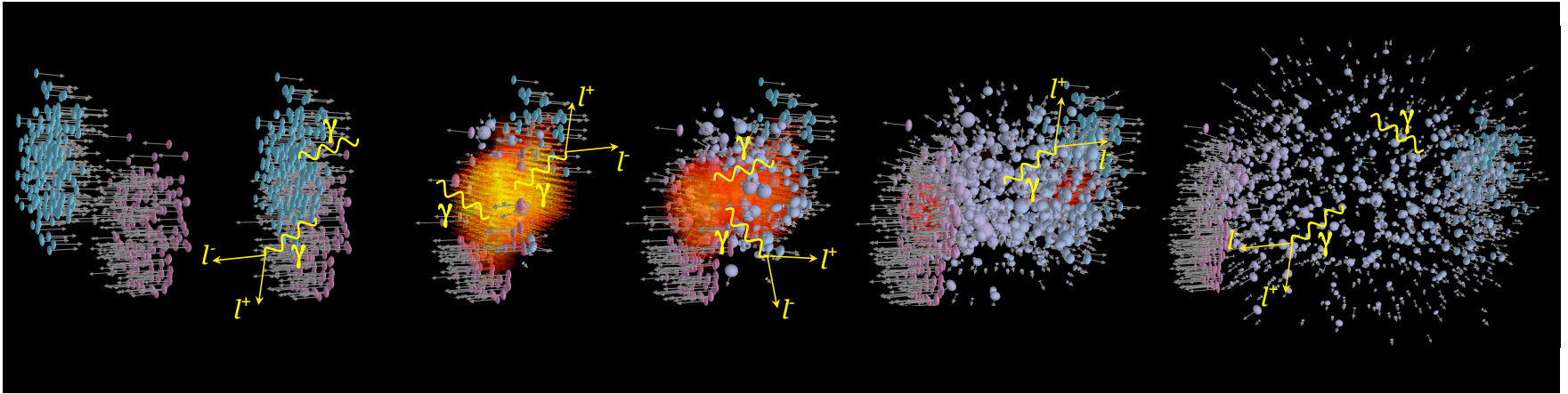
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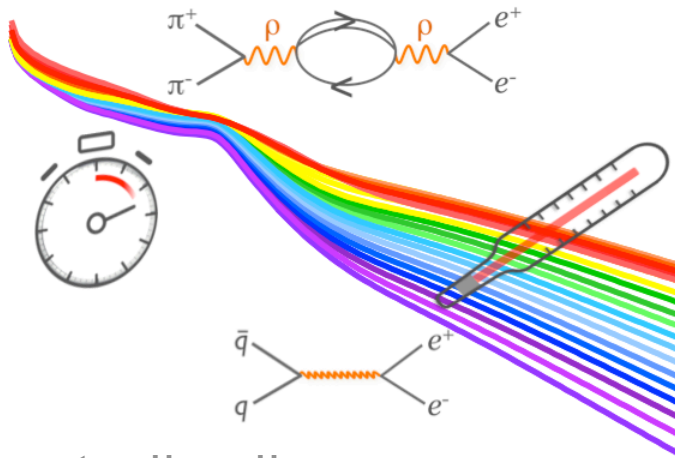
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MESON2021 Workshop, Uniwersytet Jagielloński,
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- Dileptons in HIC
- GiBUU transport simulations for SIS18 energies
- Spectral functions from aFRG flows
- Vector and axial-vector spectral functions
- Effective hadronic theory for dense nuclear matter
- Summary and Outlook



T. Galatyuk et al., Physik Journal 17 (2018) no. 10

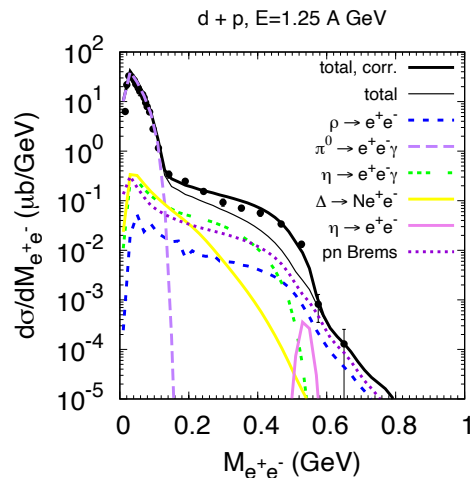
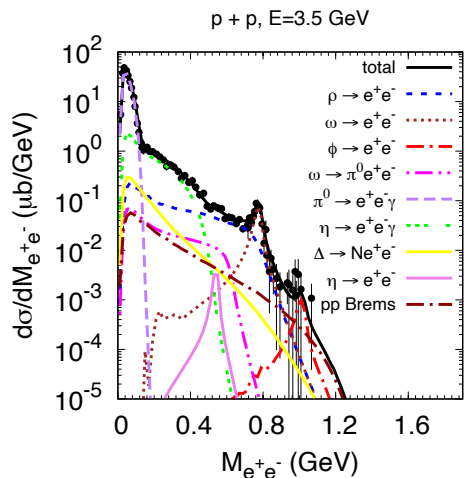
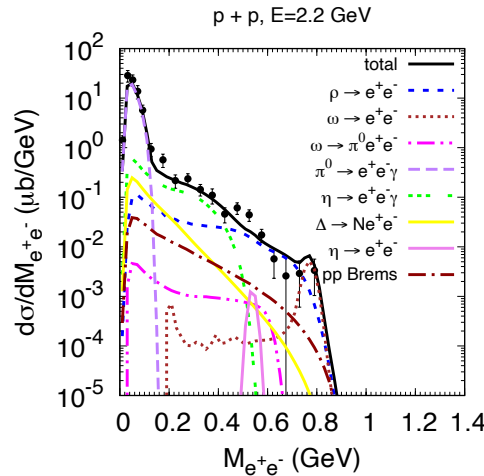
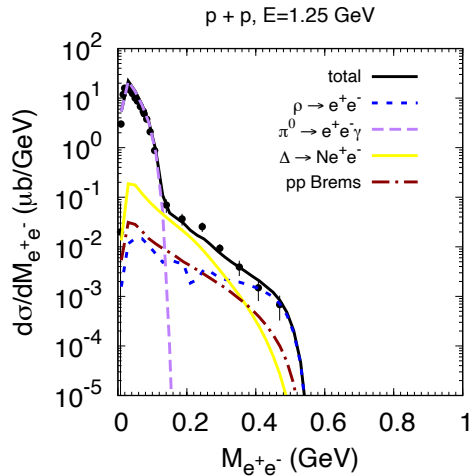


courtesy H. van Hees

- from all stages of the collision
- measure temperature in QGP, lifetime of fireball...

• simulate with GiBUU transport model

O. Buss, T. Gaitanos, K. Gallmeister et al.
(A. Larionov), Phys. Rept. 512 (2012) 1

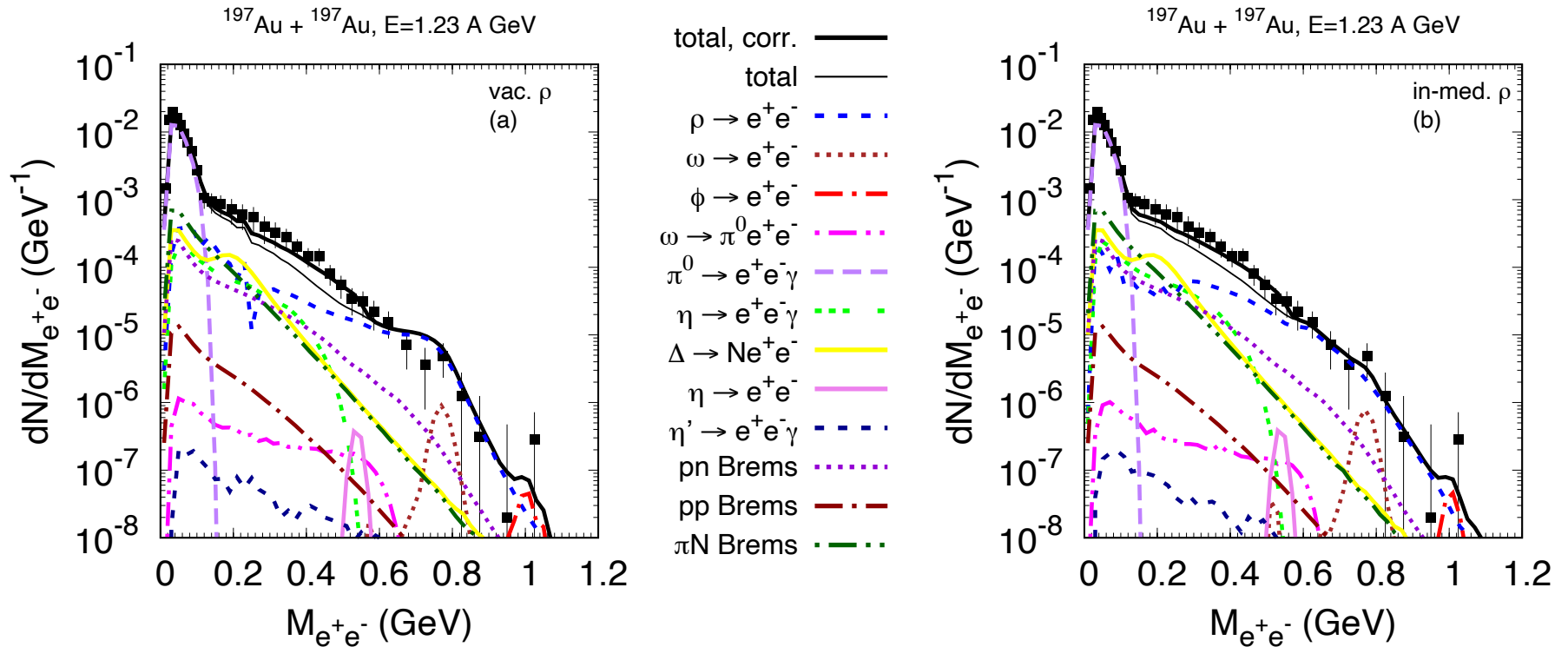


Data: HADES collaboration,
PLB 690 (2010) 118
PRC 85 (2012) 054005
EPJA 48 (2012) 64
PRC 95 (2017) 065205

A. Larionov, U. Mosel & L.v.S.,
Phys. Rev. C 102 (2020) 064913

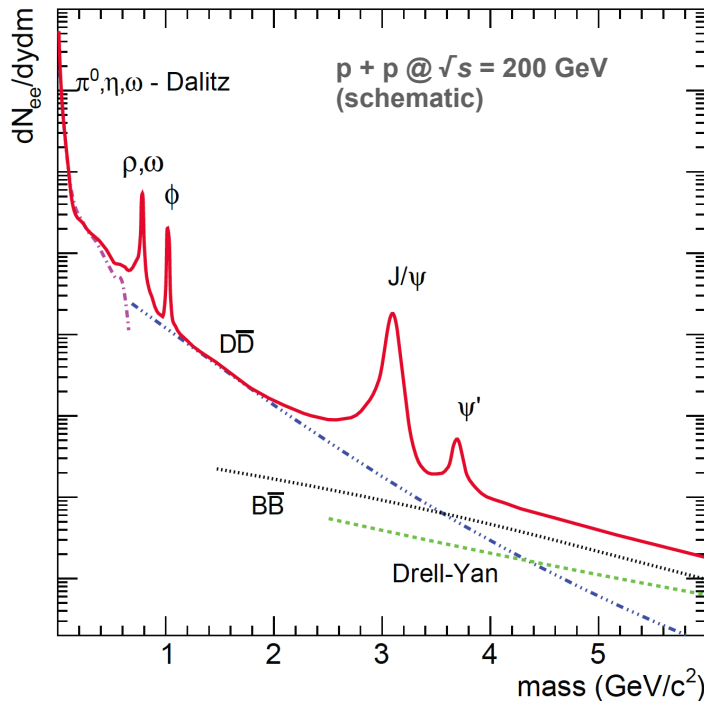
• Dileptons from GiBUU

Data: HADES Collaboration,
Nature Physics 15 (2019) 1040



A. Larionov, U. Mosel & L.v.S., Phys. Rev. C 102 (2020) 064913

A. Drees, NPA 830 (2009) 435



dilepton rate (local thermal equilibrium):

$$\frac{dN_{ll}}{d^4x d^4q} = -\frac{\alpha_{\text{em}}^2}{\pi^3 M^2} \frac{1}{3} g_{\mu\nu} \text{Im} \Pi_{\text{em}}^{\mu\nu}(M, |\vec{q}|; \mu, T)$$

electromagnetic correlator:

$$\Pi_{\text{em}}^{\mu\nu}(q; \mu, T) = -i \int d^4x e^{iqx} \theta(x^0) \langle [j_{\text{em}}^\mu(x), j_{\text{em}}^\nu(0)] \rangle$$

vector meson dominance
& quark counting:

$$\text{Im} \Pi_{\text{em}}^{\mu\nu}(M \leq 1\text{GeV}) \sim \text{Im} D_\rho^{\mu\nu} + \frac{1}{9} \text{Im} D_\omega^{\mu\nu} + \frac{2}{9} \text{Im} D_\phi^{\mu\nu}$$

- some theory basics



[courtesy L. Holicki]

commutator of interacting fields:

$$\langle [\phi(x), \phi(0)] \rangle = \int_0^\infty dm^2 \rho(m^2) i\Delta(x; m^2) \quad \leftarrow \text{free fields}$$

Fourier transform:

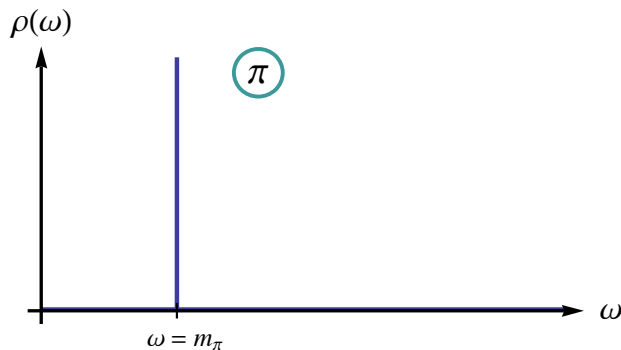
$$\rho(\omega, \vec{p}) = \int d^4x e^{ipx} i\langle [\phi(x), \phi(0)] \rangle$$

spectral function:

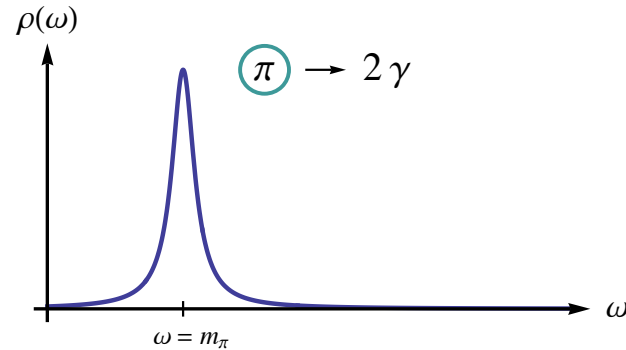
$$\Rightarrow \quad \rho(\omega, \vec{p}) = 2\pi i \epsilon(\omega) \theta(p^2) \rho(p^2)$$

$$\rho(p^2) = (2\pi)^3 \sum_{\psi} \delta^4(p - q_{\psi}) |\langle \Omega | \phi(0) | \psi \rangle|^2, \quad p_0 > 0$$

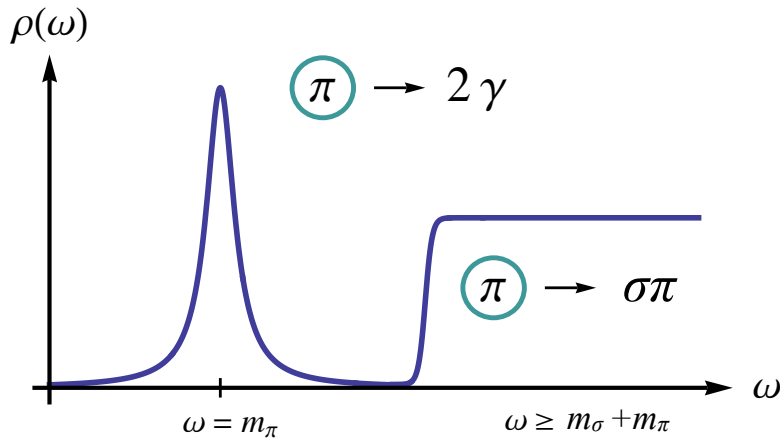
free fields (stable pion):



finite lifetime/width

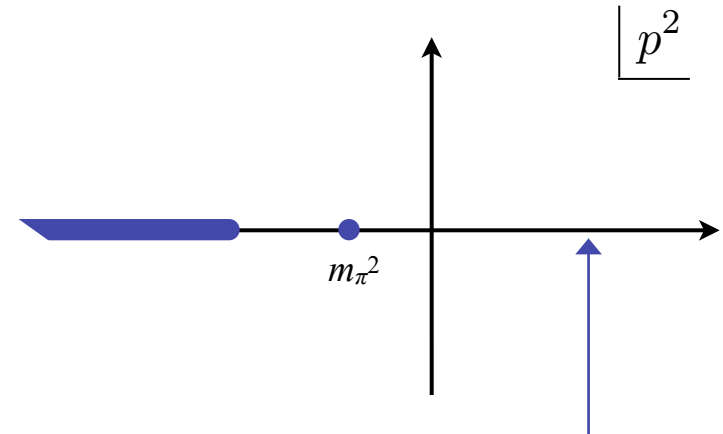


two-particle thresholds:



discontinuity at cut of propagator:

$$D(p) = \int_0^\infty dm^2 \rho(m^2) \frac{1}{p^2 + m^2}$$



retarded, imaginary part: $\rho(p^2) = -\frac{1}{\pi} \text{Im} D_R(p)$

Euclidean space:

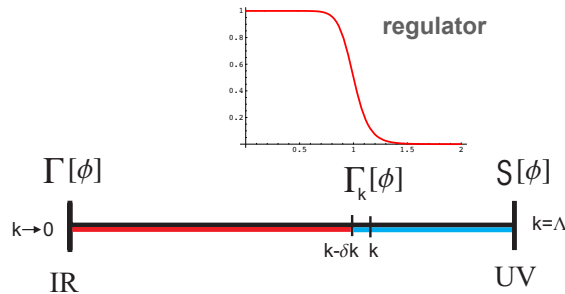
$$p^2 > 0$$

Euclidean data: $D(t, \vec{p} = 0) = \int_0^\infty dm \rho(m^2) \exp\{-mt\}$

(inverse Laplace, try e.g. MEM, but ill-posed numerical problem)

- compute effective (average) action:

Ch. Wetterich, PLB 301 (1993) 90

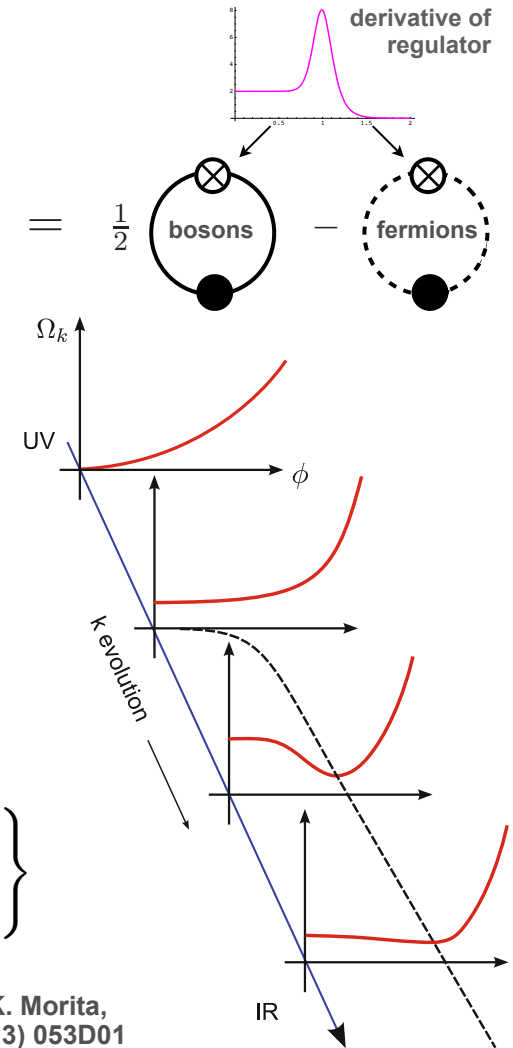


$$k \partial_k \Gamma_k[\phi] = \frac{1}{2} \left(\text{bosons} - \text{fermions} \right)$$

- flow of Landau free energy density:

(quark-meson model, leading order derivative expansion)

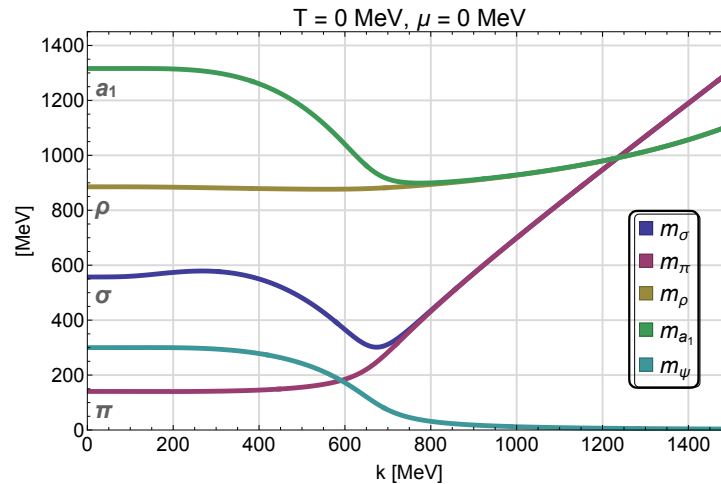
$$\partial_k \Omega_k(T, \mu; \phi^2) = \frac{k^4}{12\pi^2} \left\{ \frac{1}{E_k^\sigma} \coth\left(\frac{E_k^\sigma}{2T}\right) + \frac{3}{E_k^\pi} \coth\left(\frac{E_k^\pi}{2T}\right) - \frac{2N_c N_f}{E_k^q} \left[\tanh\left(\frac{E_k^q - \mu}{2T}\right) + \tanh\left(\frac{E_k^q + \mu}{2T}\right) \right] \right\}$$



- include mixing with density fluctuations

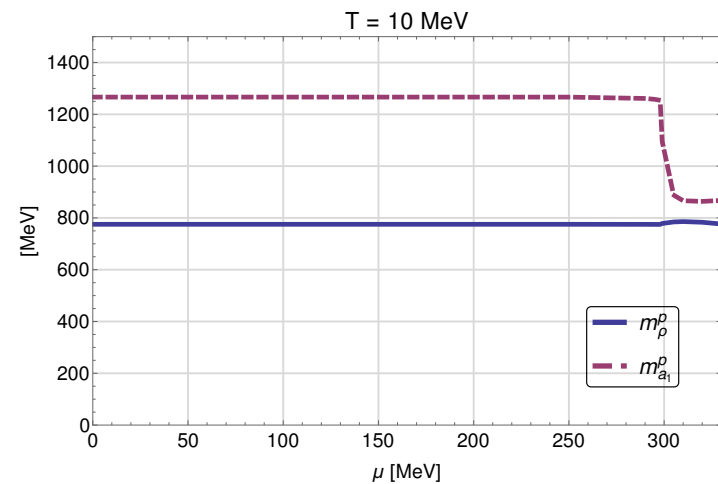
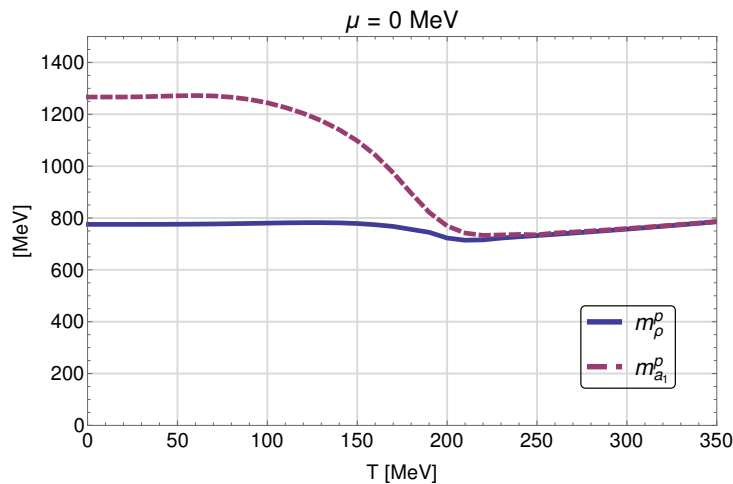
K. Kamikado, T. Kunihiro, K. Morita, A. Ohnishi, PTEP 2013 (2013) 053D01

Flow of Euclidean
(curvature) masses
including
(axial-)vector mesons



• physical pole masses:

chiral restoration
at finite T and μ



Ch. Jung & L.v.S., PRD 100 (2019) 116009

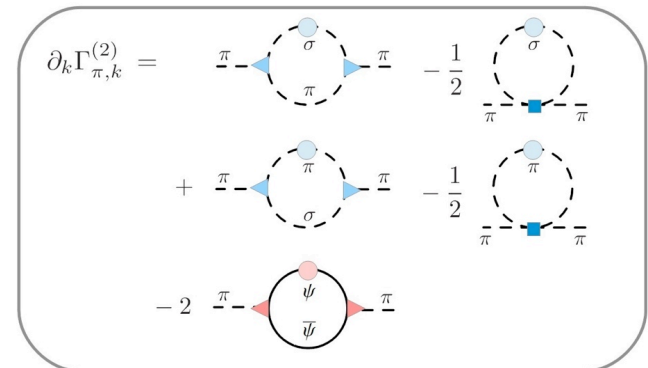
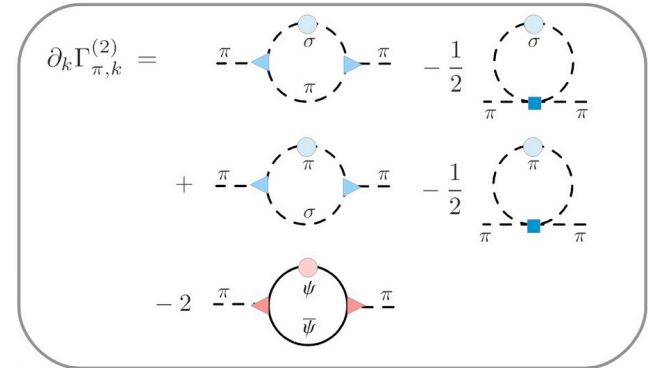
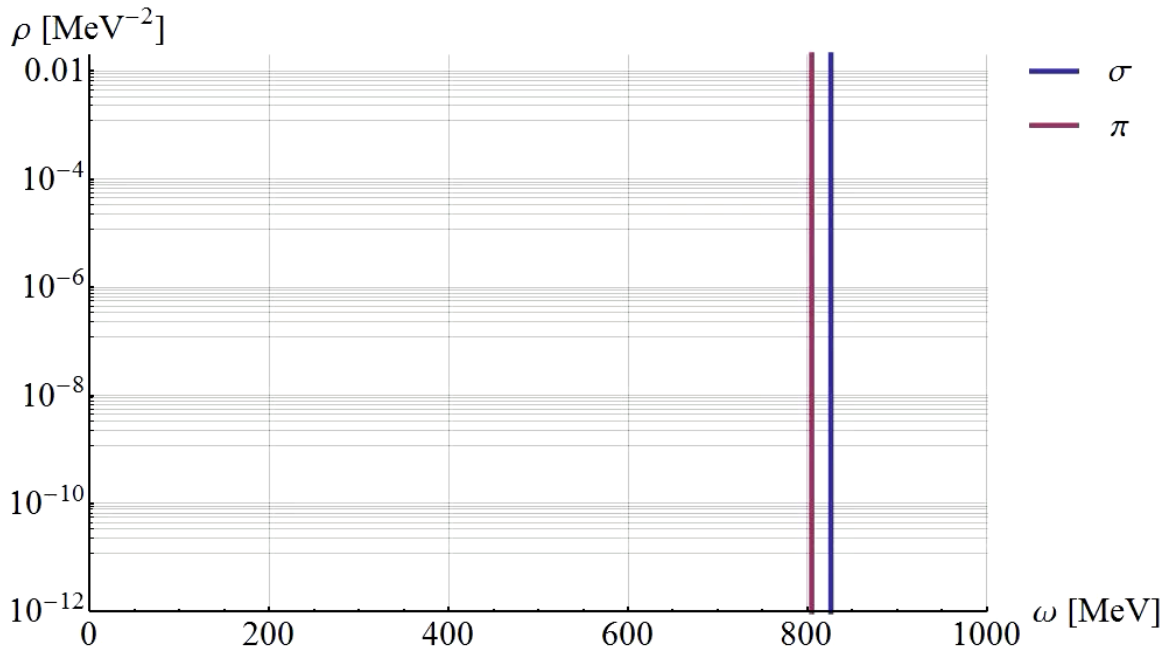
- quark-meson model, $T = \mu = 0$:

$$p_0 = -i(\omega + i\varepsilon) \quad (\text{retarded})$$

$$\rho(\omega, \vec{p}) = -\frac{1}{\pi} \text{Im} D^R(\omega, \vec{p})$$

- for $\varepsilon \rightarrow 0$:

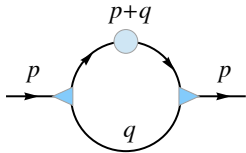
$k=1000 \text{ MeV}$



$$p_0 = -i(\omega + i\varepsilon)$$

K. Kamikado, N. Strodthoff, L.v.S. & J. Wambach, EPJC 74 (2014) 2806

- analytic continuation not unique:
exploit one-loop structure, 3-dim. regulators

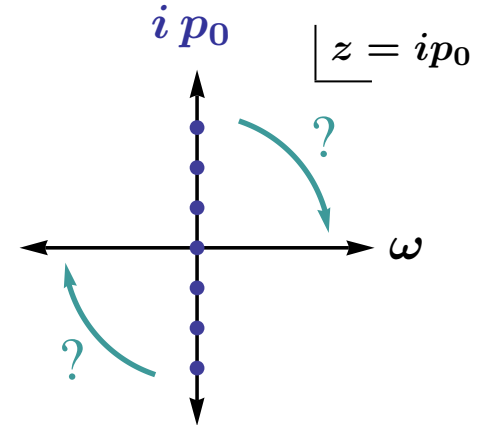


1) Use periodicity in external energy $p_0 = n 2\pi T$:

$$n_{B,F}(E + ip_0) \rightarrow n_{B,F}(E)$$

2) Substitute p_0 by continuous real frequency:

$$\Gamma^{(2),R}(\omega) = - \lim_{\epsilon \rightarrow 0} \Gamma^{(2),E}(p_0 = i\omega - \epsilon)$$



- for $\epsilon \rightarrow 0$:

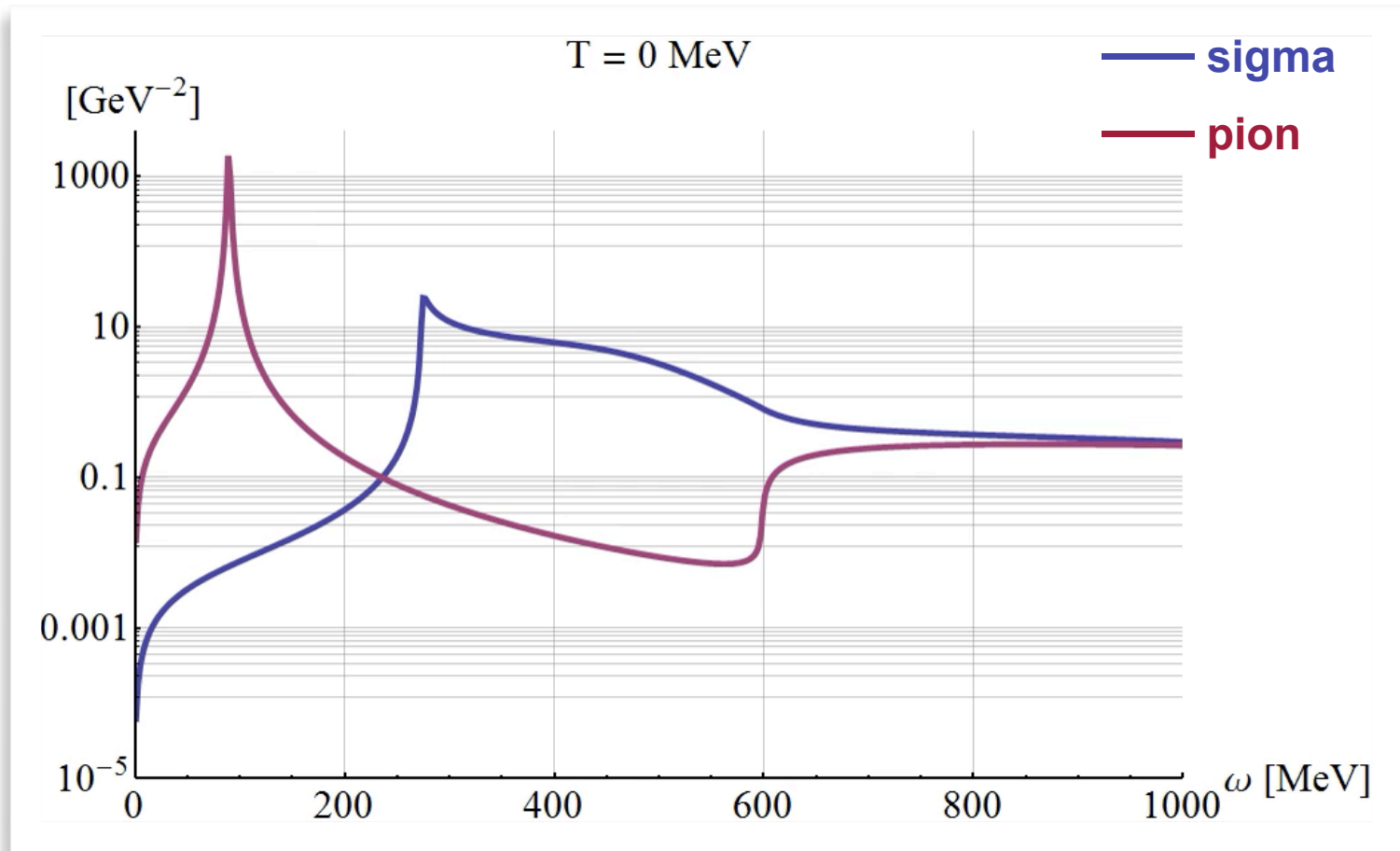
$$p_0 = -i(\omega + i\epsilon)$$

\leadsto Baym-Mermin physical b.c.'s satisfied

$$\rho(\omega, \vec{p}) = -\frac{1}{\pi} \text{Im} D^R(\omega, \vec{p}) = \frac{1}{\pi} \frac{\text{Im} \Gamma^{(2),R}(\omega, \vec{p})}{(\text{Re} \Gamma^{(2),R}(\omega, \vec{p}))^2 + (\text{Im} \Gamma^{(2),R}(\omega, \vec{p}))^2}$$

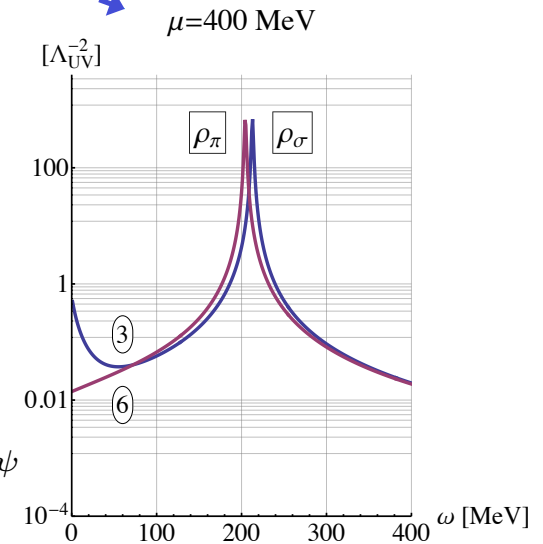
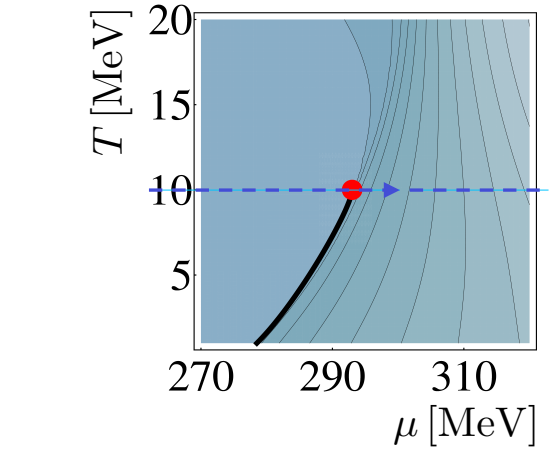
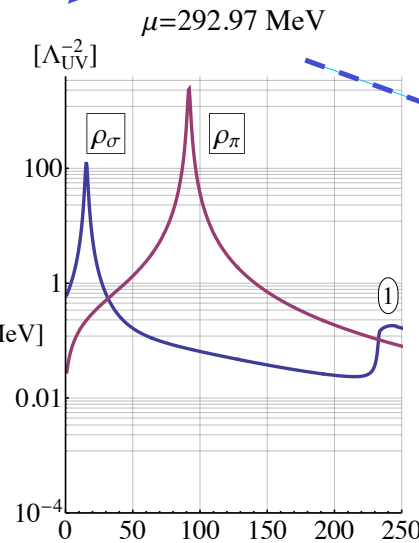
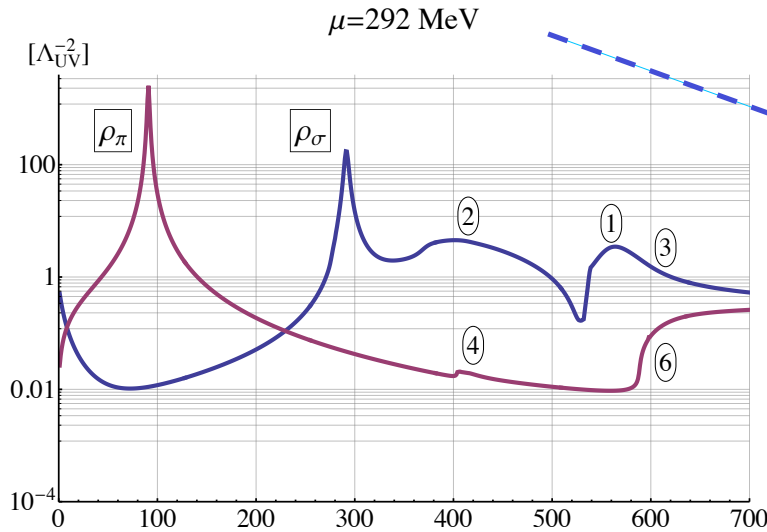
- quark-meson model:

$$\mu = 0$$



A. Tripolt, N. Strodthoff, L.v.S. & J. Wambach, PRD 89 (2014) 34010

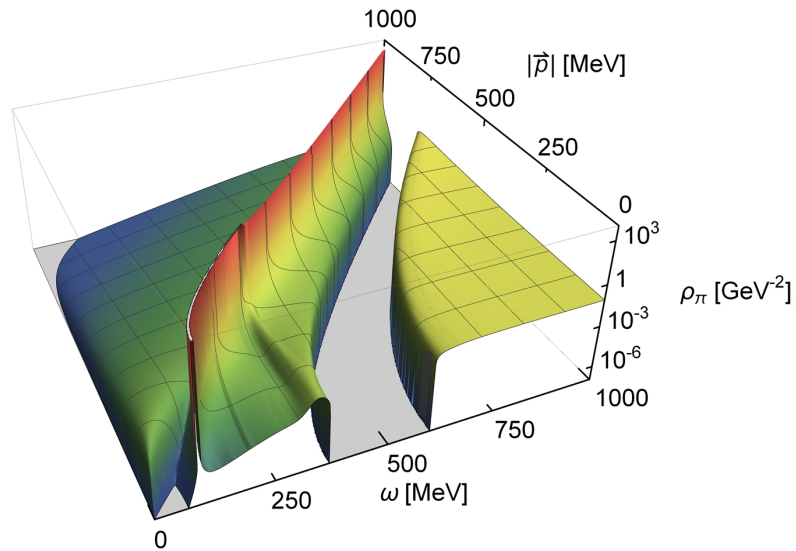
- quark-meson model:
 $T = 10 \text{ MeV}$



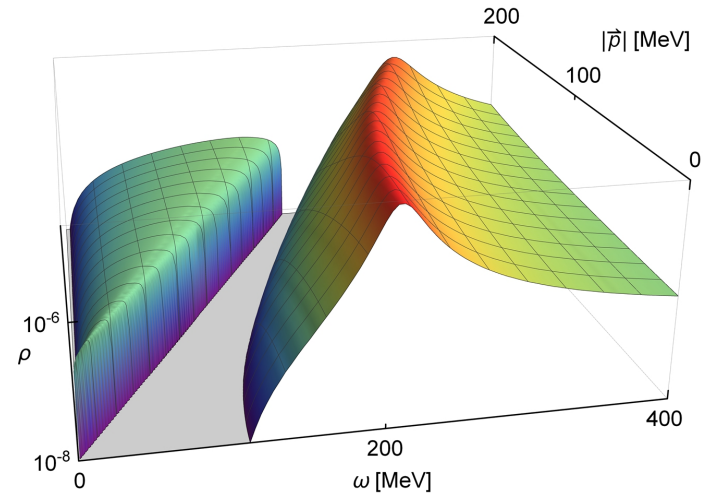
1: $\sigma^* \rightarrow \sigma\sigma$, 2: $\sigma^* \rightarrow \pi\pi$, 3: $\sigma^* \rightarrow \bar{\psi}\psi$, 4: $\pi^* \rightarrow \sigma\pi$, 5: $\pi^* \pi \rightarrow \sigma$, 6: $\pi^* \rightarrow \bar{\psi}\psi$

A. Tripolt, N. Strodthoff, L.v.S. & J. Wambach, PRD 89 (2014) 34010

pion SF $\rho(\omega, \vec{p})$ below T_c



sigma meson SF $\rho(\omega, \vec{p})$ above T_c

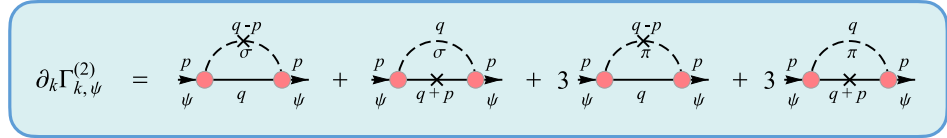


↪ transport coefficients

A. Tripolt, L.v.S. & J. Wambach, PRD 90 (2014) 074031

- aFRG for fermionic two-point functions

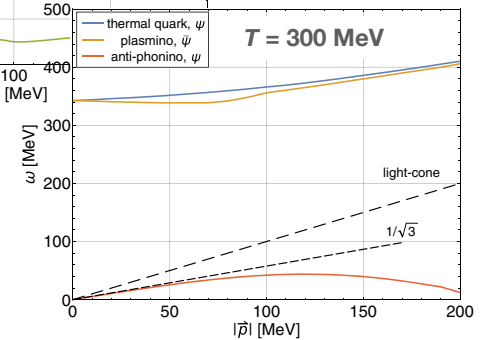
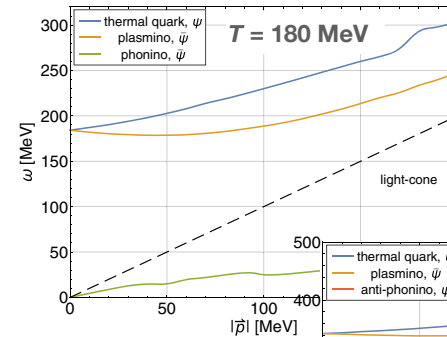
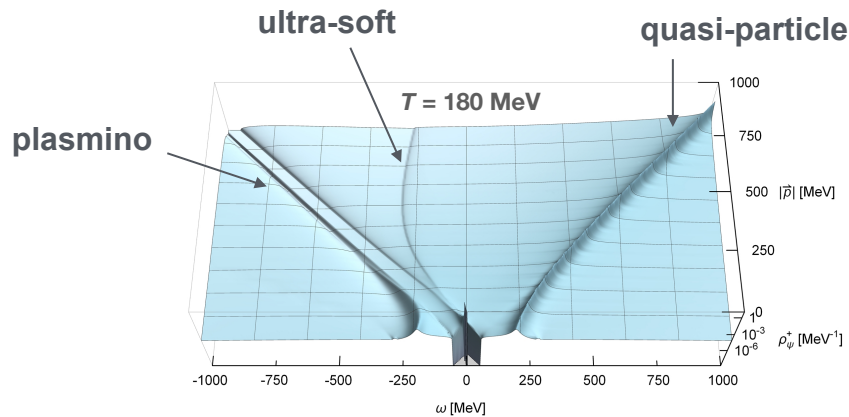
$$\Gamma_{k,\psi}^{(2)}(\omega, \vec{p}) = \gamma_0 C_k(\omega, \vec{p}) + i\vec{\gamma} \cdot \hat{p} A_k(\omega, \vec{p}) - B_k(\omega, \vec{p})$$



- spectral functions

$$\rho_{k,\psi}(\omega, \vec{p}) = \gamma_0 \rho_{k,\psi}^{(C)}(\omega, \vec{p}) + i\vec{\gamma} \cdot \hat{p} \rho_{k,\psi}^{(A)}(\omega, \vec{p}) + \rho_{k,\psi}^{(B)}(\omega, \vec{p})$$

- describe fermionic excitations at finite T



A. Tripolt, J. Weyrich, L.v.S. & J. Wambach, PRD 98 (2018) 094002
 A. Tripolt, D. Rischke, L.v.S. & J. Wambach, PRD 101 (2020) 094010

- **extended linear-sigma model with quarks:**

Ch. Jung, F. Rennecke, A. Tripolt, L.v.S. & J. Wambach, PRD95 (2017) 036020

- **electromagnetic SF from gauging and mixing:**

A. Tripolt, Ch. Jung, N. Tanji, L.v.S. & J. Wambach, NPA 982 (2019) 775

- **include fluctuating (axial-)vectors in aFRG flows for SFs:**

Ch. Jung & L.v.S., PRD 100 (2019) 116009



- **(axial-)vector SFs in hadronic effective theory for dense nuclear matter:**

A. Tripolt, Ch. Jung, L.v.S. & J. Wambach, arXiv:2105.00861 [hep-ph]

- aFRG flows for ρ and a_1 at finite T and μ :

$$\partial_k \Gamma_{\rho\rho,k}^{(2)} =$$

$$+ \text{with fluctuating } \rho \text{ and } a_1$$

$$\partial_k \Gamma_{a_1 a_1, k}^{(2)} =$$

- spectral representation of conserved current:

$$\langle T_{\text{cov}} j_\mu(x) j_\nu(0) \rangle = -i \int_0^\infty ds \frac{\rho(s)}{s} \int \frac{d^4 p}{(2\pi)^4} e^{-ipx} \frac{p^2 g_{\mu\nu} - p_\mu p_\nu}{p^2 - s + i\epsilon}$$

- current-field identity, transverse vector propagator:

$$D_{\mu\nu}^{T,V}(p) = -i \frac{Z}{m_v^2} \frac{p^2 g_{\mu\nu} - p_\mu p_\nu}{p^2 - m_v^2 + i\epsilon} + \dots$$

- Euclidean two-point function, single-particle contribution:

$$\Gamma_{\mu\nu}^{(2)T}(p) = -\frac{m_0^2}{p^4} (p^2 + m_v^2) (p^2 \delta_{\mu\nu} - p_\mu p_\nu) \quad m_{0,k}^2 = m_{v,k}^2 / Z_k$$

- (axial-)vectors from (anti-)selfdual field strengths:

$$\mathcal{L}_0^\rho = -\frac{1}{4} \text{tr} (\partial_\mu \rho_{\mu\nu}) \partial_\sigma \rho_{\sigma\nu} + \frac{m_v^2}{8} \text{tr} \rho_{\mu\nu} \rho_{\mu\nu}$$

- (axial-)vectors from (anti-)selfdual field strengths:

$$\rho_{\mu\nu} = \vec{\rho}_{\mu\nu}^+ \cdot \vec{T}_R + \vec{\rho}_{\mu\nu}^- \cdot \vec{T}_L$$

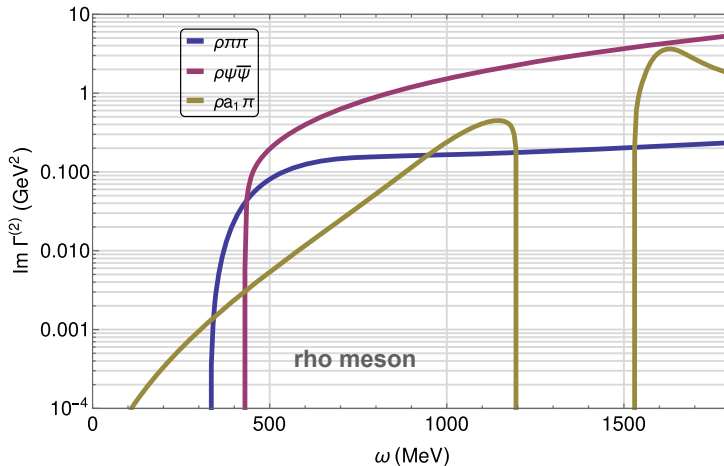
$$\vec{\rho}_\mu = \frac{1}{2m_\nu} \text{tr}(\partial_\sigma \rho_{\sigma\mu} \vec{T}_V)$$

$$\vec{a}_{1\mu} = \frac{1}{2m_\nu} \text{tr}(\partial_\sigma \rho_{\sigma\mu} \vec{T}_A)$$

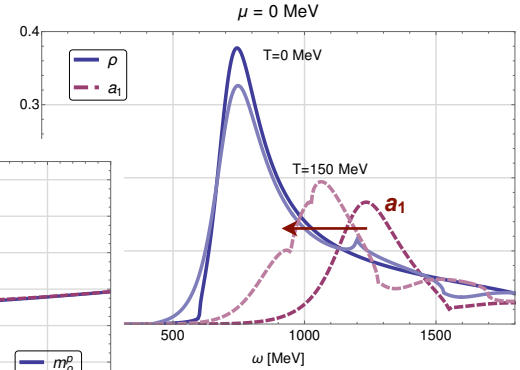
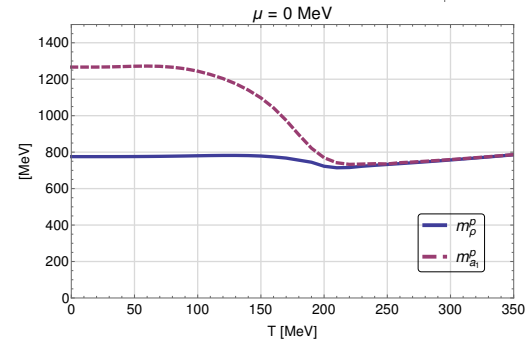
- new processes / imaginary parts:

$\text{Im } \Gamma^{(2)}_\rho(\omega)$

$T = 150 \text{ MeV}, \mu = 0 \text{ MeV}$

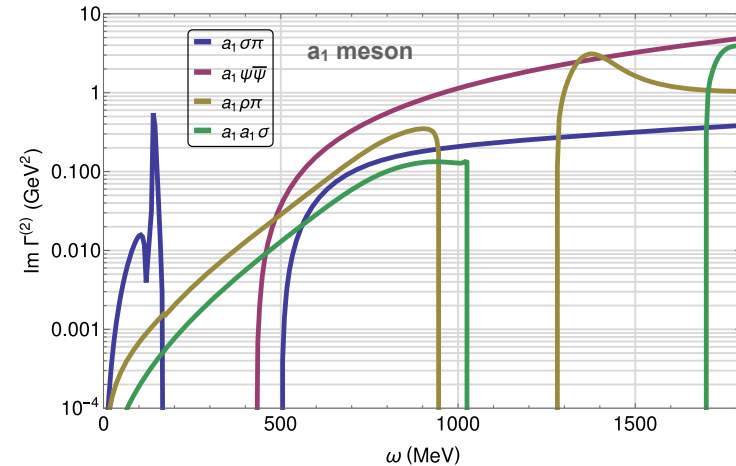


pole masses



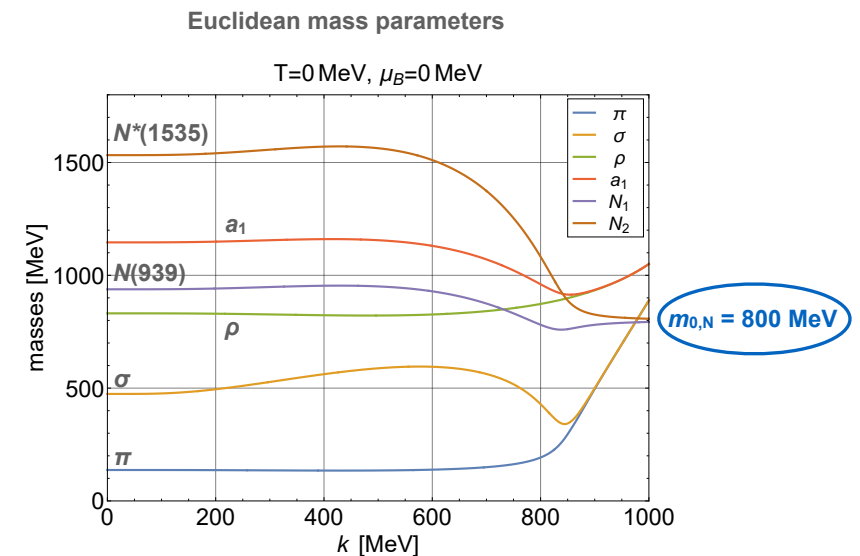
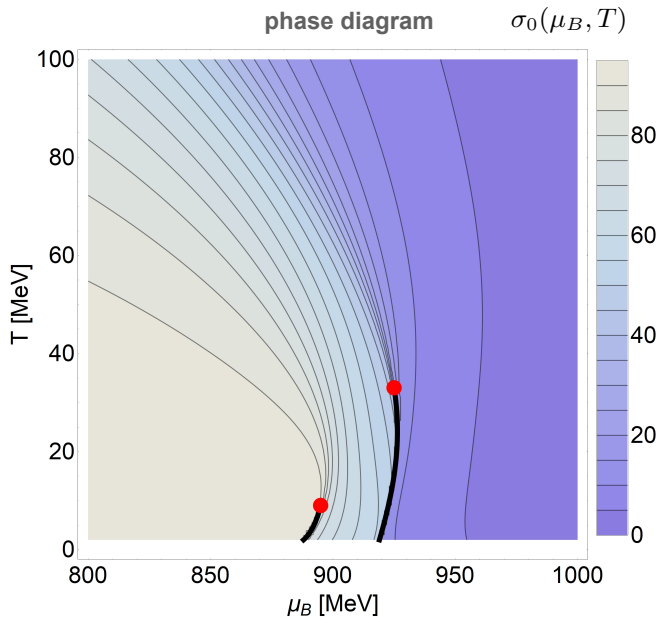
$\text{Im } \Gamma^{(2)}_a(\omega)$

a_1 meson



- gauged linear-sigma model with $N(939)$ and $N^*(1535)$ iso-doublets:

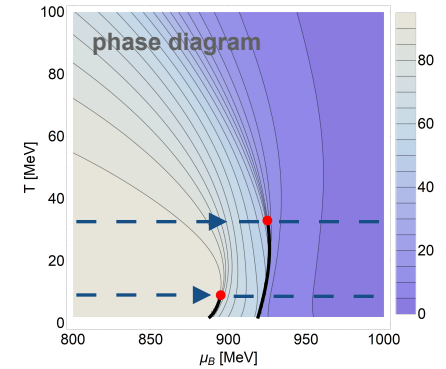
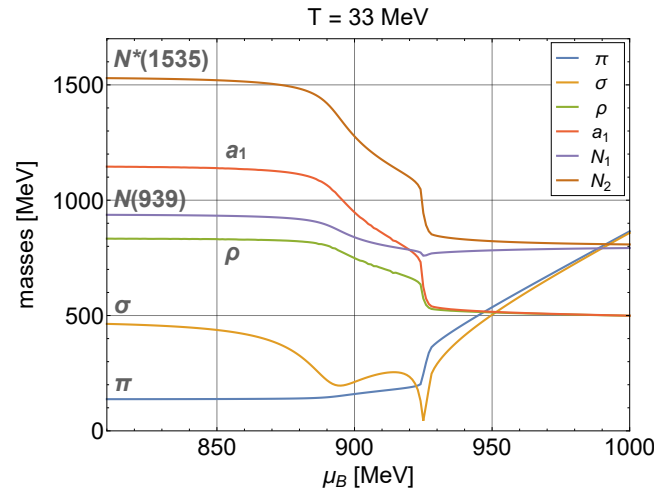
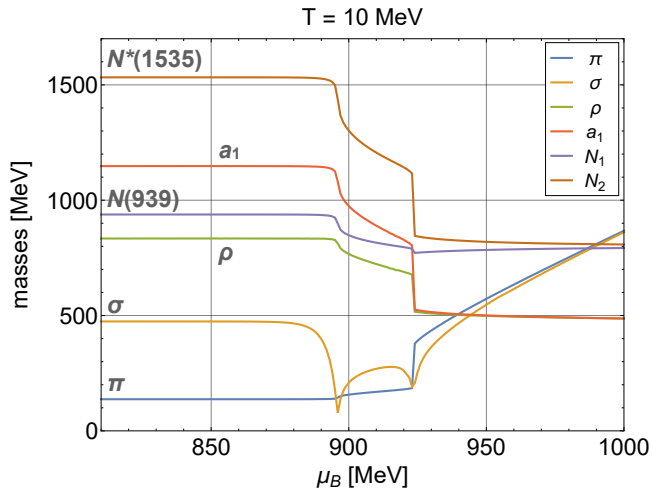
$$\Gamma_k = \int d^4x \left\{ \bar{N}_1 (\not{\partial} - \mu_B \gamma_0 + h_{s,1}(\sigma + i\vec{\tau} \cdot \vec{\pi} \gamma^5) + h_{v,1}(\gamma_\mu \vec{\tau} \cdot \vec{\rho}_\mu + \gamma_\mu \gamma^5 \vec{\tau} \cdot \vec{a}_{1,\mu})) N_1 \right. \\ \left. + \bar{N}_2 (\not{\partial} - \mu_B \gamma_0 + h_{s,2}(\sigma - i\vec{\tau} \cdot \vec{\pi} \gamma^5) + h_{v,2}(\gamma_\mu \vec{\tau} \cdot \vec{\rho}_\mu - \gamma_\mu \gamma^5 \vec{\tau} \cdot \vec{a}_{1,\mu})) N_2 + m_{0,N} (\bar{N}_1 \gamma^5 N_2 - \bar{N}_2 \gamma^5 N_1) \right. \\ \left. + U_k(\phi^2) - c\sigma + \frac{1}{2} (D_\mu \phi)^\dagger D_\mu \phi - \frac{1}{4} \text{tr} \partial_\mu \rho_{\mu\nu} \partial_\sigma \rho_{\sigma\nu} + \frac{m_v^2}{8} \text{tr} \rho_{\mu\nu} \rho_{\mu\nu} \right\}$$



J. Weyrich, N. Strodthoff & L.v.S., PRC 92 (2015) 015214

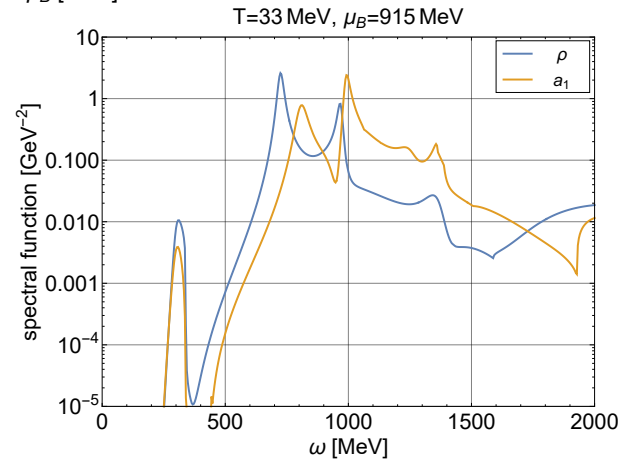
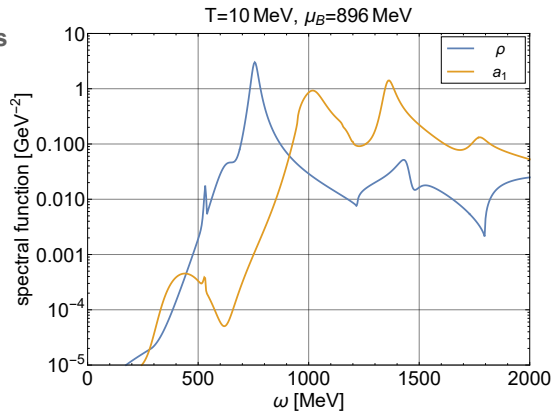
A. Tripolt, Ch. Jung, L.v.S. & J. Wambach, arXiv:2105.00861 [hep-ph]

- mass parameters at finite T and μ :



- (axial-)vector SFs inside nuclear matter:

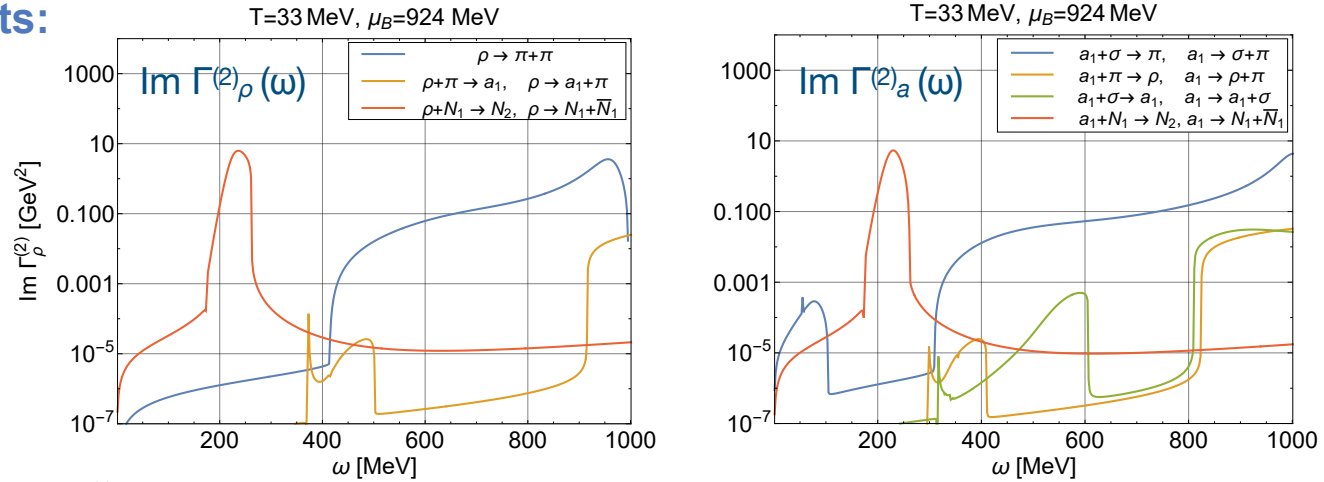
nuclear liquid-gas CEP



approaching chiral CEP inside dense nuclear matter

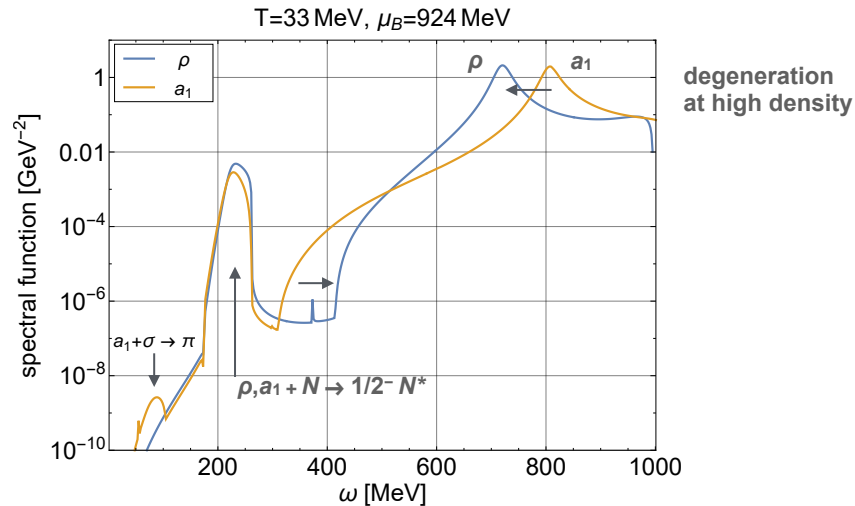
A. Tripolt, Ch. Jung, L.v.S. & J. Wambach, arXiv:2105.00861 [hep-ph]

• imaginary parts:

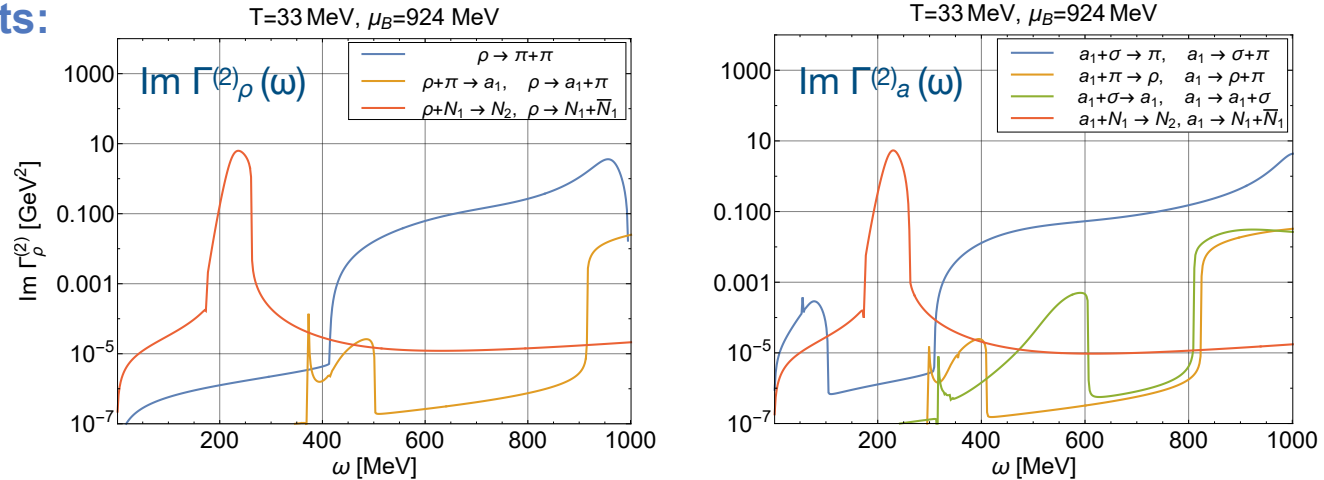


near chiral CEP
inside dense nuclear matter

• (axial-)vector SFs:



• imaginary parts:



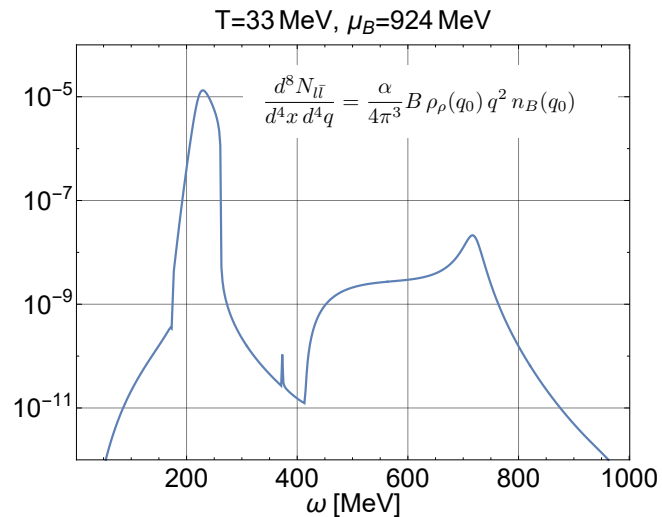
near chiral CEP
inside dense nuclear matter

• electromagnetic SF:

convert to thermal dilepton rate

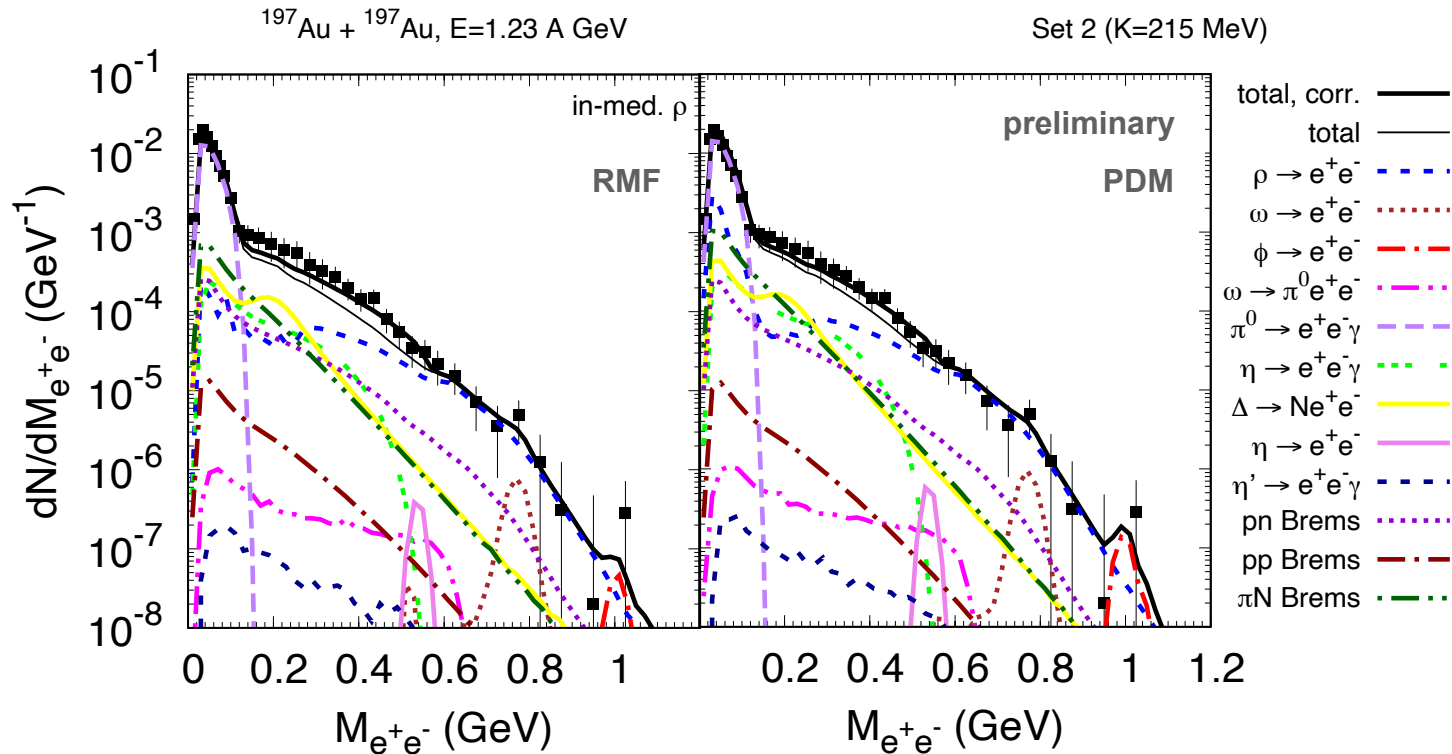
$$\frac{d^8 N_{l\bar{l}}}{d^4 x d^4 q} = \frac{\alpha}{12\pi^3} \left(1 + \frac{2m^2}{q^2}\right) \left(1 - \frac{4m^2}{q^2}\right)^{1/2} q^2 (2\rho_T + \rho_L) n_B(q_0)$$

Weldon, PRD 42 (1990) 2385



- RMF mode: replace mean fields by those of PDM

PDM parameters: M. Kim, S. Jeon, Y.-M. Kim, Y. Kim, & C.-H. Lee, PRC 101 (2020) 064614

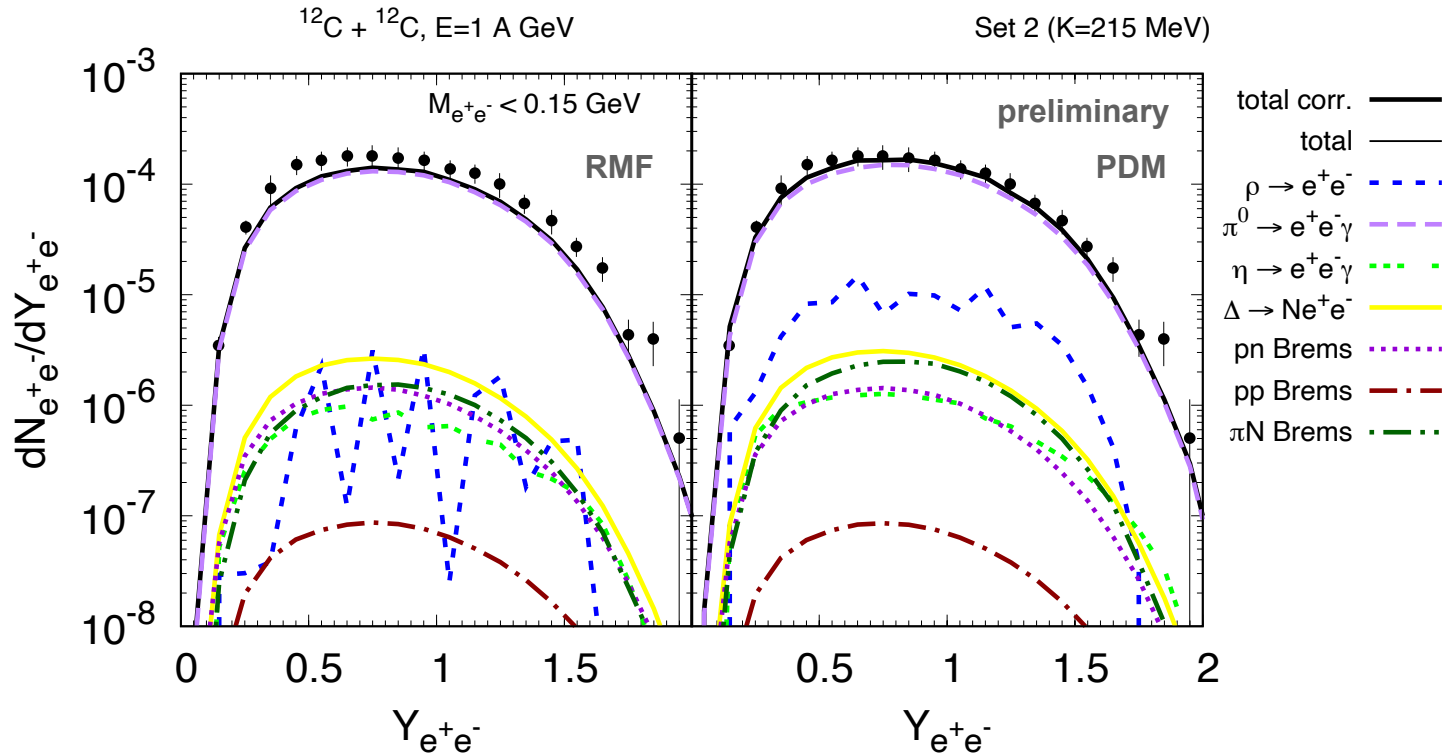


A. Larionov, U. Mosel & L.v.S.,
Phys. Rev. C 102 (2020) 064913

A. Larionov & L.v.S., in preparation

- RMF mode: replace mean fields by those of PDM

Data: Y. Pachmayer, PhD thesis, GU Frankfurt, 2008

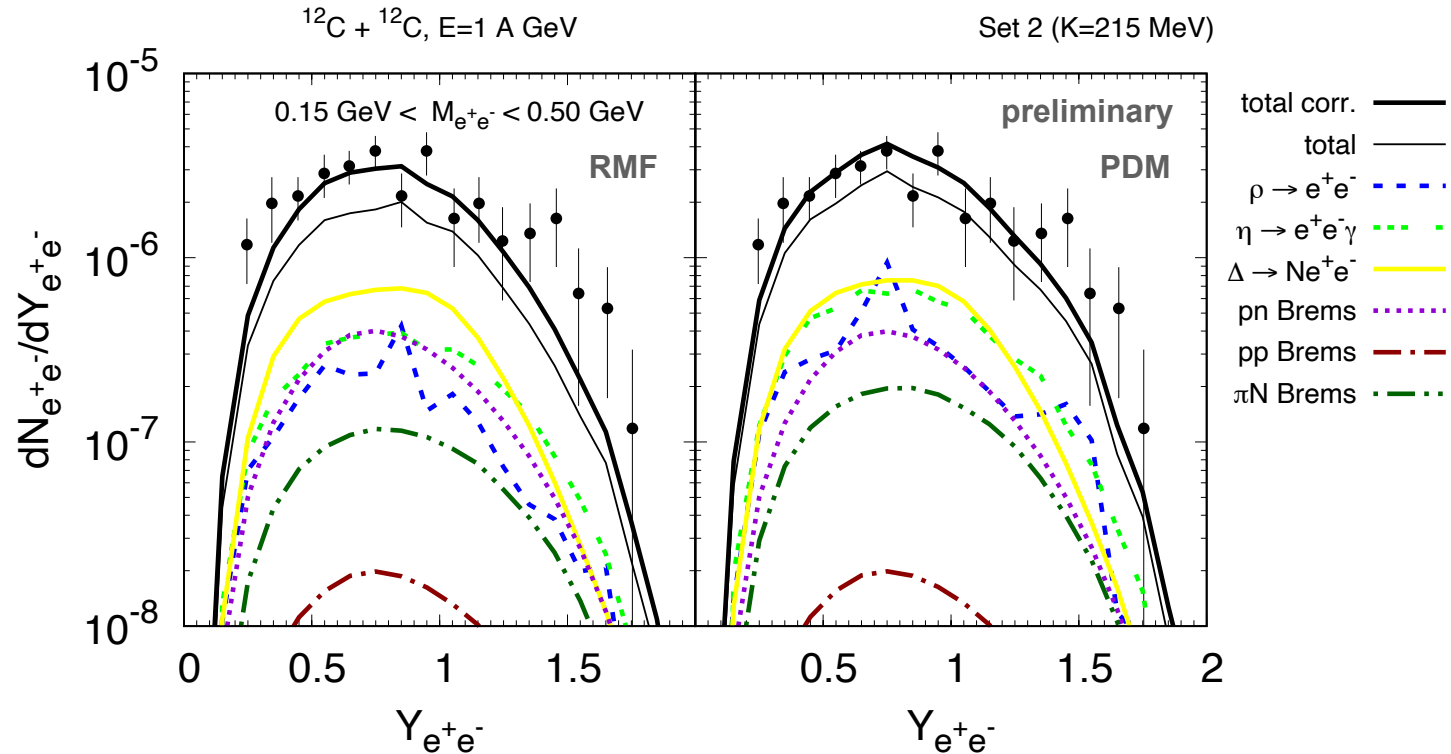


A. Larionov, U. Mosel & L.v.S.,
Phys. Rev. C 102 (2020) 064913

A. Larionov & L.v.S., in preparation

- RMF mode: replace mean fields by those of PDM

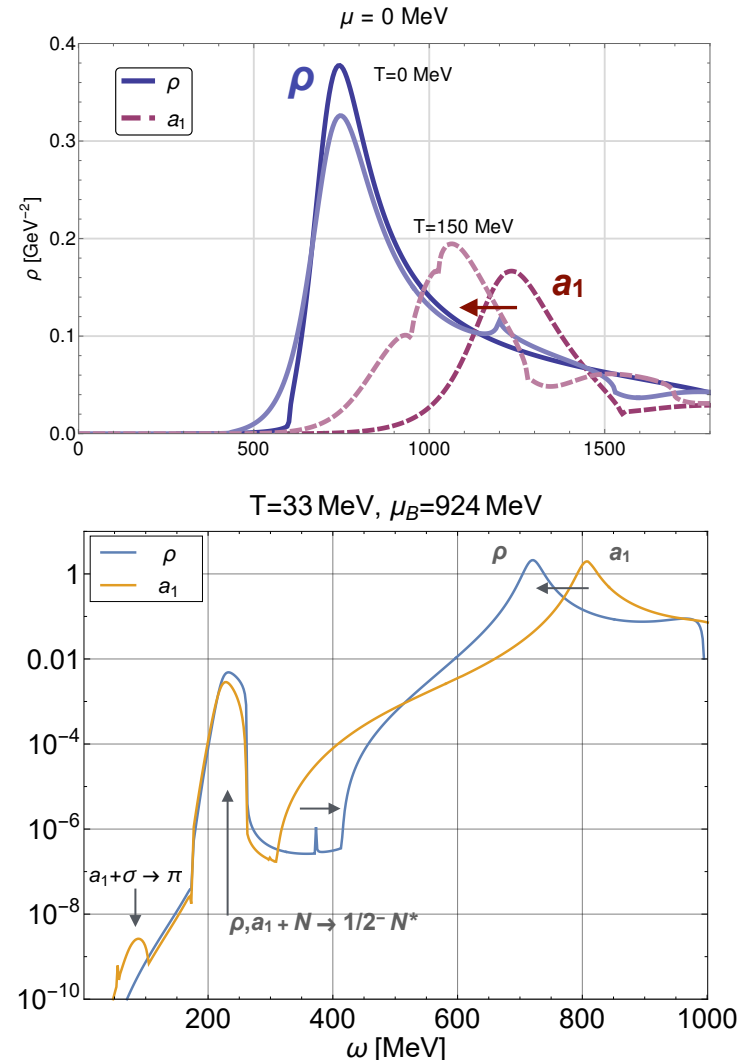
Data: Y. Pachmayer, PhD thesis, GU Frankfurt, 2008



A. Larionov, U. Mosel & L.v.S.,
Phys. Rev. C 102 (2020) 064913

A. Larionov & L.v.S., in preparation

- Spectral functions from analytically contd. aFRG flows effective theories (chiral, linear)
- Vector and axial-vector SFs at finite T and μ melting-rho scenario
- Electromagnetic spectral function U(1) gauging, mixing
- Fermionic spectral functions use for baryonic SFs in dense matter
- (Axial-)Vector SFs in nuclear matter, parity doubling effective hadronic theory with chiral PT
- Parity-doublet chiral MFT in GiBUU enhanced low energy ρ and η signals



- **parity-doublet model with fluctuating ω and ρ**
symmetric nuclear and neutron matter

- **ρ - a_1 mixing and signatures of CEP in HIC**

electromagnetic \rightarrow dilepton rates
weak \rightarrow neutron star mergers...

- **self-consistent spectral functions**

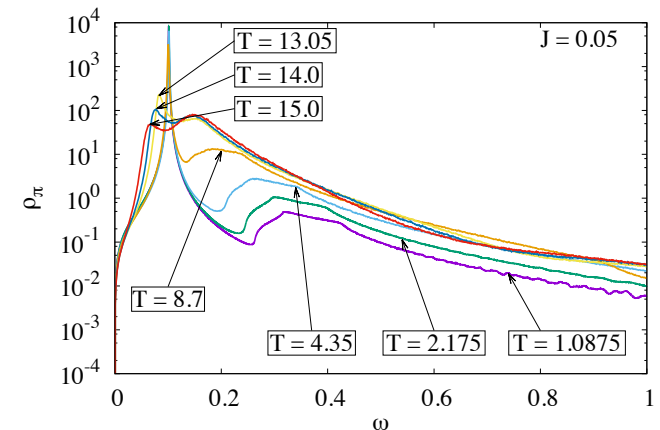
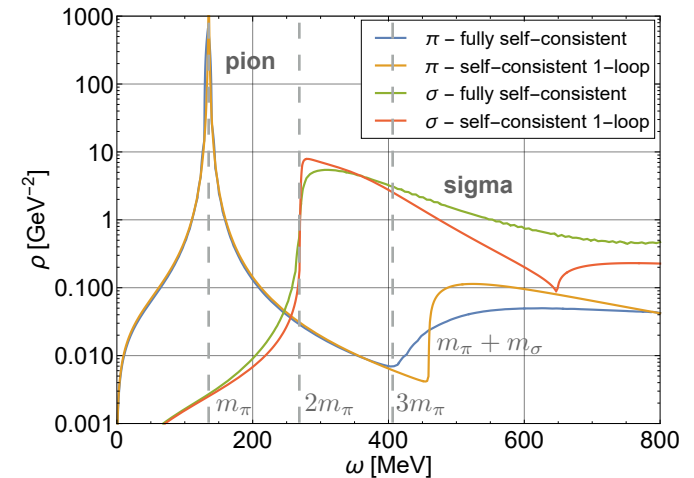
O(4)-model, in preparation

- **universal critical SFs from classical-statistical simulations**

O(4)-model, S. Schlichting, D. Smith & Lv.S,
NPB 950 (2020) 114868

universal dynamic scaling functions,
S. Schlichting, D. Schweitzer & Lv.S, NPB 960 (2020) 115165

- **Gaussian-state approximation, real-time FRG...**



Thank you for your attention!